

Lifetimes of ultralong-range Rydberg molecules in a dense Bose-Einstein condensate

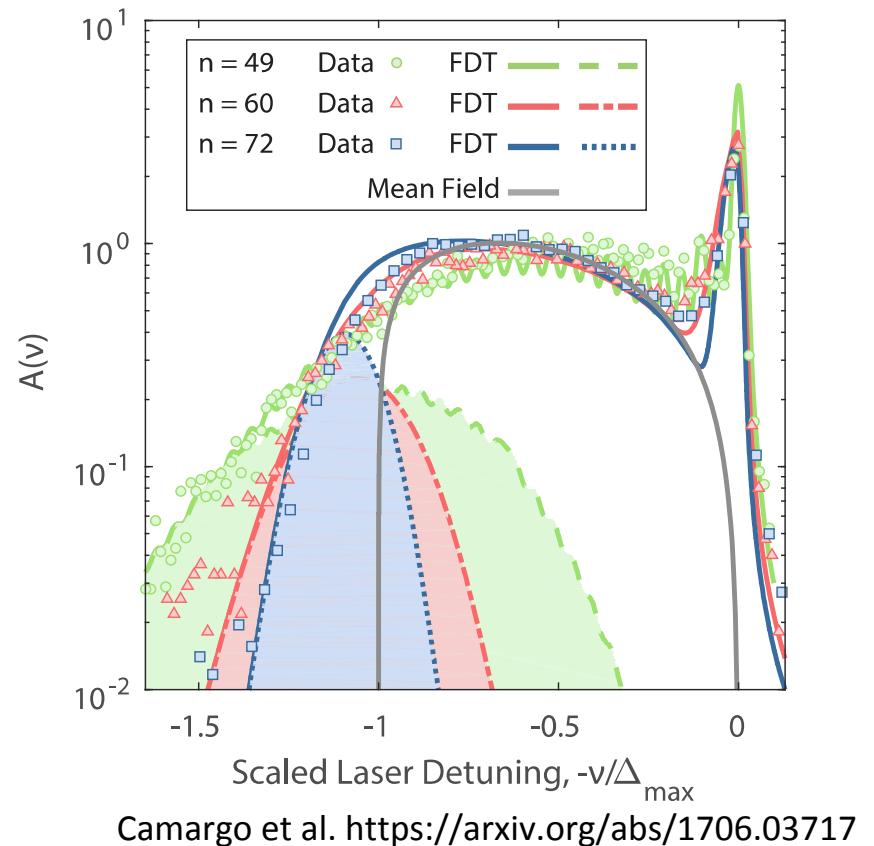
Experiment: **Joseph D. Whalen**, F. Camargo, R. Ding, T. C. Killian, F. B. Dunning

Theory: J. Pérez-Ríos, S. Yoshida, J. Burgdörfer



Cold Rydberg gases

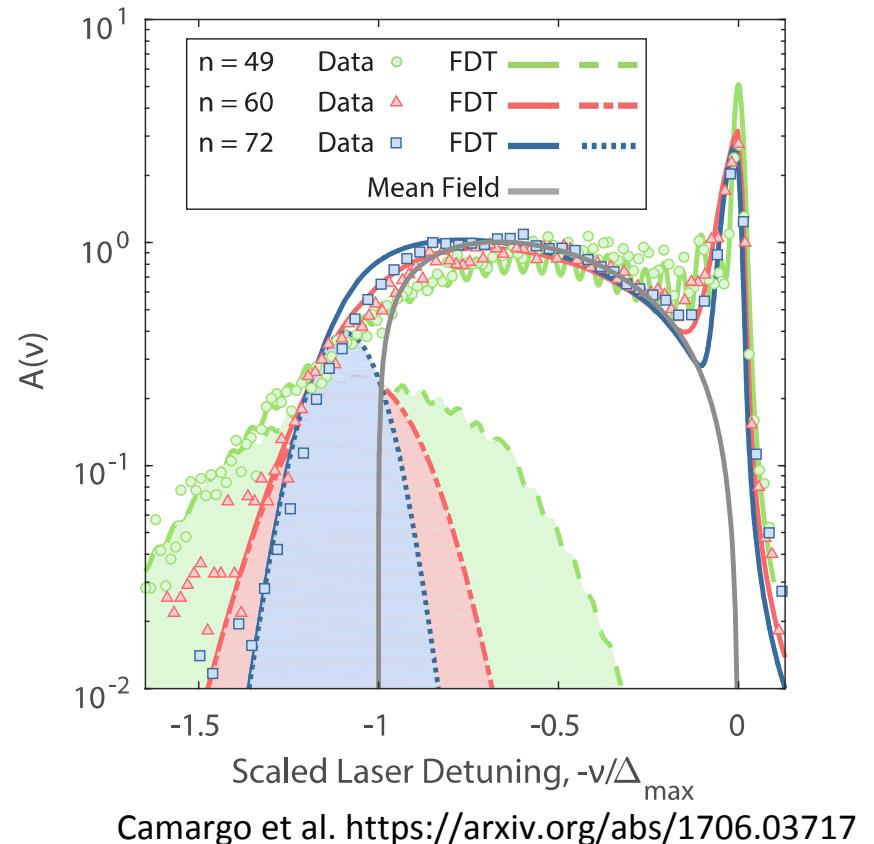
- Using highly excited atoms embedded in BEC or cold thermal gas we can study:
 - Tunable long-range dipolar interactions
 - Quantum Information
 - Electron-atom scattering
 - Impurity physics and polarons
 - Novel bound states



Camargo et al. <https://arxiv.org/abs/1706.03717>

Cold Rydberg gases

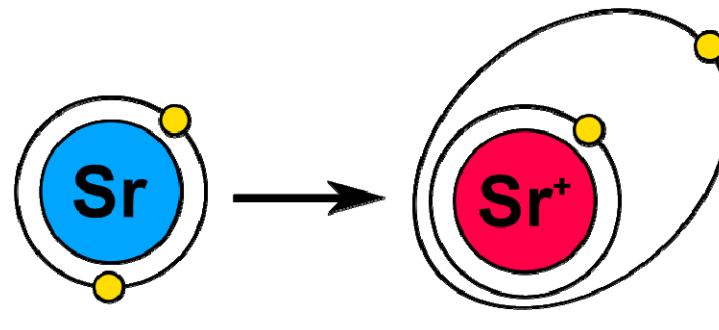
- Using highly excited atoms embedded in BEC or cold thermal gas we can study:
 - Novel bound states
- Question: How long do these excitations live in a dense gas?
- Motivated by studies in Rb
 - Low density – 2011 JPB **44** 184004
 - High density – PRX **6**, 031020
- Our work in Sr
 - Low density – PRA **93**, 022702, 2015
 - High density (this talk) – arXiv:1707.02354



Camargo et al. <https://arxiv.org/abs/1706.03717>

Rydberg atoms

- Atomic state of high principal quantum number n
- Electron orbits far away from ionic core
- Highly exaggerated properties that scale with n
 - Massive size
 - Long lifetime
 - Huge dipole moment



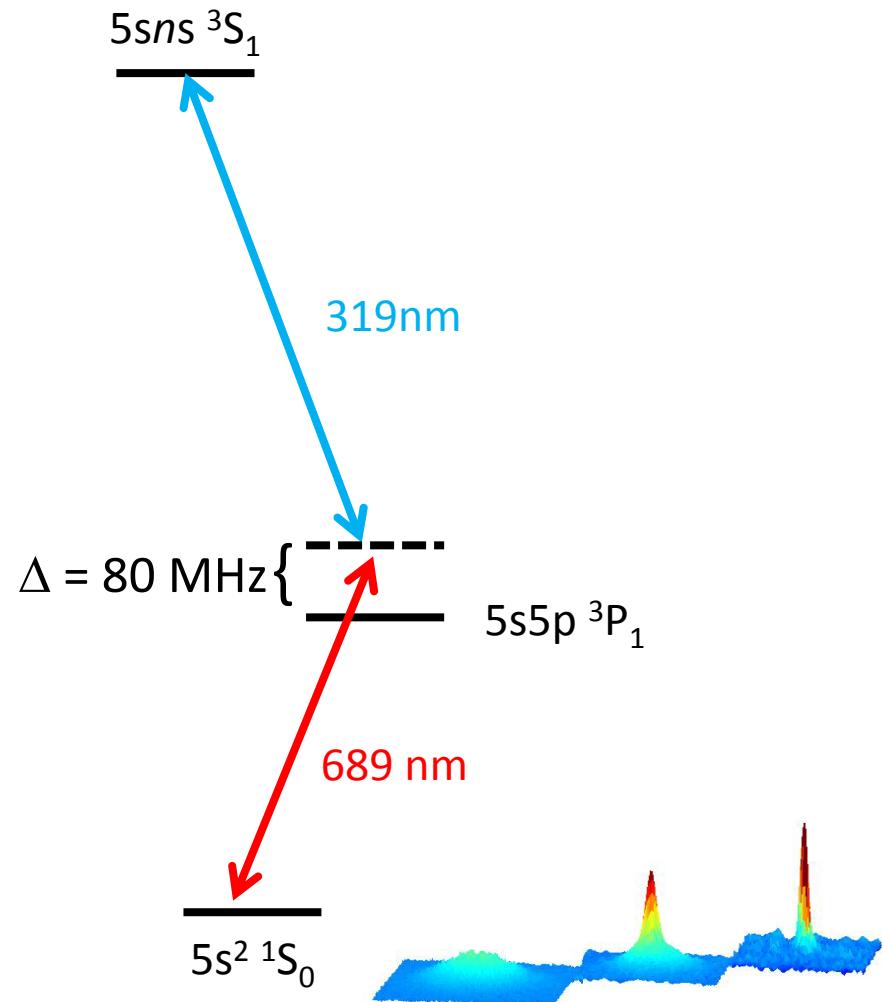
$5s^2 \ ^1S_0$

$5sns \ ^3S_1$

Property	n Dependence	Sr $5s50s \ ^3S_1$
Binding Energy	n^{-2}	50 cm^{-1}
Radiative Lifetime	n^3	$38 \mu\text{s}$
Rydberg Radius $\langle r \rangle$	n^2	200 nm
Ionization Field	n^{-4}	65 V/cm
Polarizability	n^7	$56 \text{ MHz cm}^2/\text{V}^2$

Experimental method

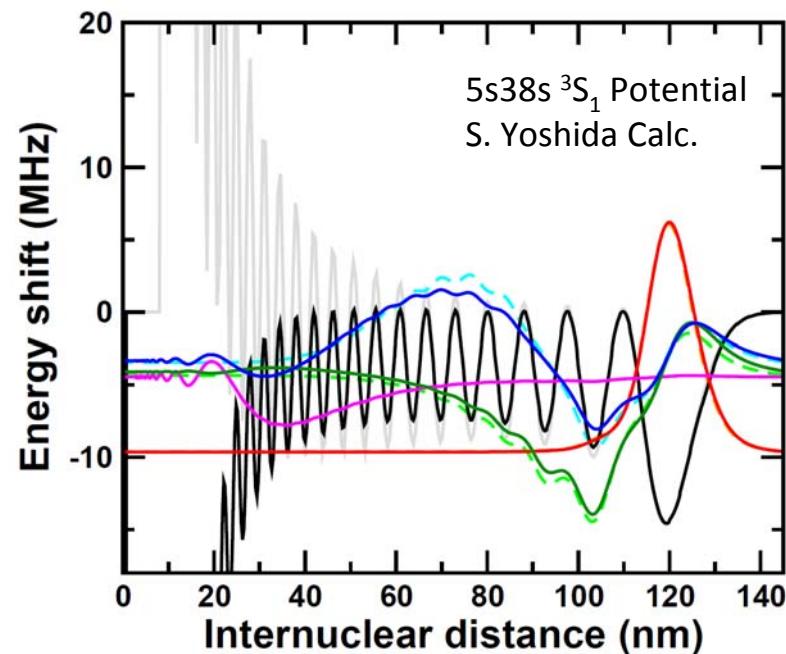
- Prepare BEC of 4×10^5 ^{84}Sr atoms
 - 1064nm ODT, harmonic potential
 - Peak density = $4 \times 10^{14} \text{ cm}^{-3}$
- Excite to triplet states via ${}^3\text{P}_1$ intercombination line
- Obtain spectrum by scanning UV photon energy
- Hold excited Rydberg atoms for 1-100 μs
- Detect Rydberg population on MCP with selective field ionization (SFI)



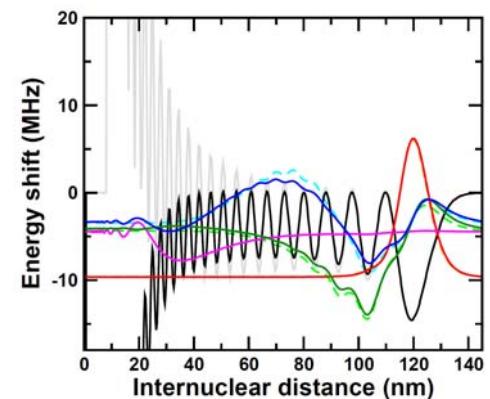
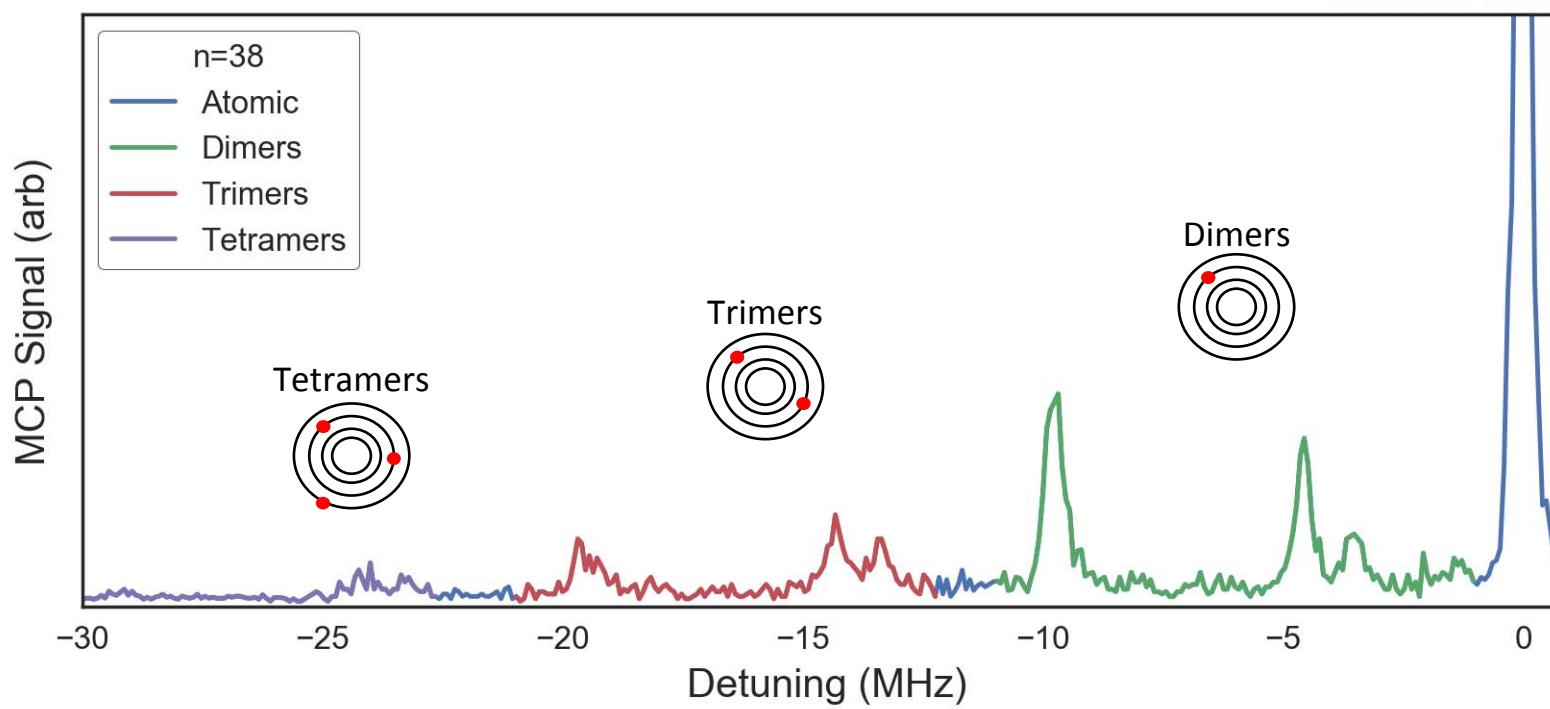
Rydberg molecules

$$V(\vec{R}) = \frac{2\pi\hbar^2 a_s(k)}{m_e} \left| \psi_{nl}(\vec{R}) \right|^2 + \frac{6\pi\hbar^2 a_p(k)^3}{m_e} \left| \vec{\nabla} \psi_{nl}(\vec{R}) \right|^2$$

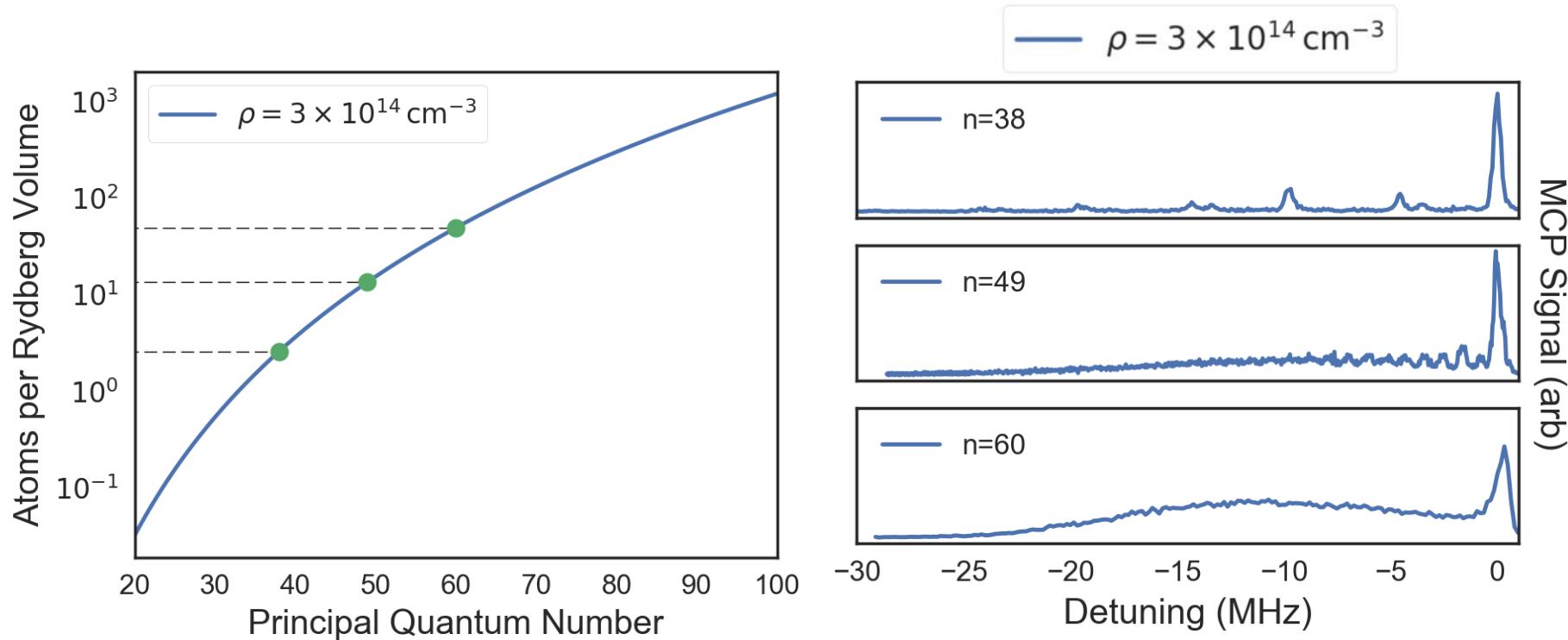
- Exciting a Rydberg atom in a cold dense gas
- Scattering between Rydberg electron and ground-state atoms can form bound states
- Predicted by Greene et. al - Phys. Rev. Lett. **85**, 2458, 2000
- First observed in 2009 by Pfau group - *Nature* **458**, 1005-1008, 2009



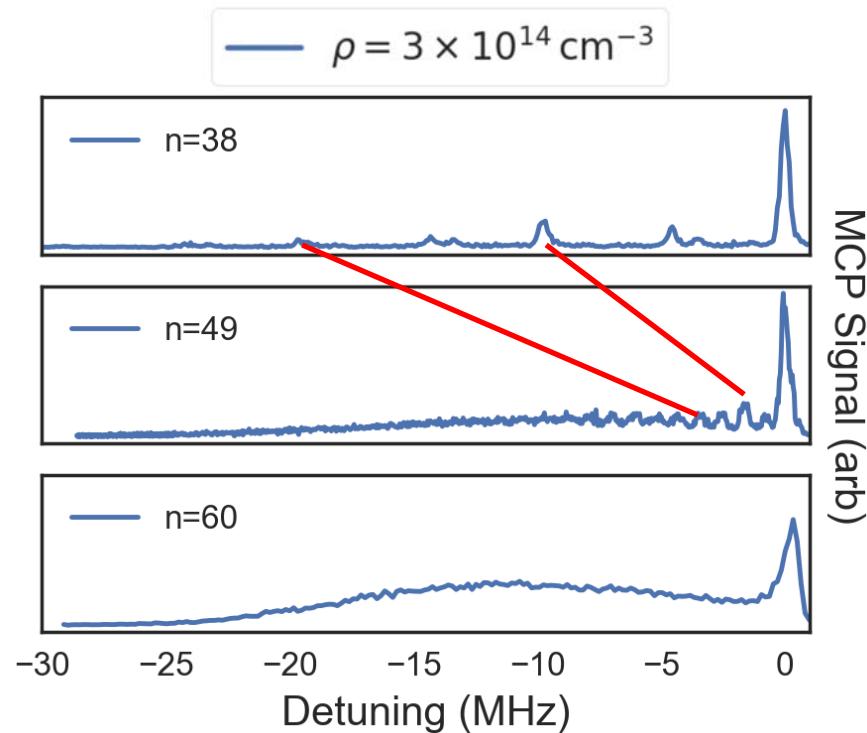
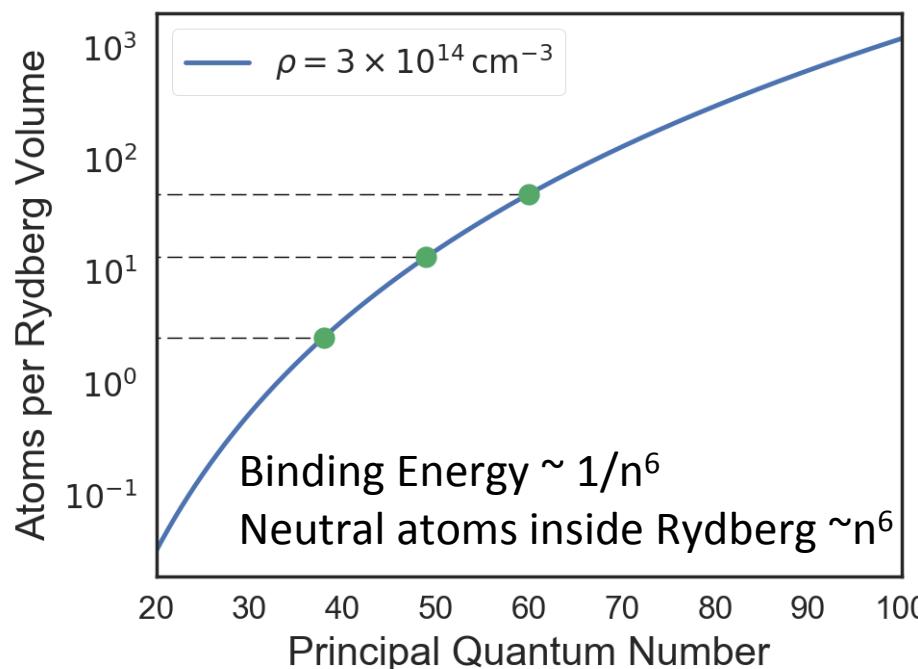
Rydberg molecules



Increasing principal quantum number



Increasing principal quantum number

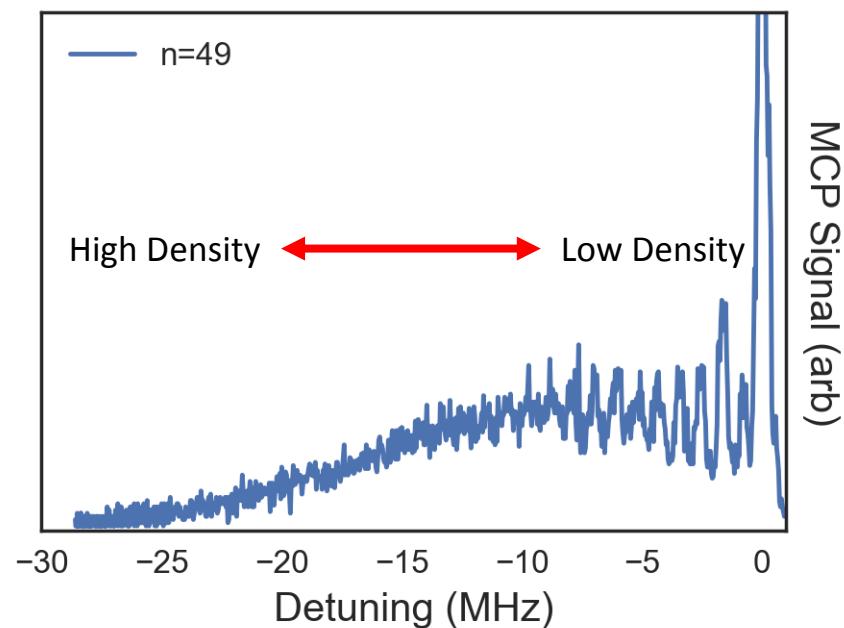


Laser detuning selects local density

- Mean field predicts shift in resonance proportional to local density

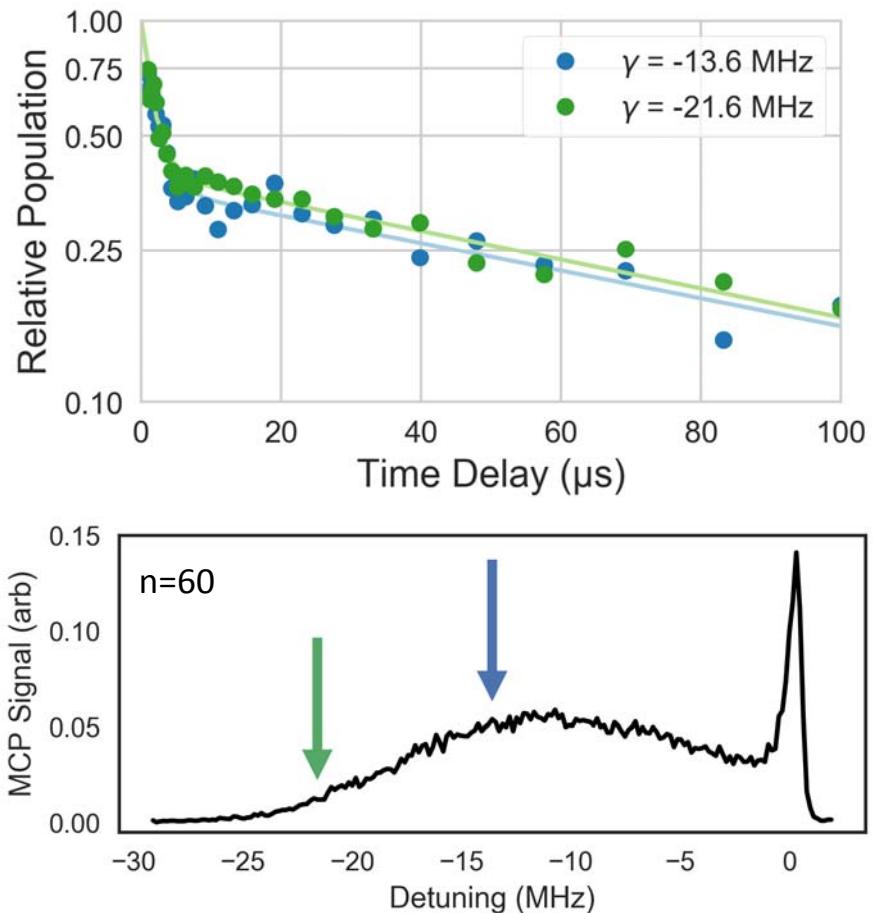
$$\Delta E = \frac{2\pi\hbar^2}{m_e} A_s(k) \rho$$

- Large red laser detuning implies excitation at higher density.
- With a narrow laser we can choose excitation density precisely



Population dynamics

- Measure decay of Rydberg population at a fixed laser detuning \square (fixed density)
- At high densities we observe fast initial loss and slow radiative decay at later times
- Indicative of molecular loss channels other than spontaneous emission
- Possible mechanisms
 - Sr_2^+ production
 - L-changing collisions

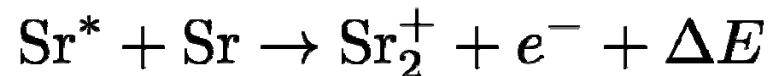
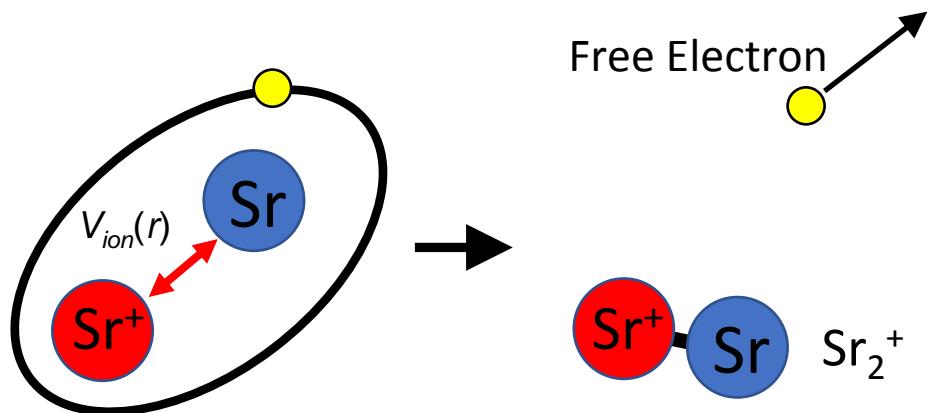
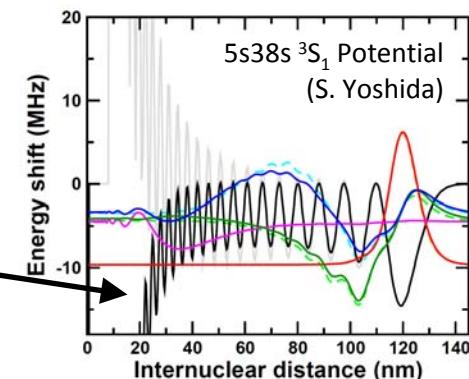


Loss mechanisms – Sr_2^+ production

- Reaction initiated by core-ion + nearest ground-state atom polarization potential
- Electron takes away binding energy and escapes the trap
- Observed as loss in molecular signal
- Rb_2^+ directly observed in other experiments

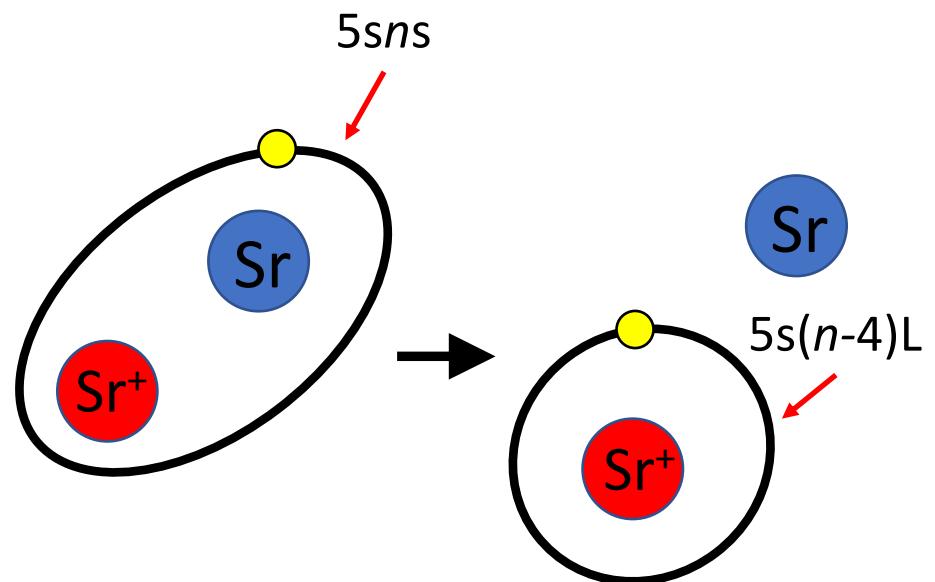
$$V_{ion}(r) \propto -\frac{C_4}{r^4}$$

Dominant at short range



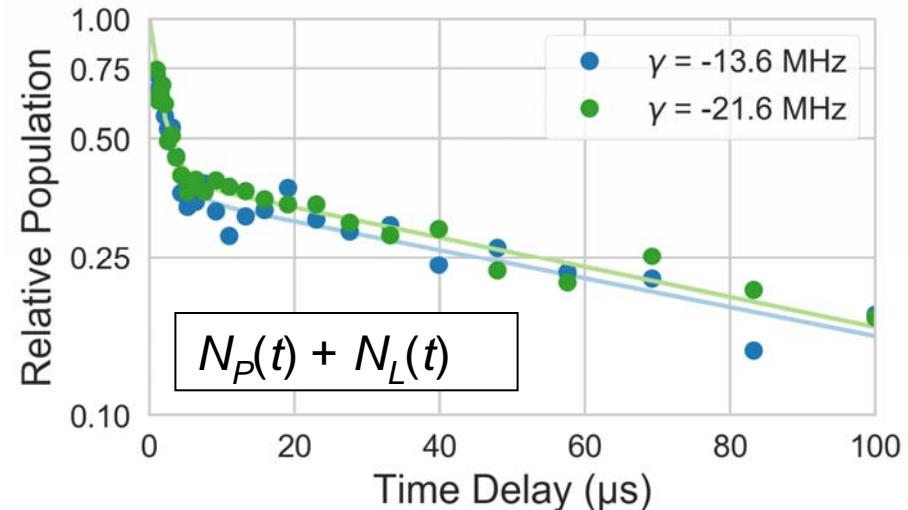
Loss mechanisms – L changing collisions

- Reaction initiated by core-ion + nearest ground-state atom polarization potential
- Rydberg atom in high L is ejected from the trap
- Ejected Ryds. undergo no more collisions, $1/e$ lifetime scales as n^3
- Observed as residual signal and characteristic shift in SFI profile



Model

- Results consistent with simple rate equation model
- Three component fit with rate constants
 - Γ_{AI} – Associative ionization rate
 - Γ_L – L-changing collision rate
 - Γ_R – Radiative decay rate
- Fit sum of parent atoms, N_p , and L-changed atoms, N_L , to total population



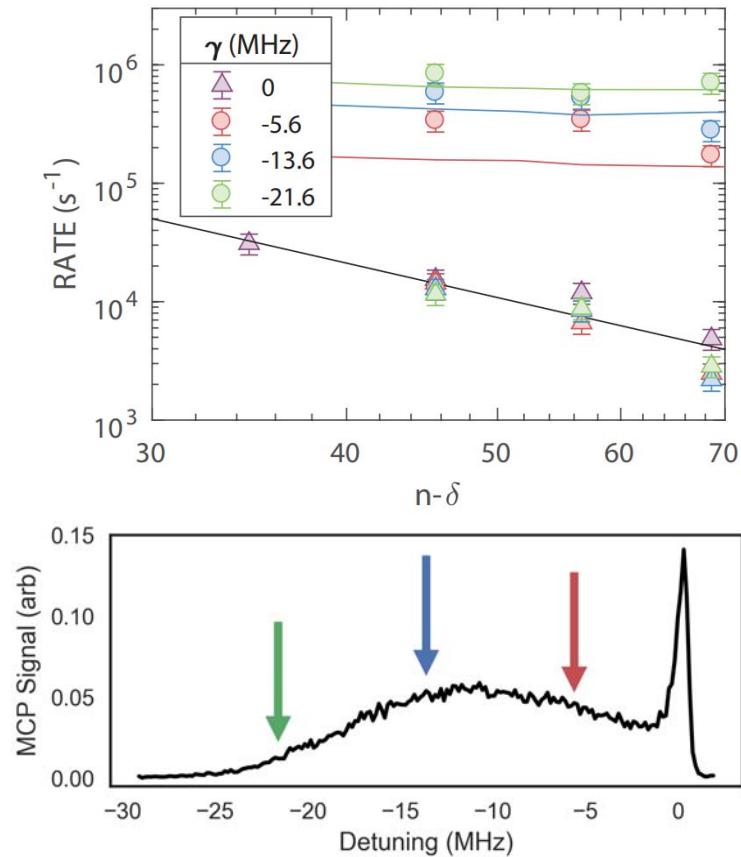
$$N_p(t) = N_0 e^{-(\Gamma_{AI} + \Gamma_L + \Gamma_R)t}$$

$$N_L(t) = N_0 \frac{\Gamma_L}{\Gamma_{AI} + \Gamma_L} e^{-\Gamma_R t} \left[1 - e^{-(\Gamma_{AI} + \Gamma_L)t} \right]$$

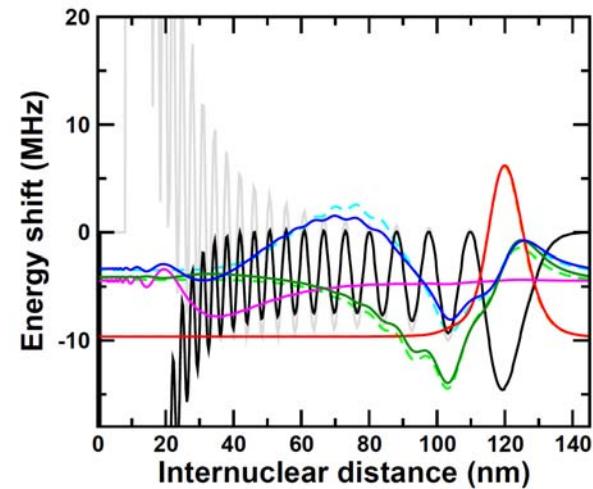
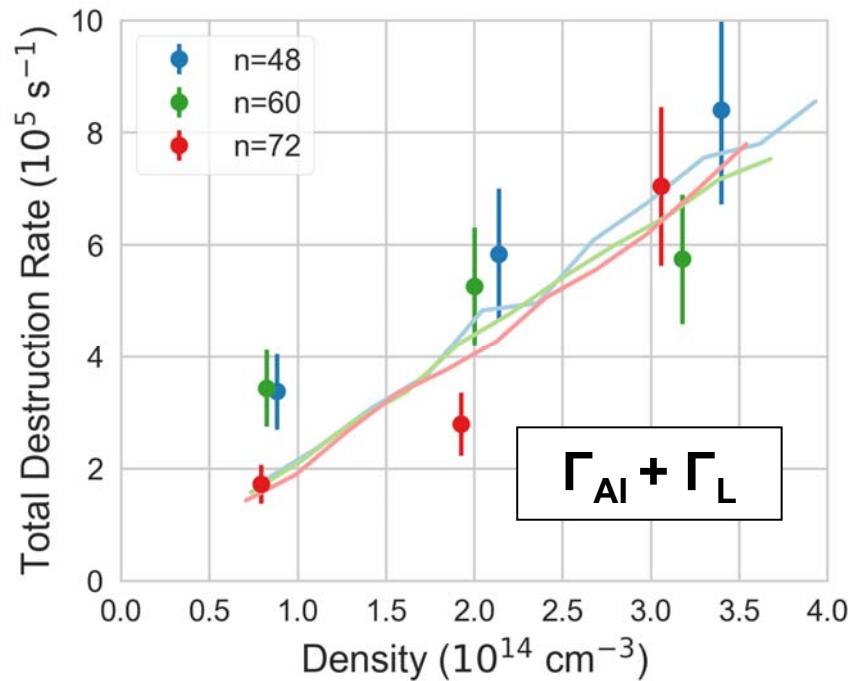
n dependence of rates

- Destruction rates largely independent of n when many atoms present in Ryd. orbital.
- Indicates presence of one ground-state atom near the core-ion leads to destruction of Ryd. molecule.
- Radiative lifetime negligible compared to collision rates

$$\Gamma_{\text{AI}} + \Gamma_{\text{L}}$$



Density dependence of rates



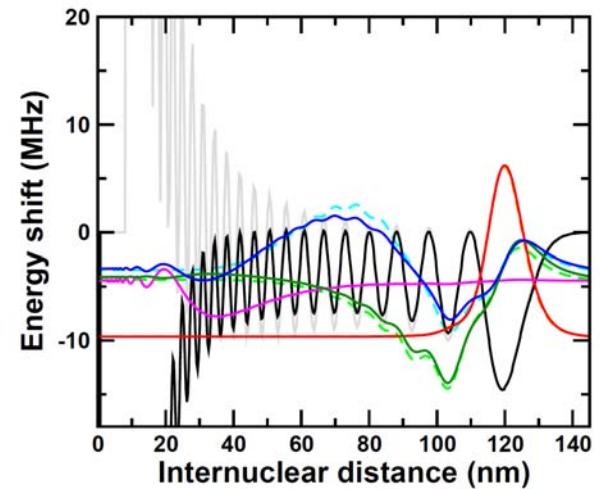
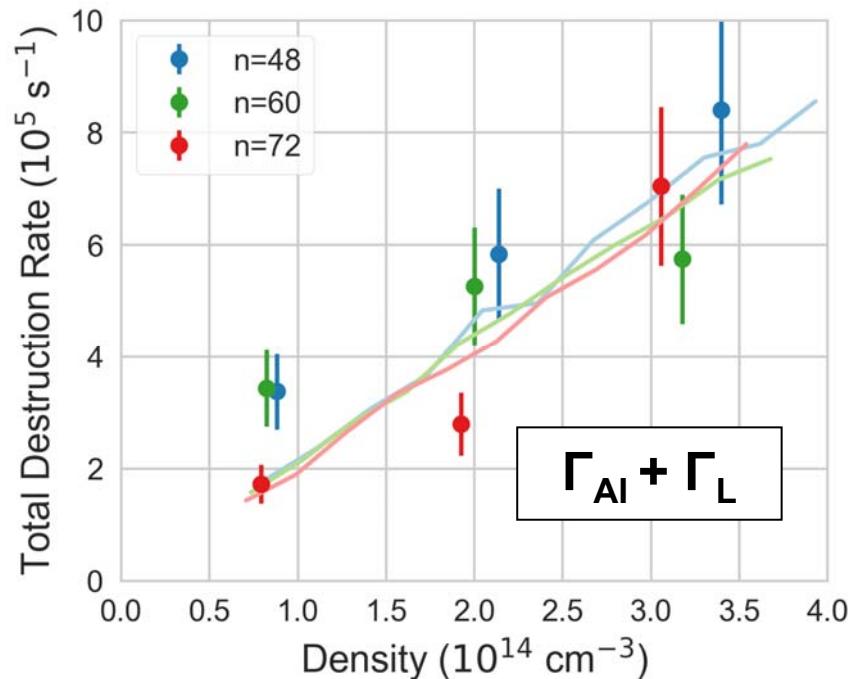
Collision time t scales inversely with density in the limit $E_{\text{coll}} \rightarrow 0$

$$t = \int_{r_i}^{r_f} \frac{dr}{v(r)} = \int_{r_i}^{r_f} \frac{dr}{\sqrt{\frac{2}{\mu} \left(E_{\text{coll}} + \frac{C_4}{r^4} \right)}}$$

$$t \propto - \int_{r_i}^{r_f} r^2 dr \propto r_i^3$$

$$r_i \propto \rho^{-1/3} \implies t \propto 1/\rho$$

Density dependence of rates



Collision time t scales inversely with density in the limit $E_{\text{coll}} \rightarrow 0$

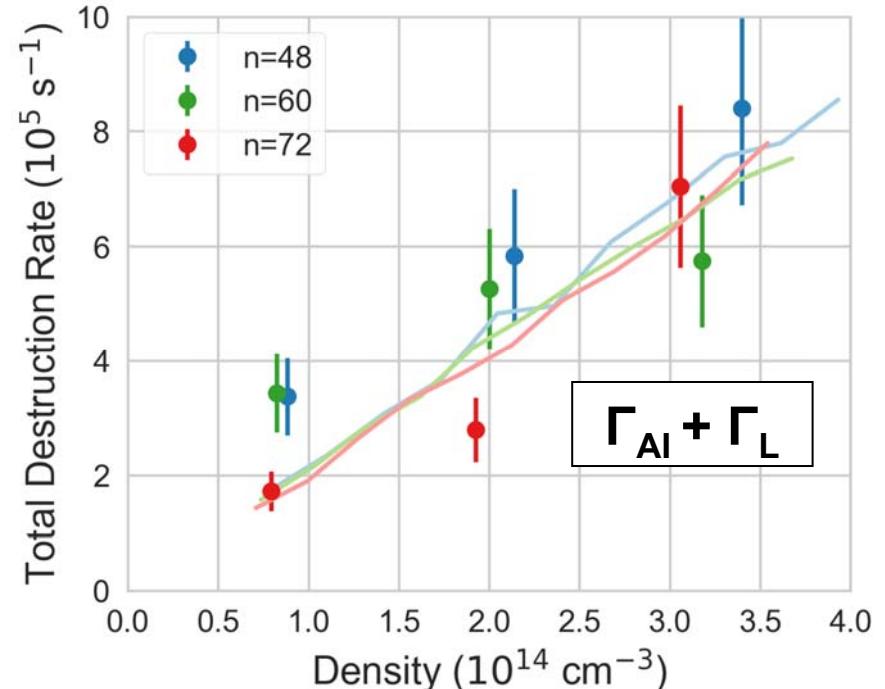
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Conclusions and outlook

- Lifetime of Rydberg excitation in dense gases is limited by collision time with nearest neighbor
- Two likely decay mechanisms
 - Sr_2^+ production
 - L-changing collisions
- Described by simple kinematic model
- Should rate decrease at higher n ?
 - Electron spends less time near the nucleus
- Do particle statistics affect molecular formation / lifetimes?
 - Experiments with ^{87}Sr underway



Thank You

Killian Group (Rice)

Rydberg:

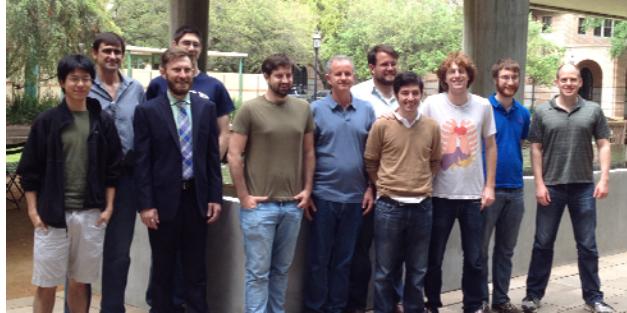
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Plasma:

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Barry Dunning



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July 26, 2017



Ultracold Atoms and Plasmas

Signature of L-changing collisions

