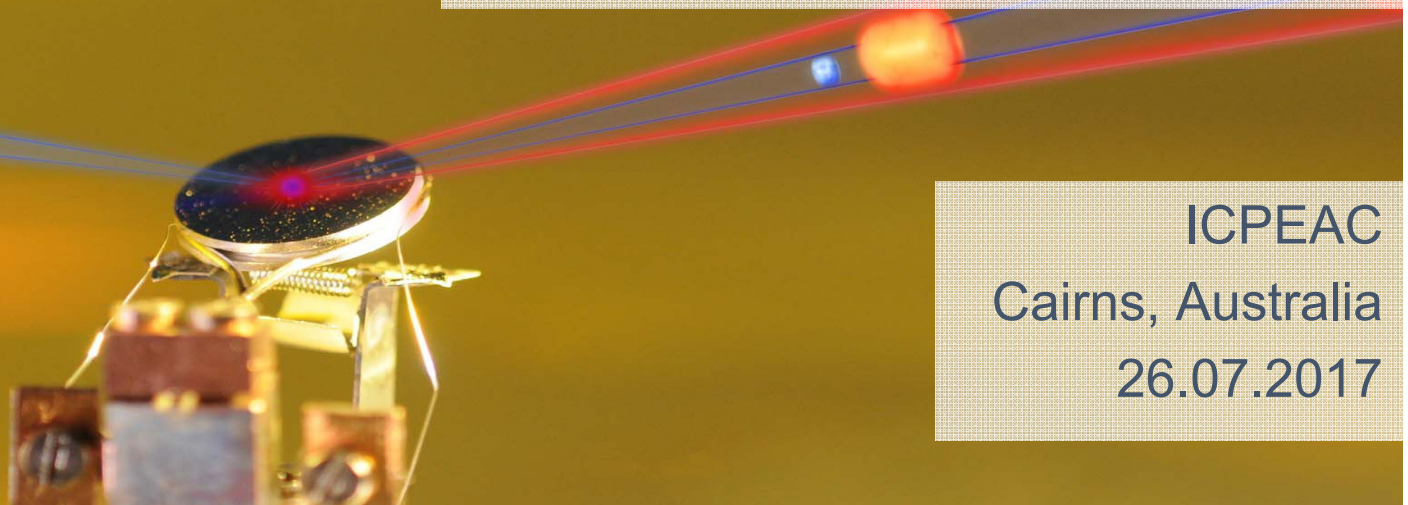


Attosecond Electron Dynamics on Surfaces and Layered Systems

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Technische Universität Munich
Max-Planck-Institute of Quantum Optics

ICPEAC
Cairns, Australia
26.07.2017





Scope of ICPEAC



The ICPEAC conference attracts several hundred scientists from all over the globe who work actively in the fields **of collisions involving photons, electrons, ions, atoms, molecules, clusters, surfaces, and exotic particles**. Recent expansions of the conference scope include **collision with cold targets and attosecond science**.

electrons – atoms (ions): High-order Harmonic Generation

photons – surfaces: photoelectroc effect

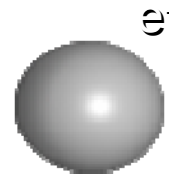
Attosecond science



attosecond science:

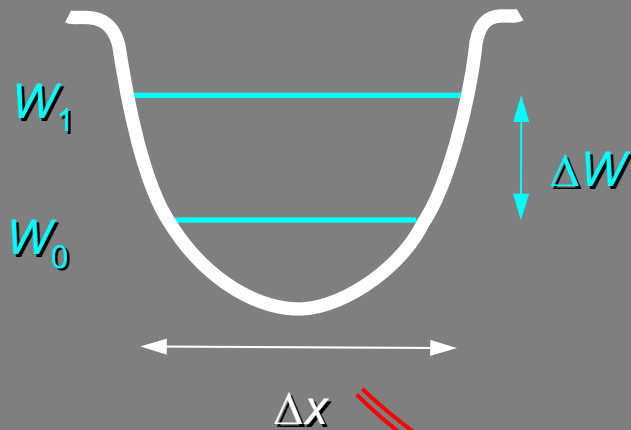


Observation and steering of electronic processes





the time & length scales of microscopic motion
are connected by the laws of **quantum mechanics**

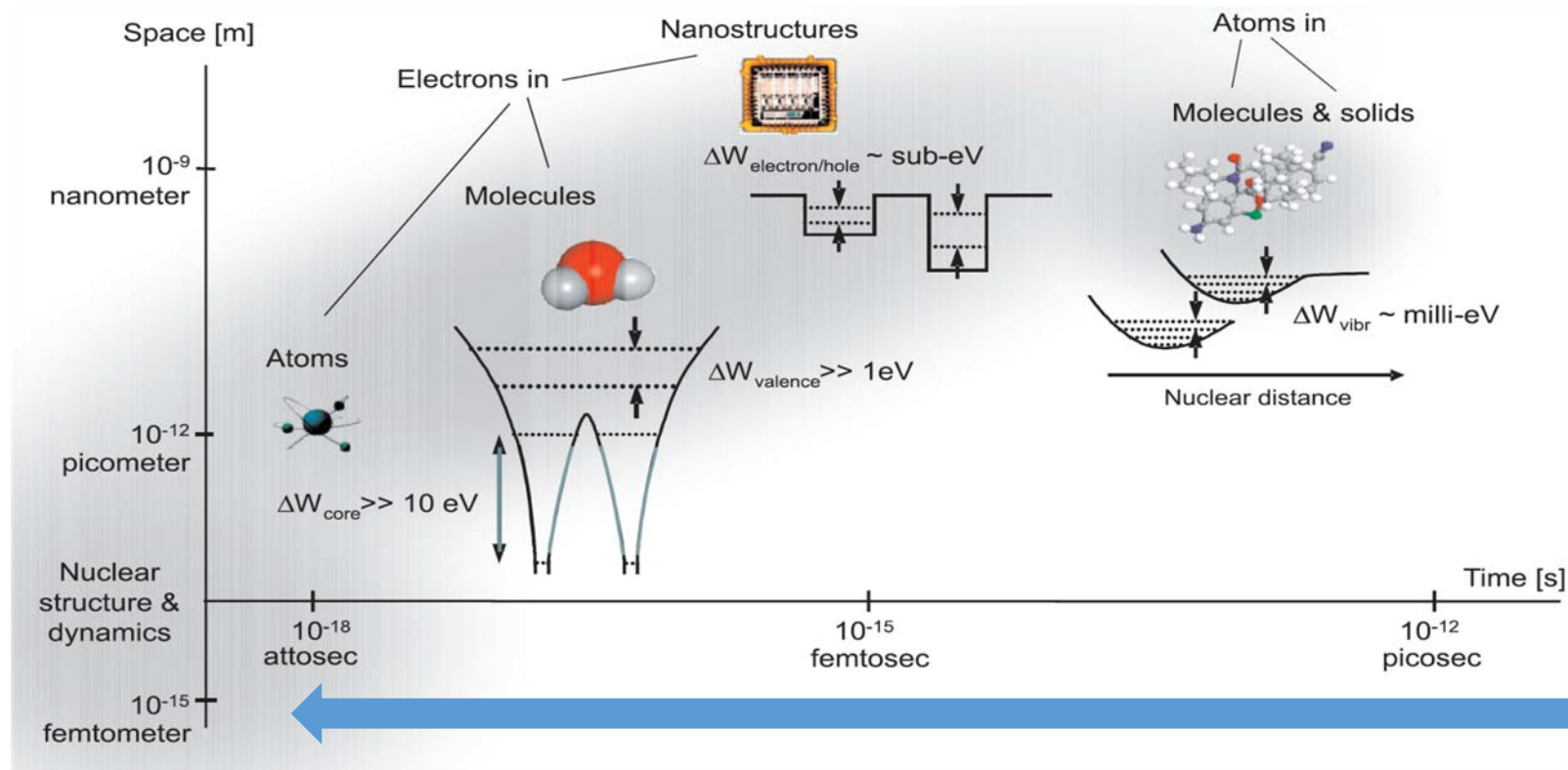


$$\Delta t \sim \frac{h}{\Delta W}$$

$$\Delta W \sim f\left(\frac{1}{\Delta x}\right)$$



Timescales of electronic dynamics in ever smaller systems





Limit in resolution: exposure time





microsecond - photography



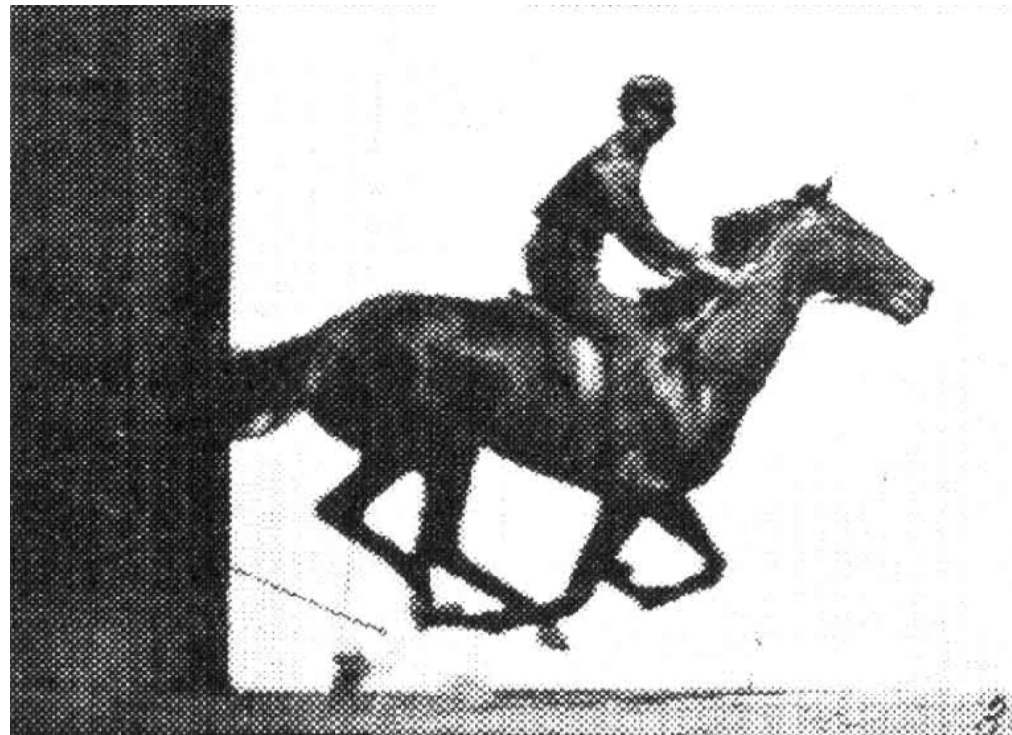
bullet in a glass block



1878: E. Muybridge, Stanford



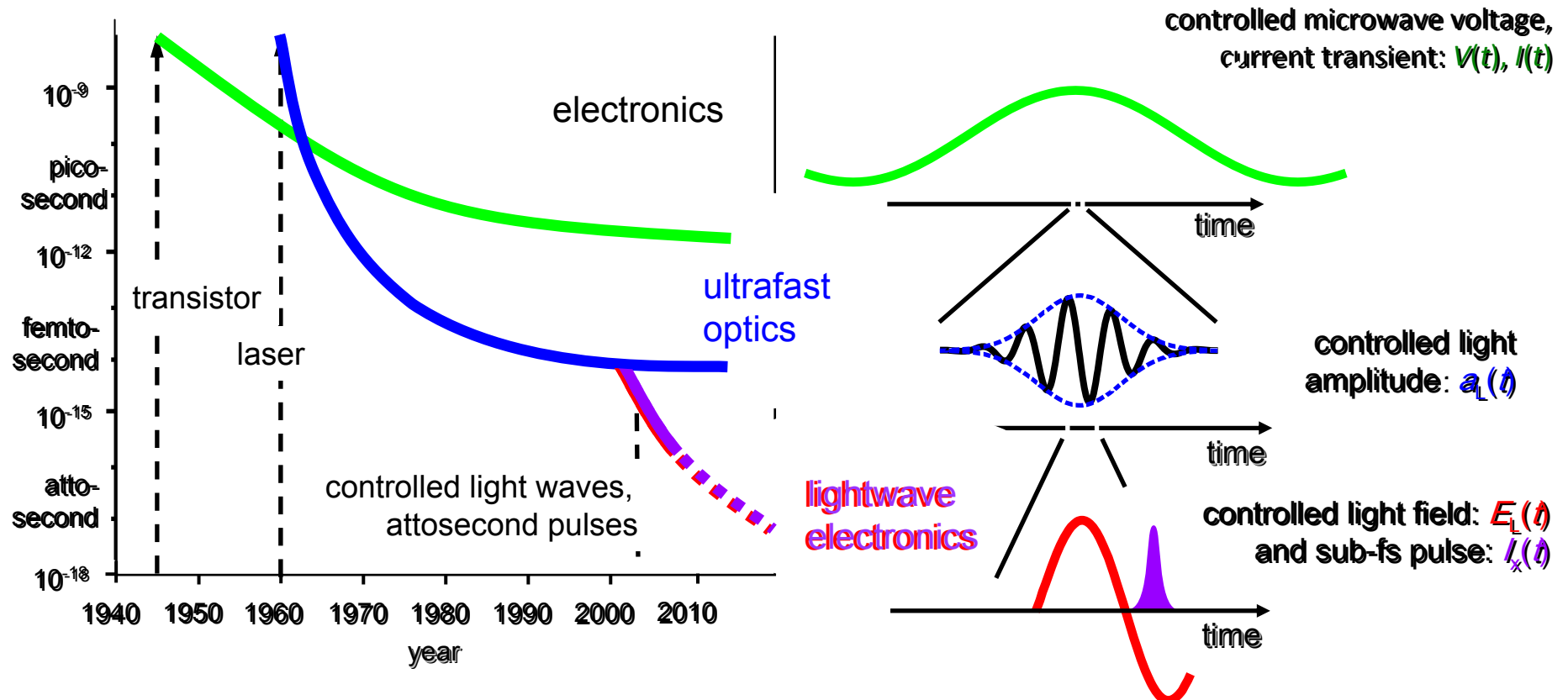
Reconstruction of the process (“slow-motion”)
by stringing together snapshots at different phases of the process

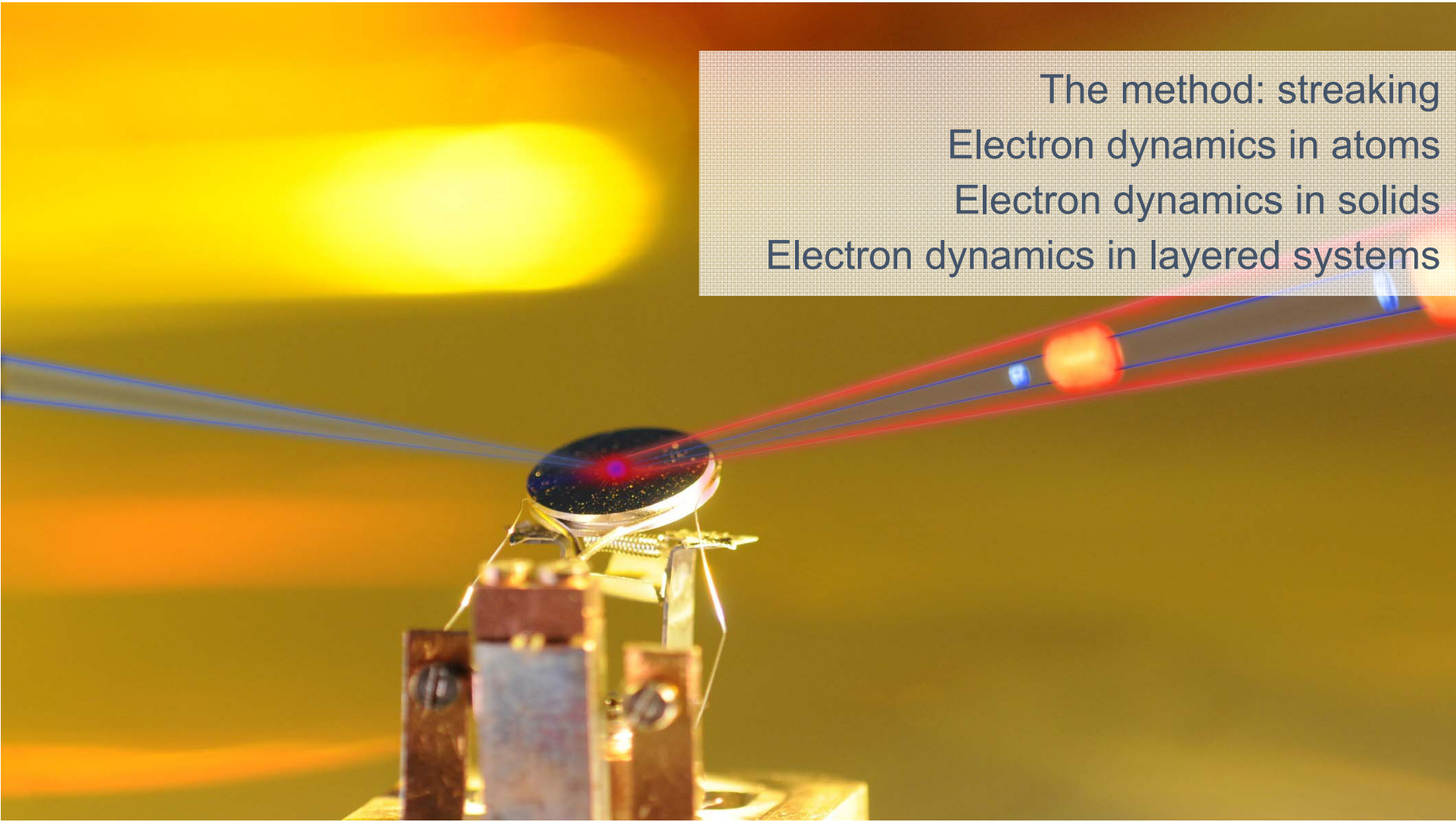


E. Muybridge, *Animals in Motion*, ed. by L. S. Brown (Dover Publ. Co., New York 1957)



from superfast electronics to lightwave driven electronics





The method: streaking

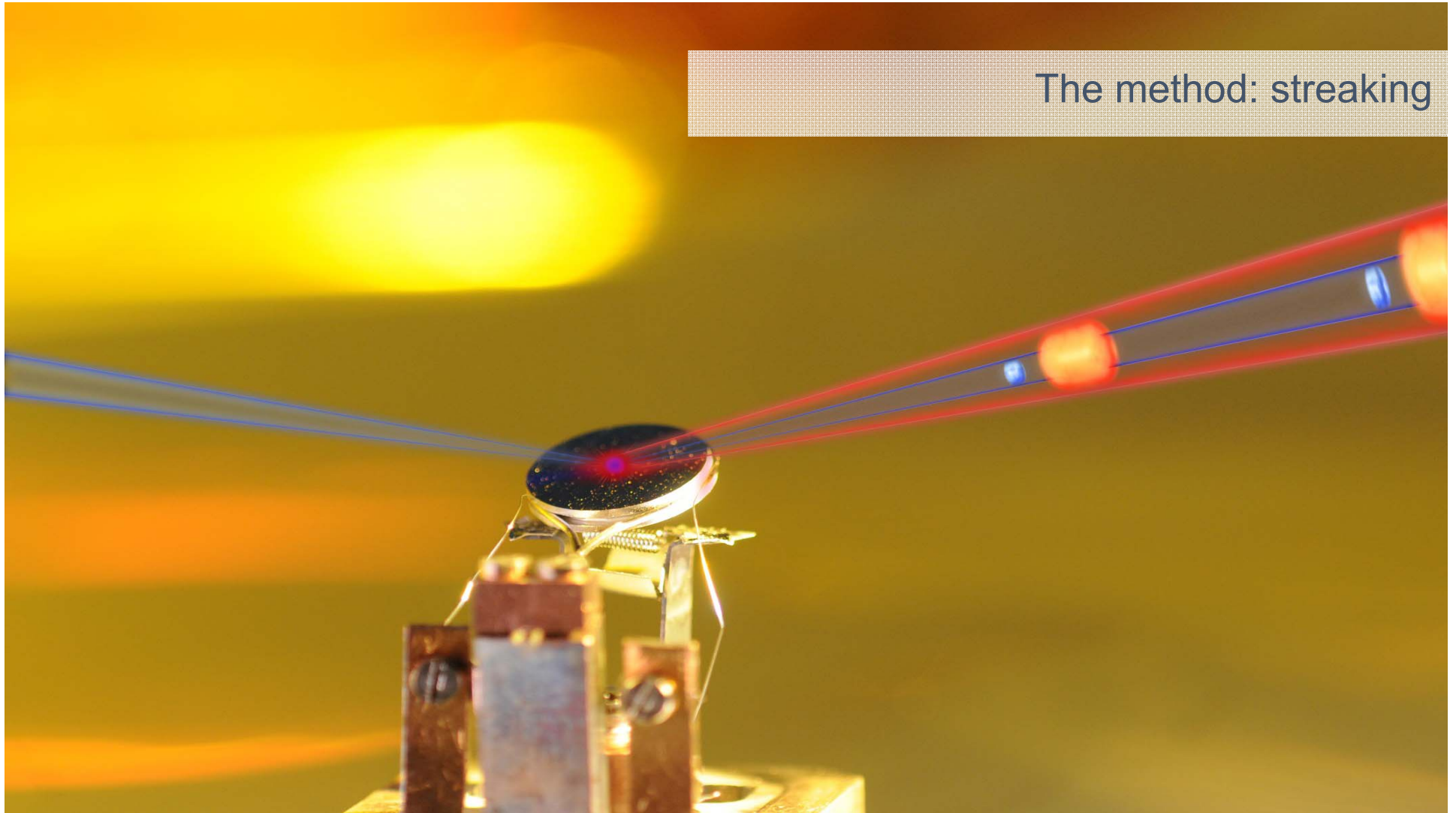
Electron dynamics in atoms

Electron dynamics in solids

Electron dynamics in layered systems

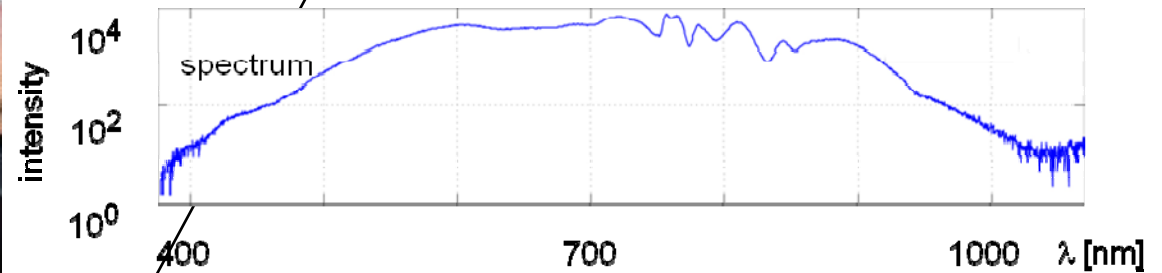
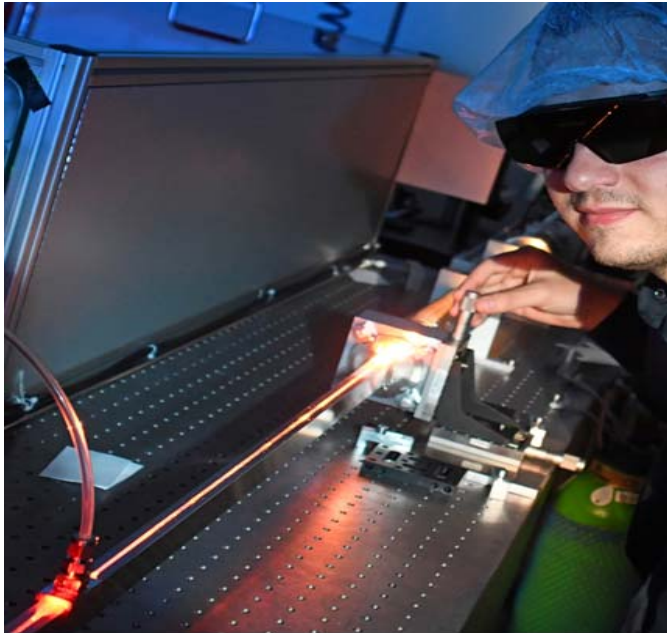
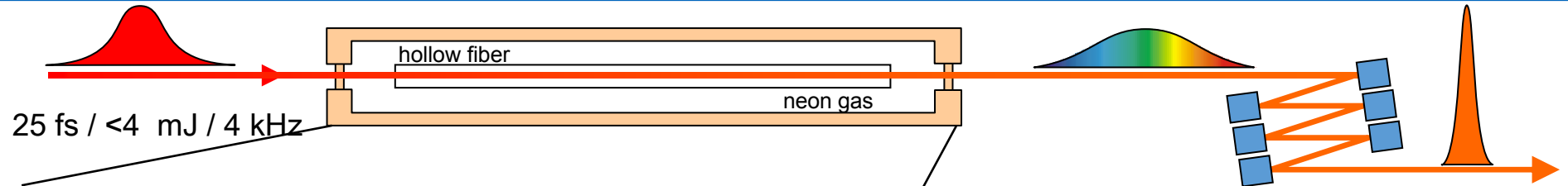
The image shows a streak camera setup. A bright yellow light source is in the background. A circular detector is mounted on a metal frame. Two blue lines represent the streaking process, showing a red dot (electron) moving along a path. The background is a warm, yellowish-orange gradient.

The method: streaking



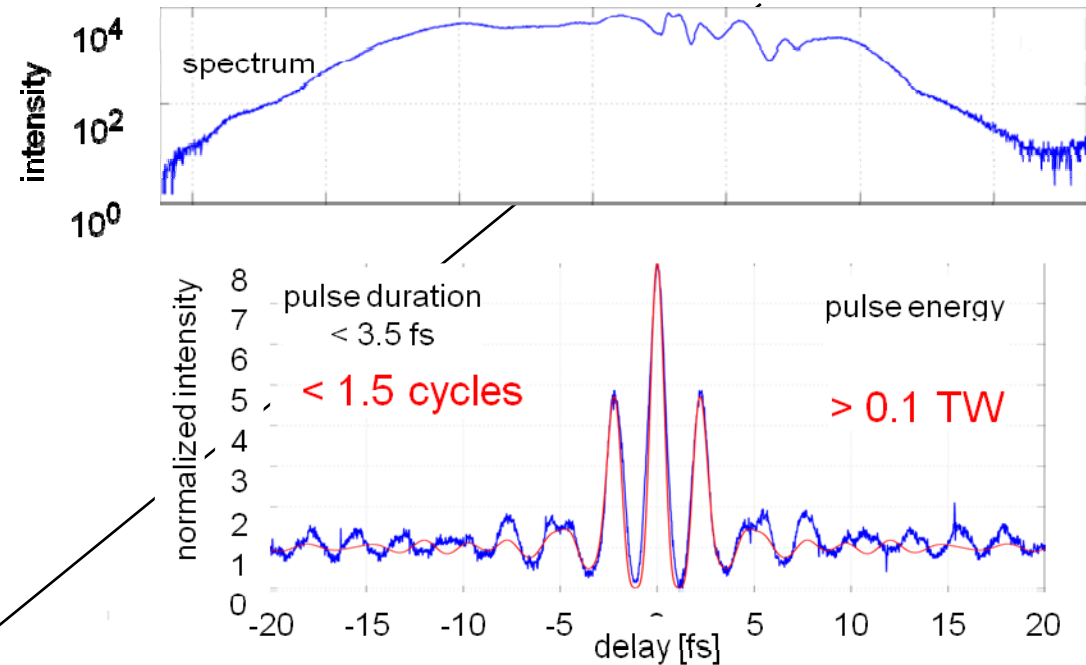
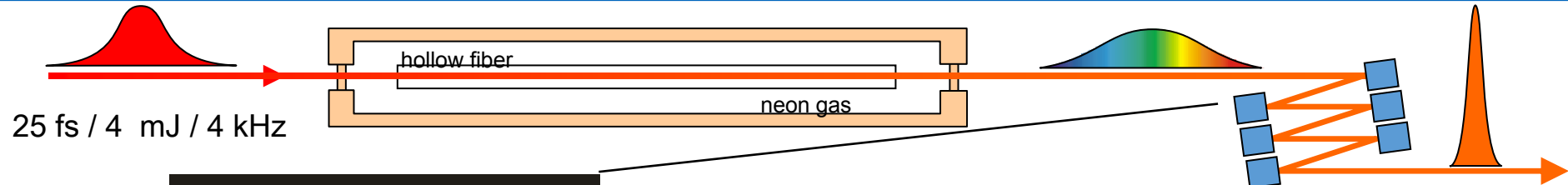


few-cycle pulses





few-cycle pulses



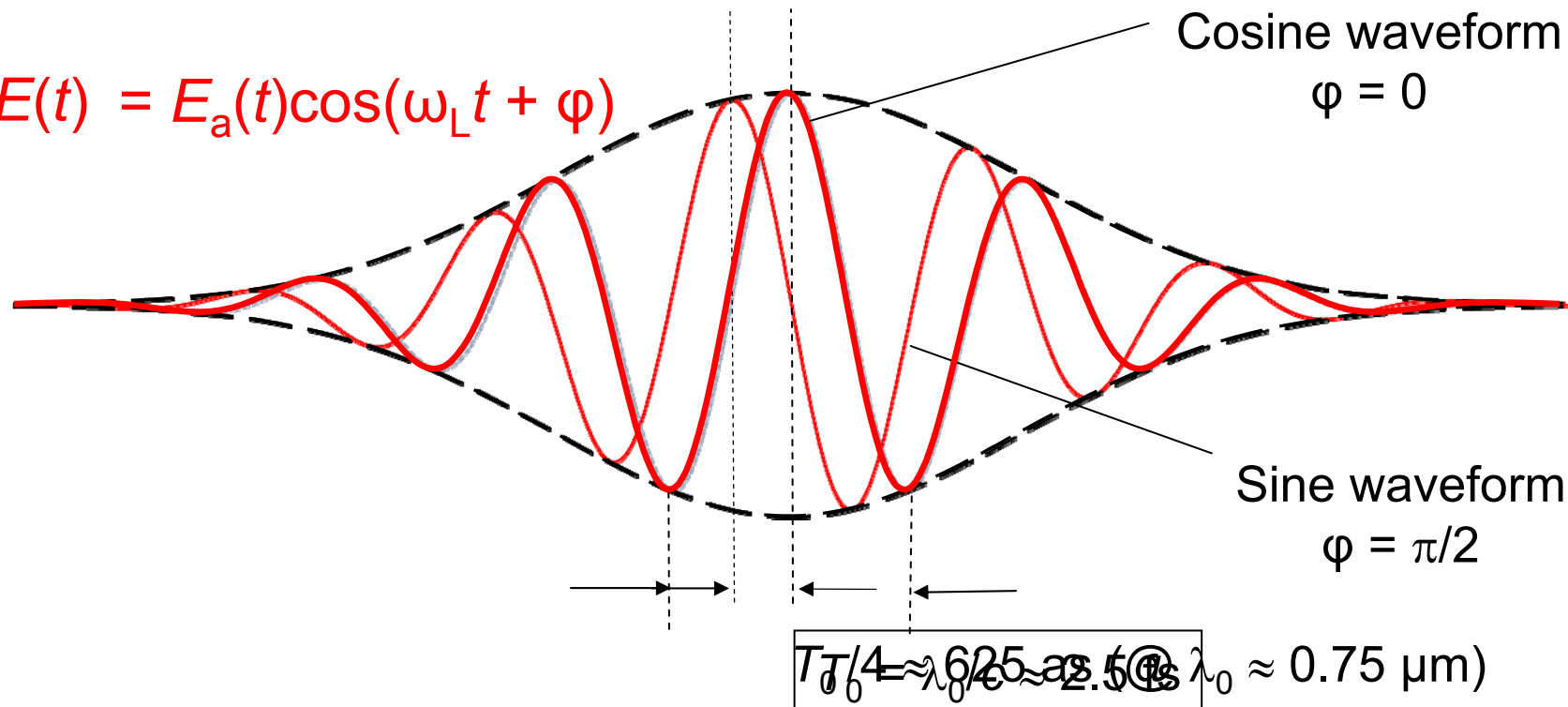




the carrier-envelope phase CEP



$$E(t) = E_a(t)\cos(\omega_L t + \varphi)$$



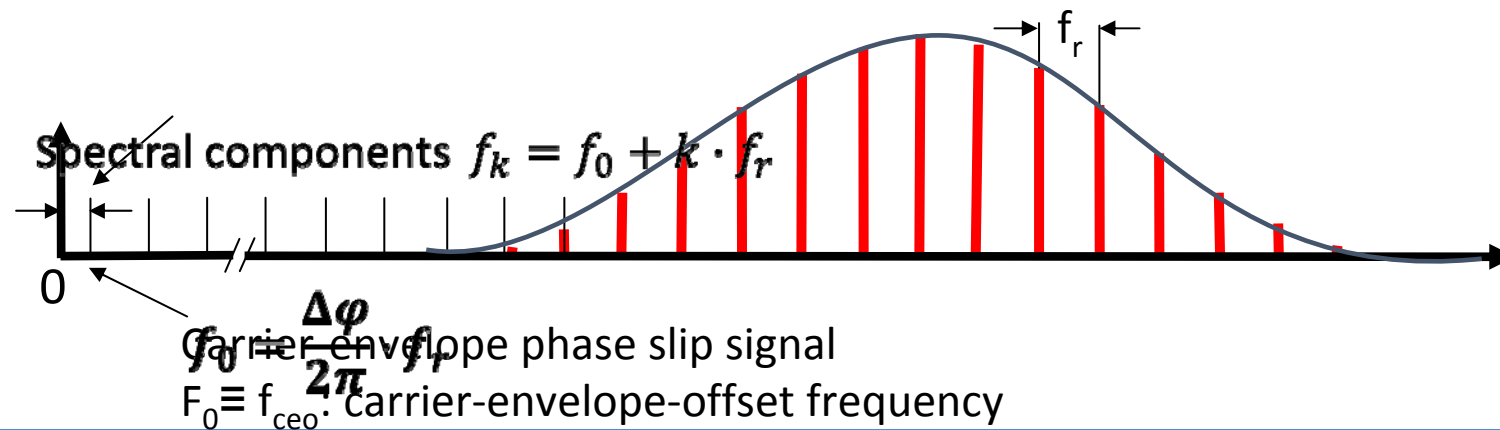
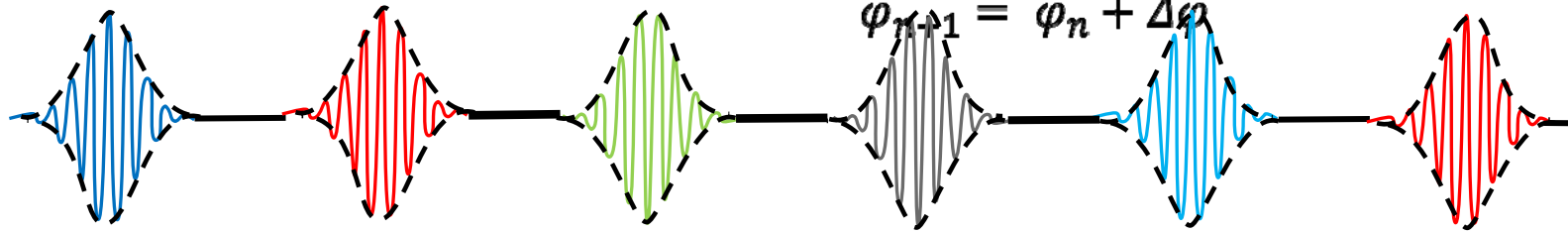


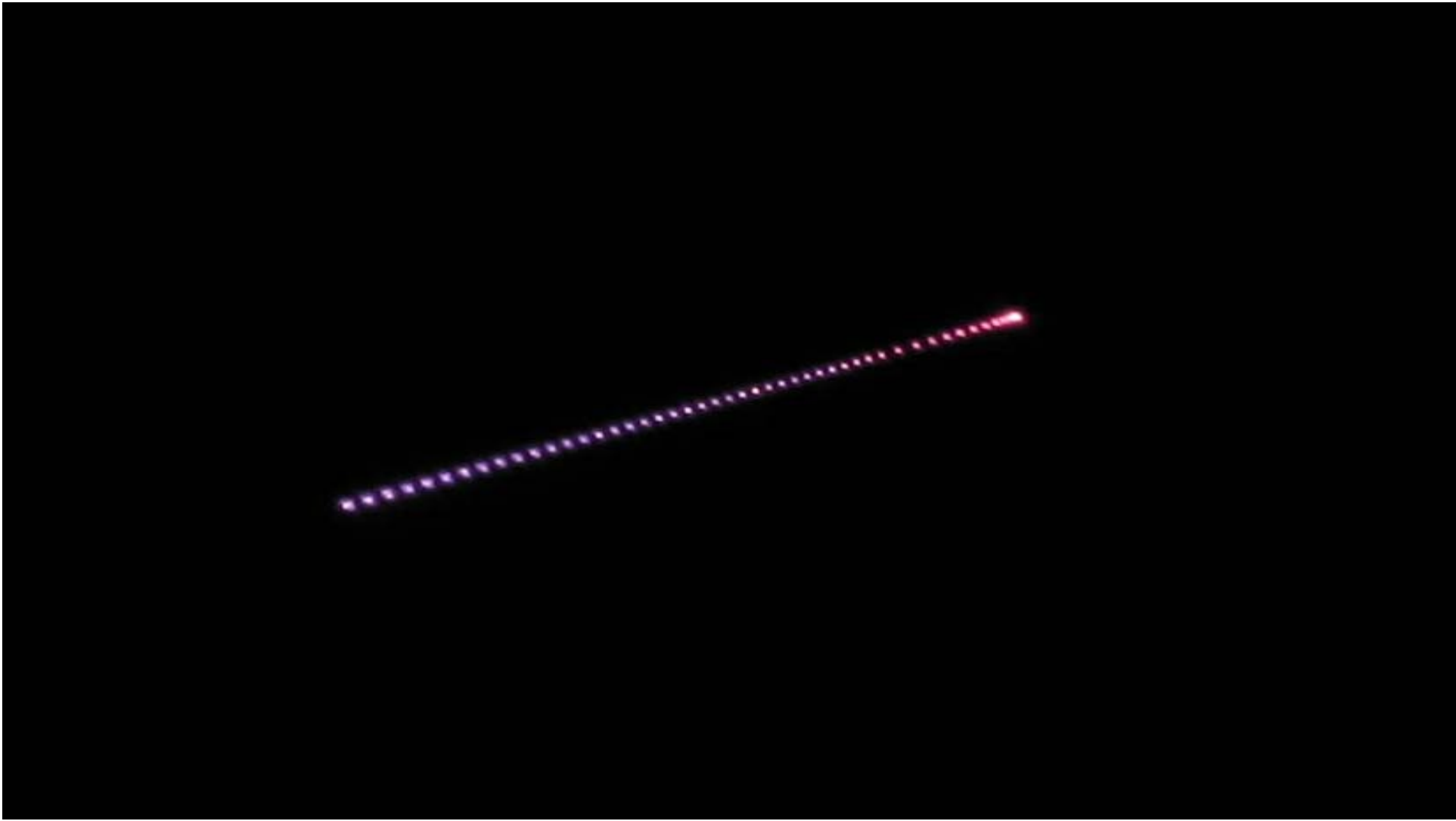
pulse train from a mode-locked laser



$$E_n(t) = A(t) \cdot e^{-i\omega_0 t + \varphi_n} + c.c.$$

$$\varphi_{n+1} = \varphi_n + \Delta\varphi$$



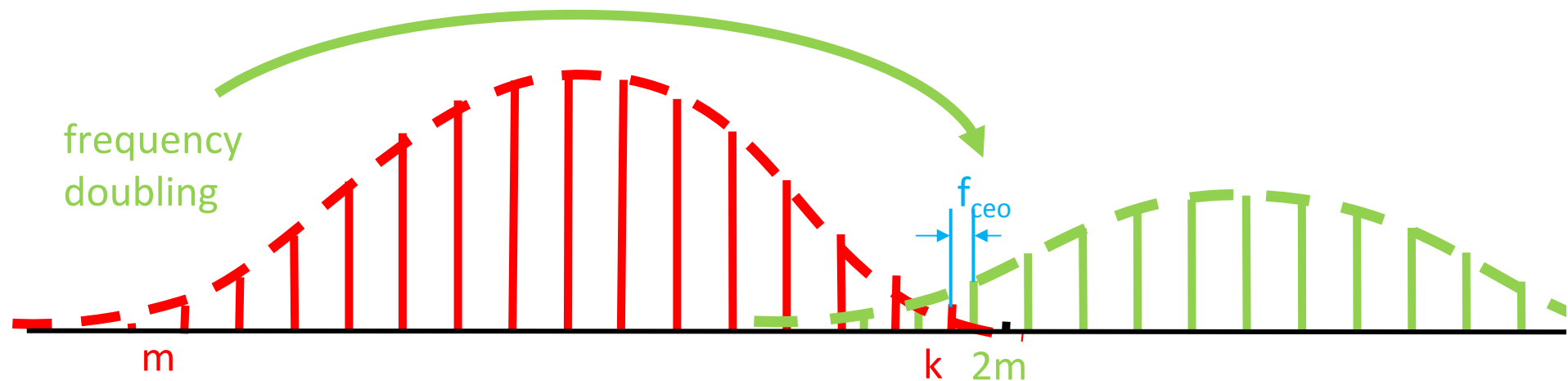




frequency-domain control of the CEP



T. W. Hänsch *et al.*, 1997, 1999; U. Keller *et al.*, 1999



Beating of the fundamental

$$f_k = f_0 + k \cdot f_r$$

and SH

$$2f_m = 2f_0 + 2m \cdot f_r$$

for $k = 2m$

$$2f_m - f_k = 2f_0 - f_0 + (2m - k) \cdot f_r = f_0 \equiv f_{ceo}$$

Beat signal yields temporal evolution of φ_n

First implementation: D. Jones *et al.*, Science 288, 635 (2000); A. Apolonski *et al.*, PRL 85, 740 (2000)

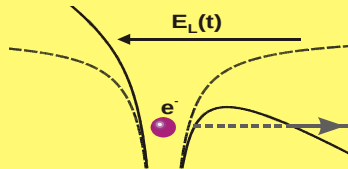


high-order harmonic generation in the gas phase



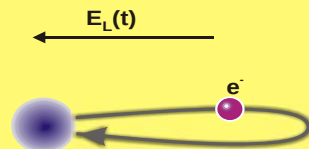
Step 1

Optical field ionization



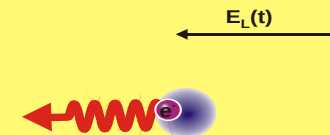
Step 2

e^- Acceleration



Step 3

XUV emission on recollision



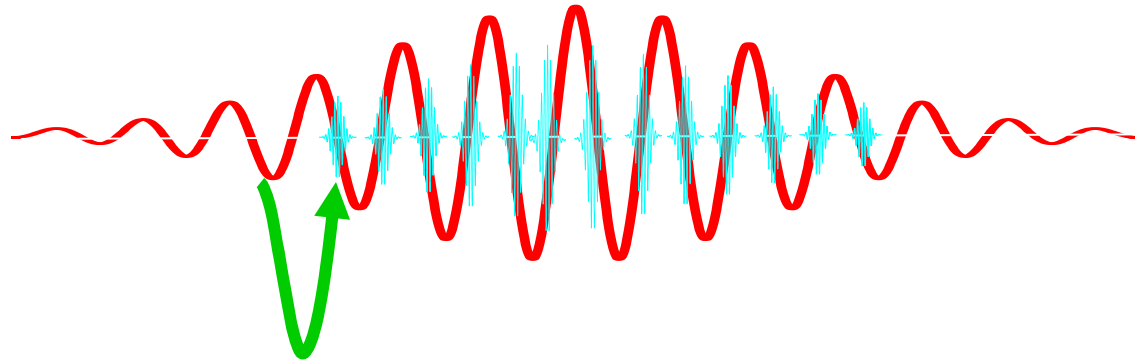
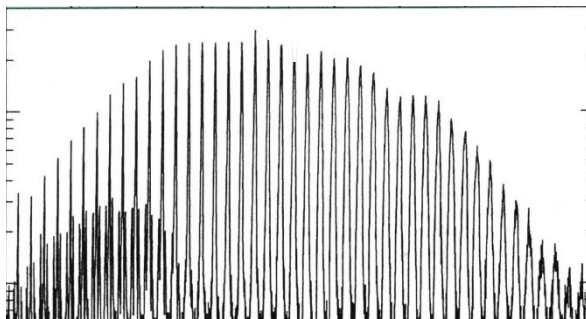
P.B. Corkum, PRL 71 (1993)



recombination emission from strongly-driven atoms



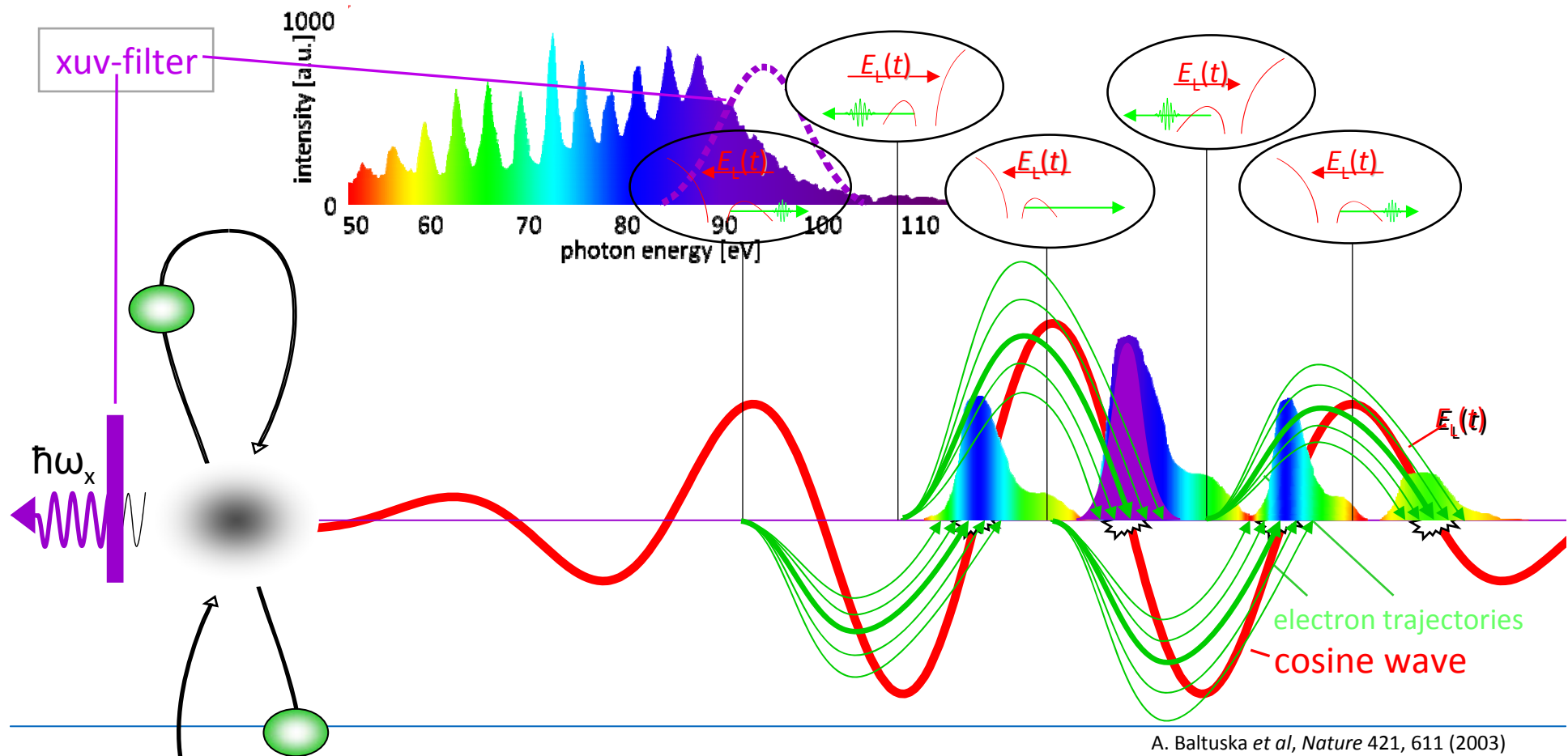
Multi-cycle driver pulse : $\tau_p \gg T_o$
High-order odd harmonics of the driver laser



Cut-off harmonics: train of attosecond bursts

L'Huillier, Balcou, 1993, *PRL* 70, 774
Macklin *et al*, 1993, *PRL* 70, 766

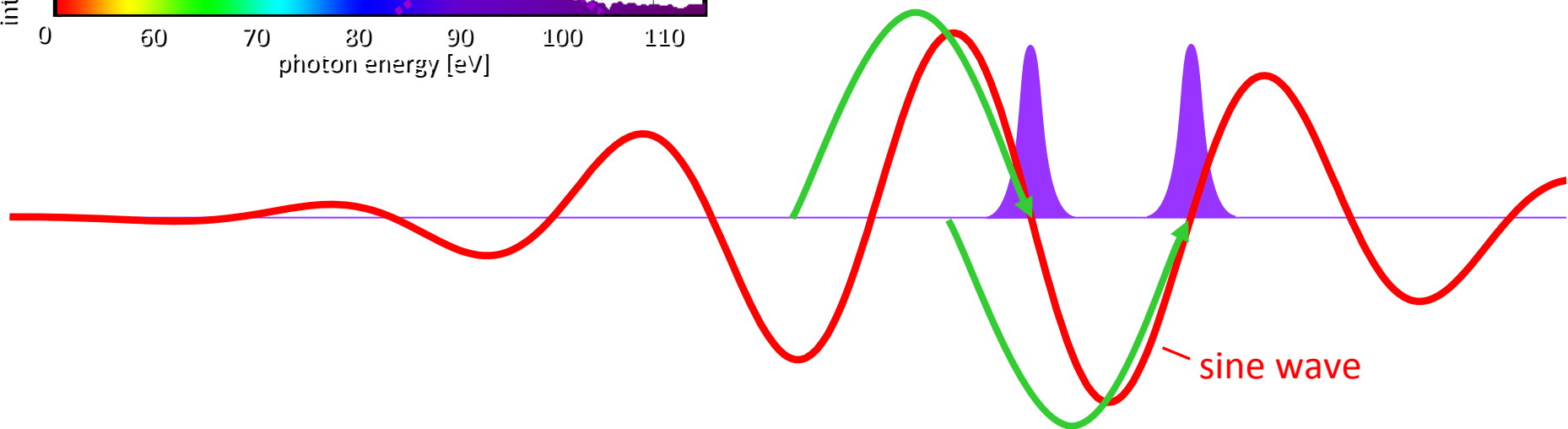
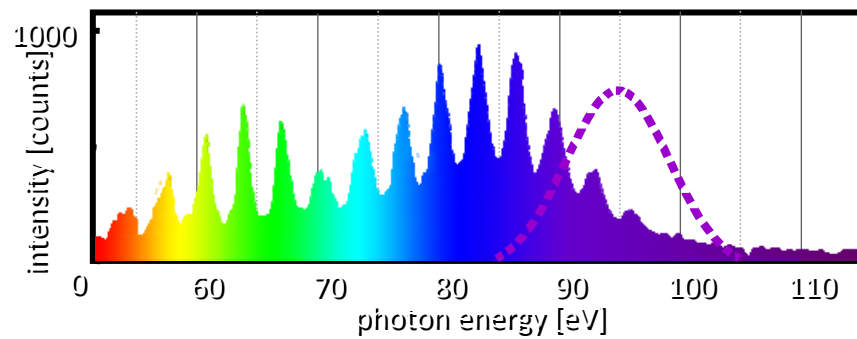
Paul *et al*, *Science* 292, 1689 (2001)
Tsakiris, Charalambidis *et al*, 2003



A. Baltuska *et al*, *Nature* 421, 611 (2003)

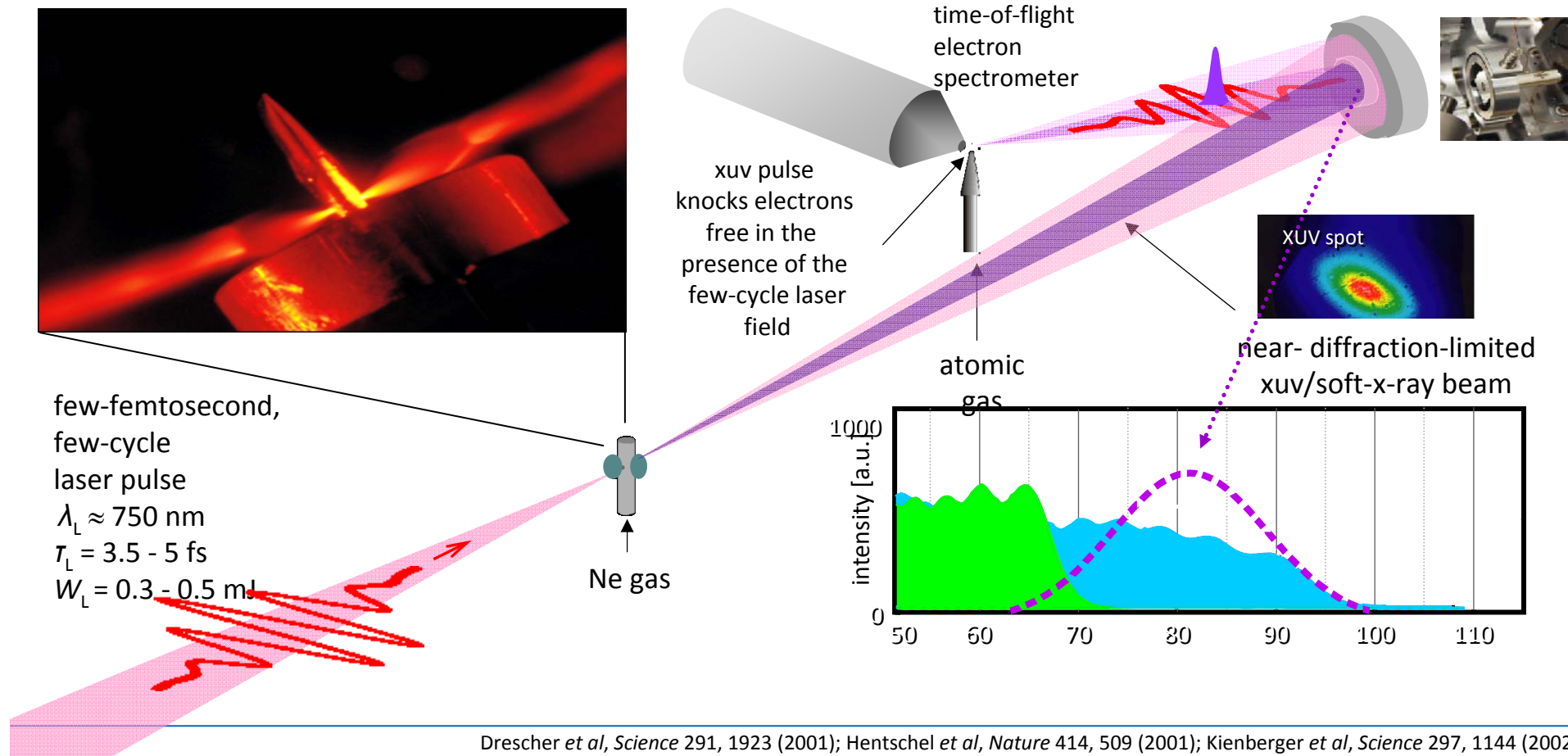


attosecond xuv/soft x-ray pulse generation depending on the CEP





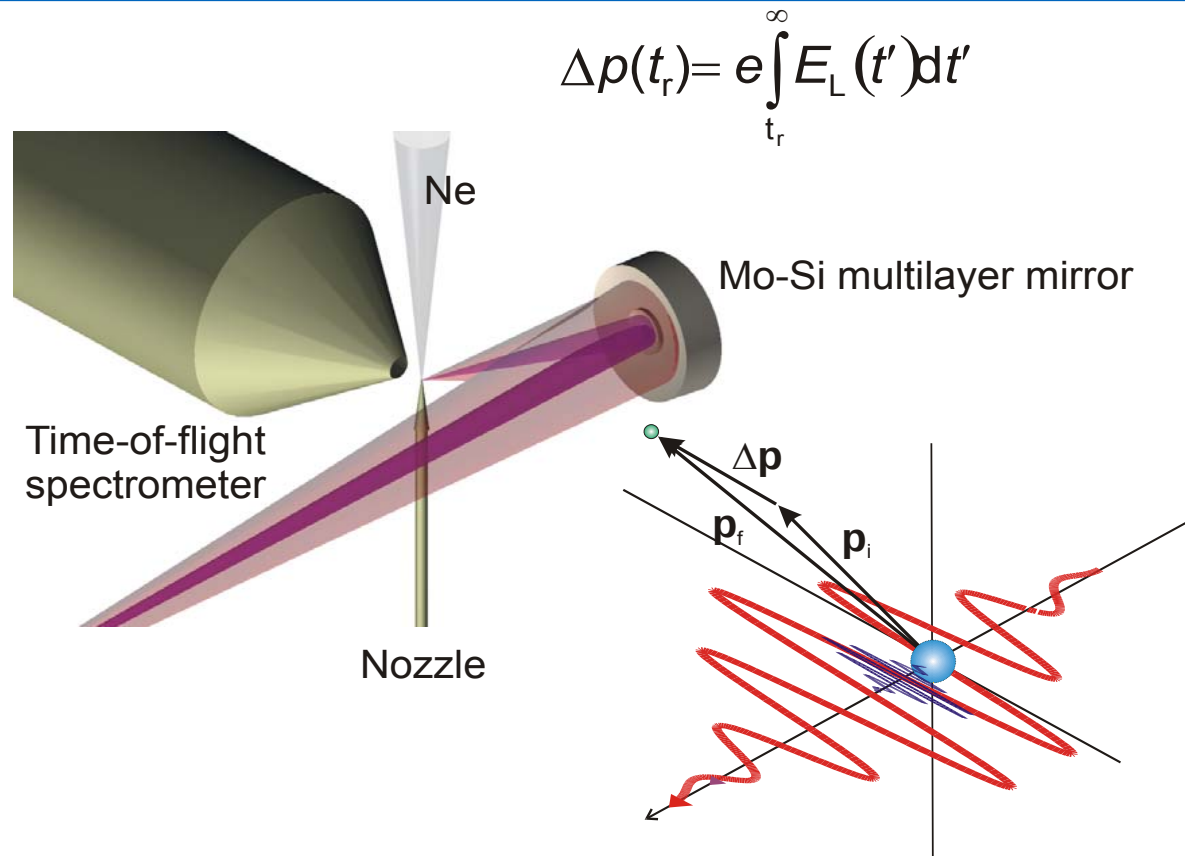
attosecond pulse generation and measurement



Drescher *et al*, *Science* 291, 1923 (2001); Hentschel *et al*, *Nature* 414, 509 (2001); Kienberger *et al*, *Science* 297, 1144 (2002)

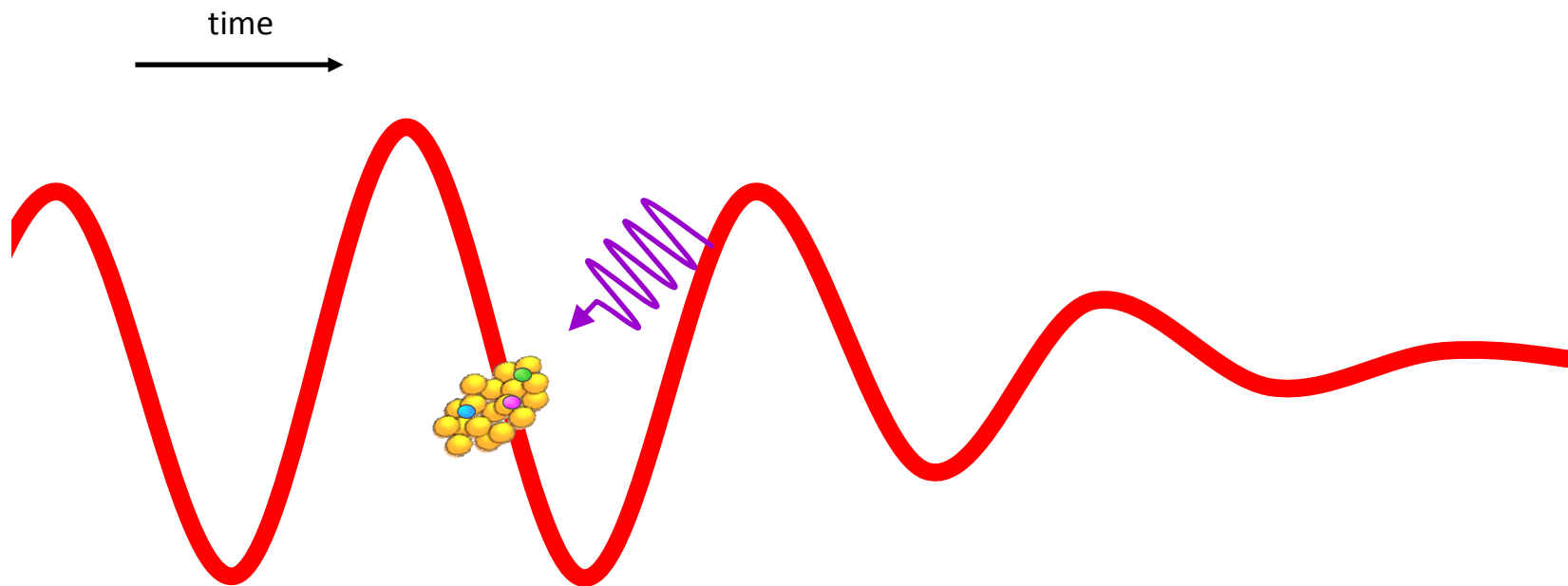


photoelectrons generated by an as-pulse



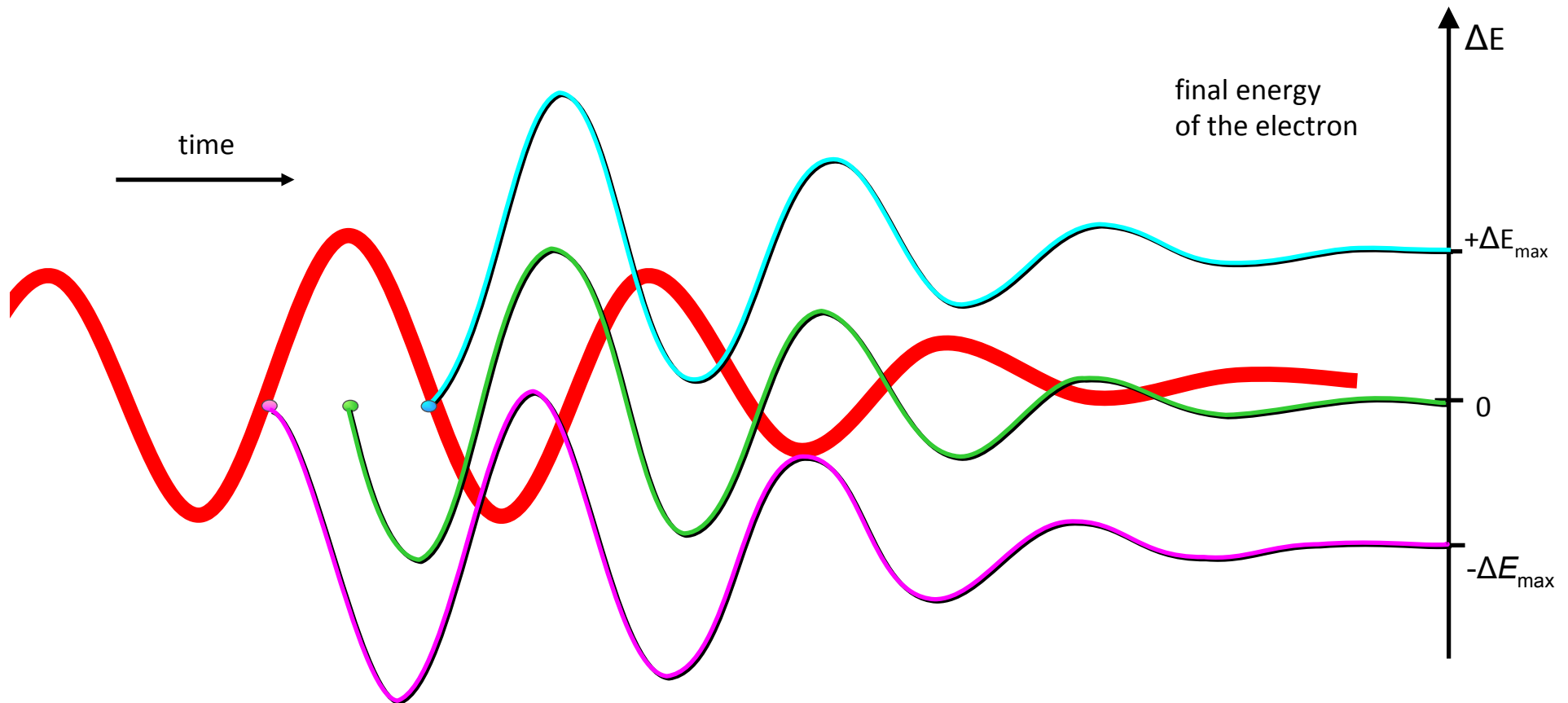


ionization at different instants of time



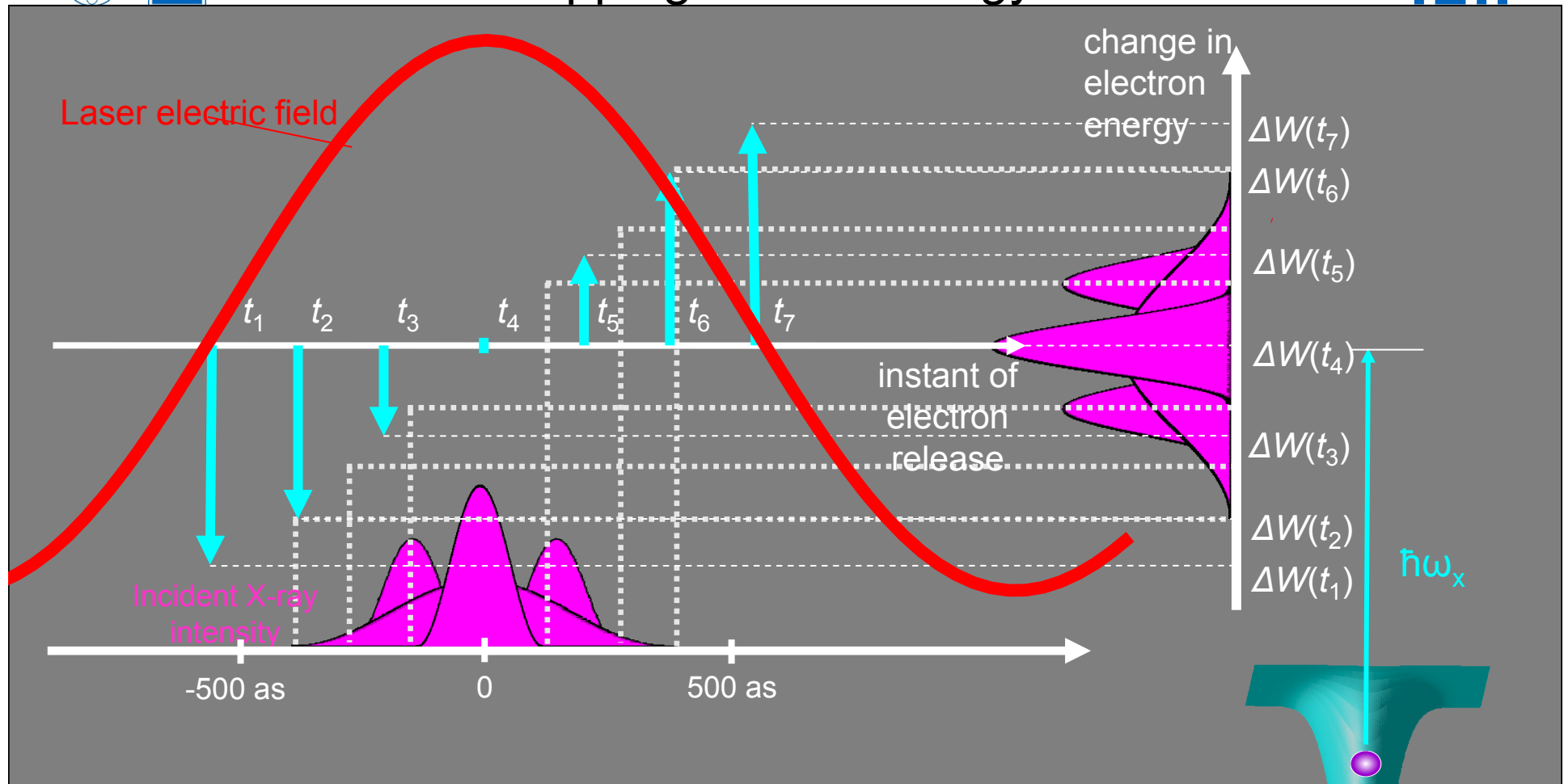


final energy of photoelectrons depends on the release time



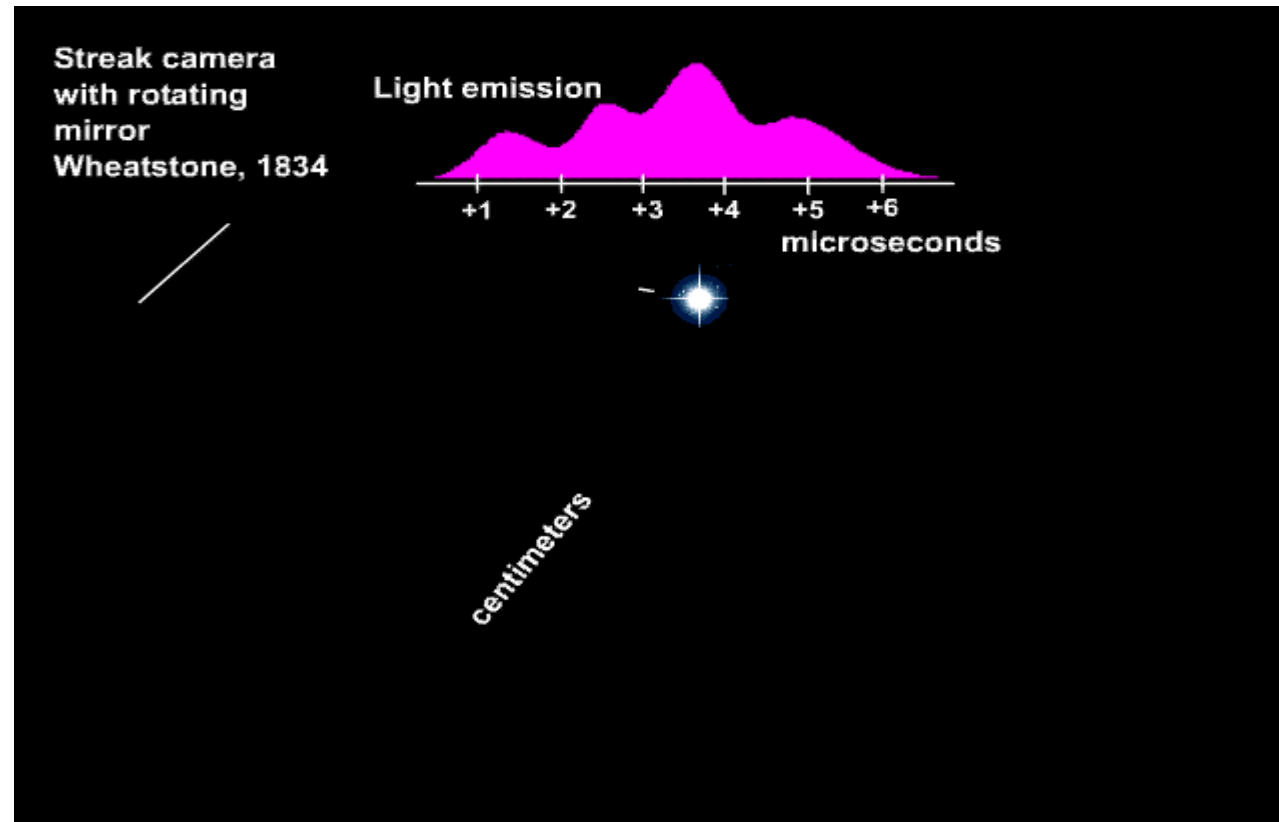


mapping time to energy



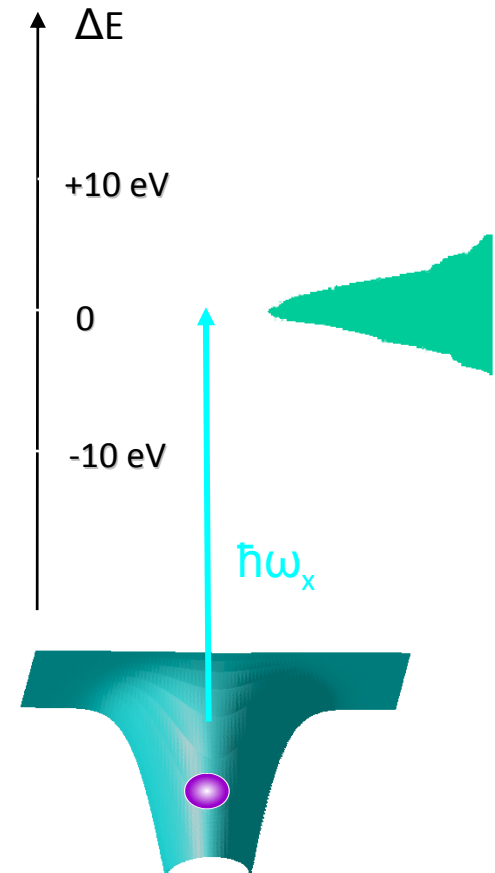
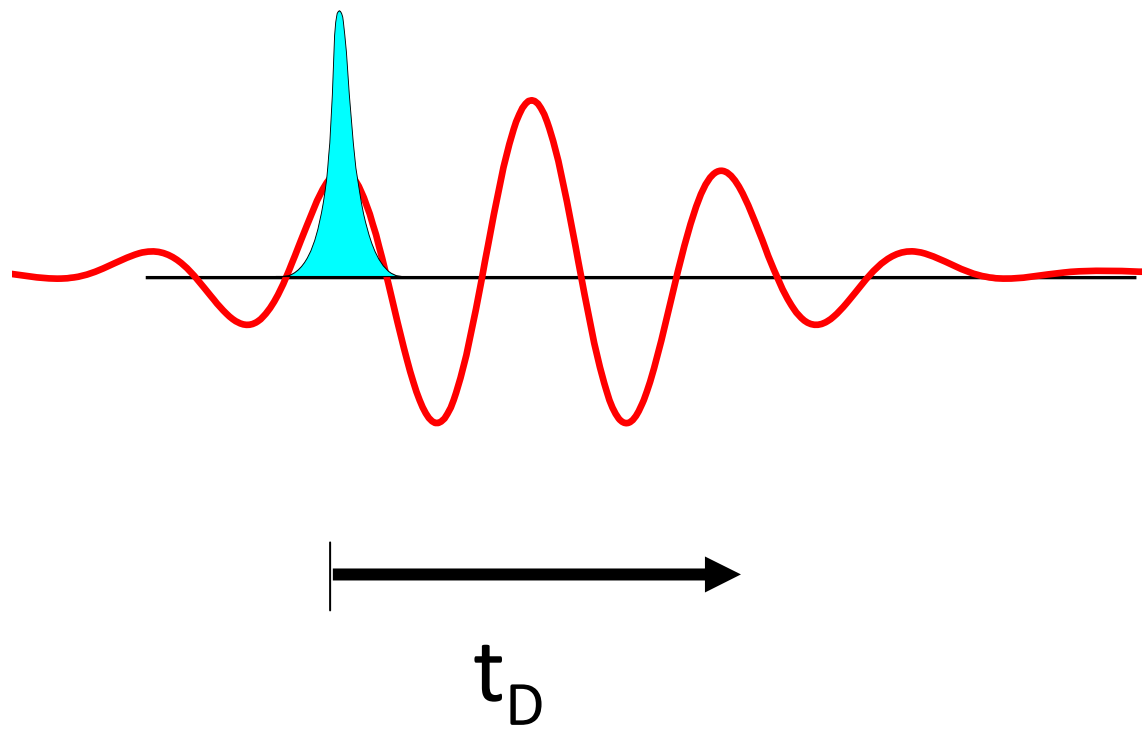


optical streak camera, 1834





sampling field oscillations



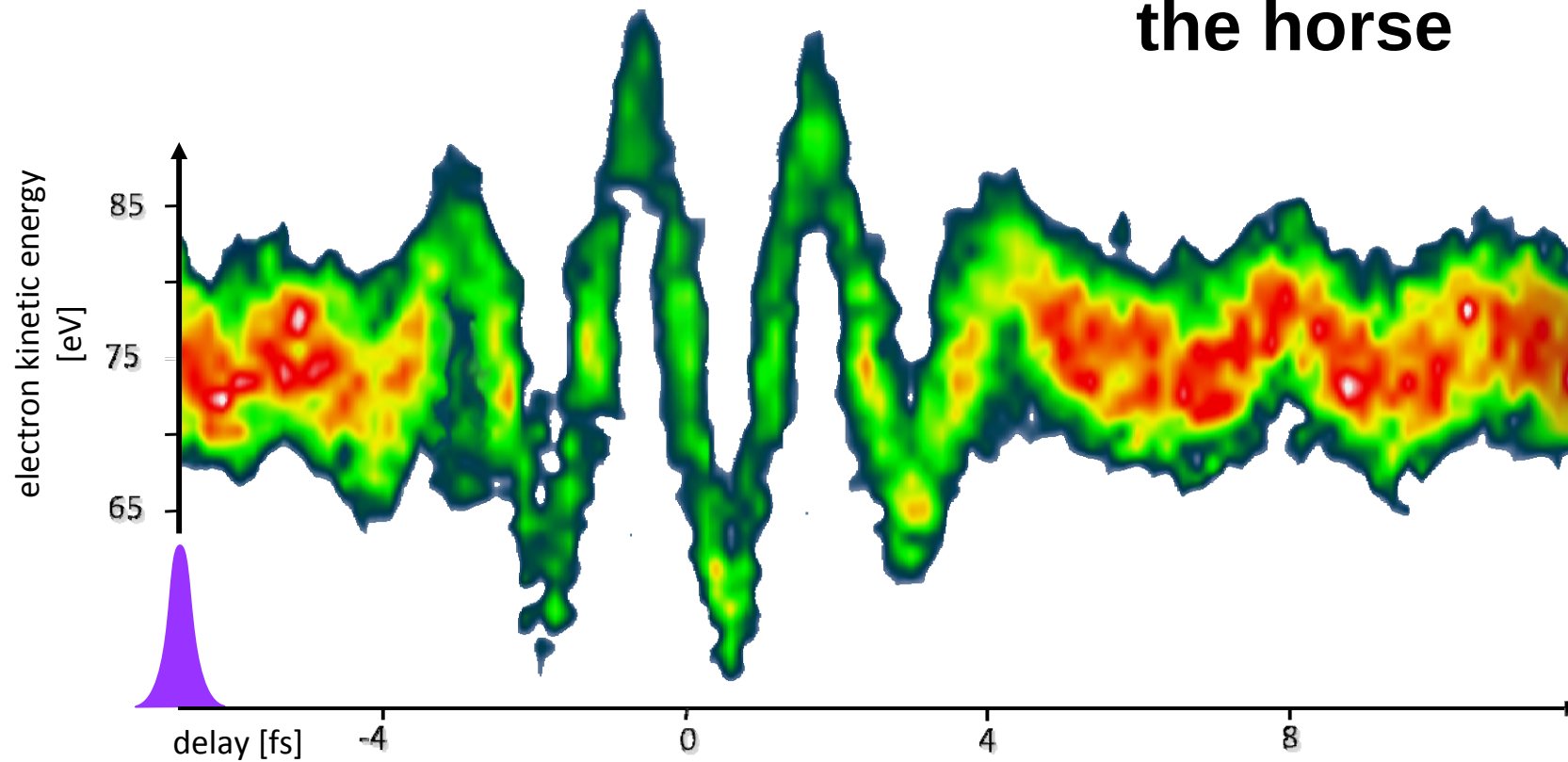
R. Kienberger et al, Nature 294 (2004)



the spectrogram: “a streaking trace”



the horse

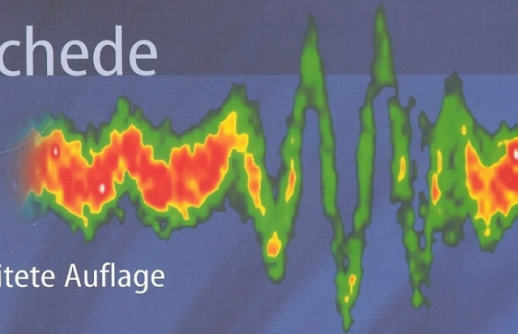




„DIE GANZE PHYSIK ZUM 21. JAHRHUNDERT“

Gerthsen Physik

D. Meschede



23., überarbeitete Auflage



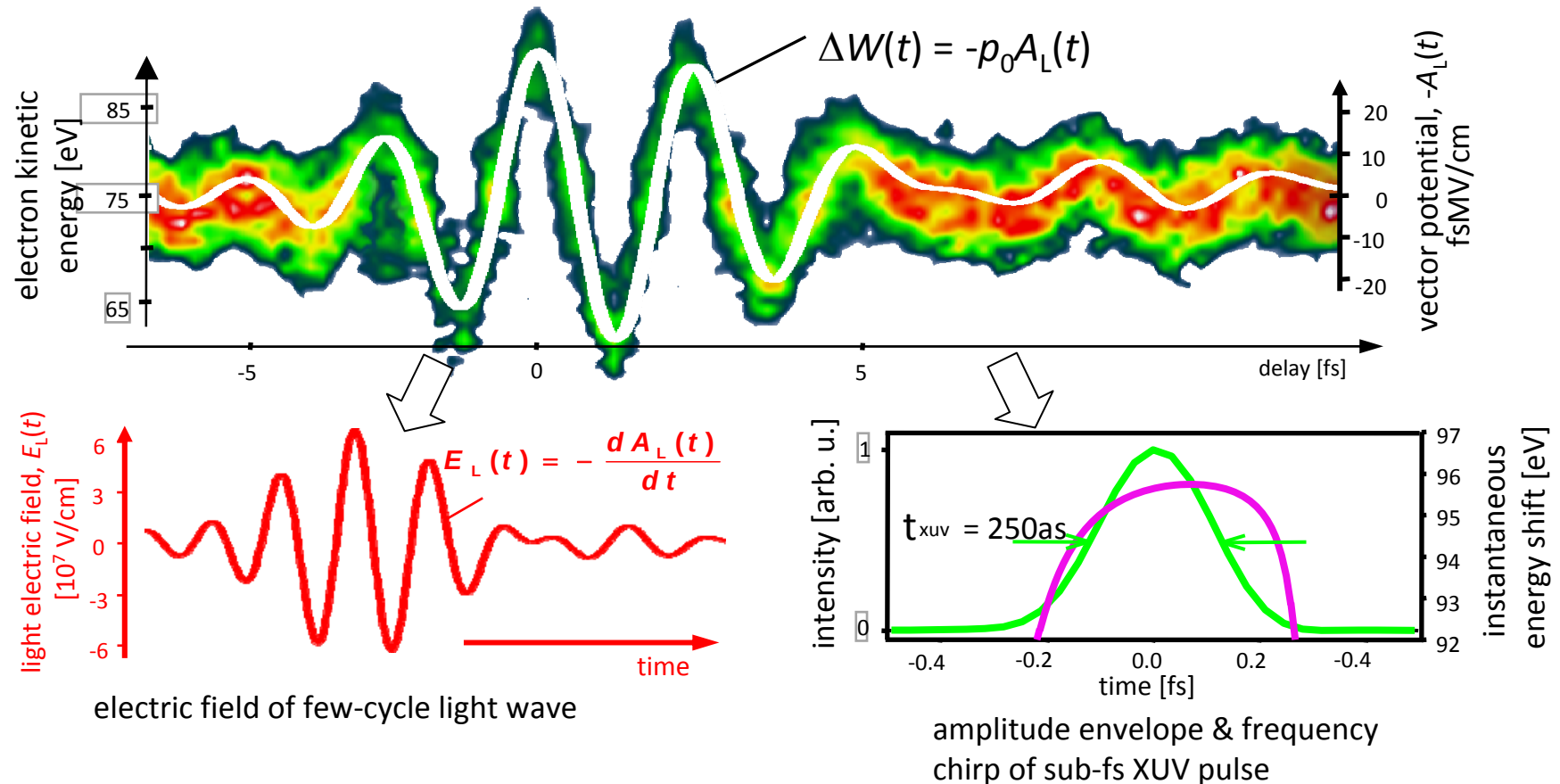
CD-ROM

1074 Aufgaben und
vollständige Lösungen auf CD-ROM.
30 Animationen zur
Relativitätstheorie auf CD-ROM.

 Springer

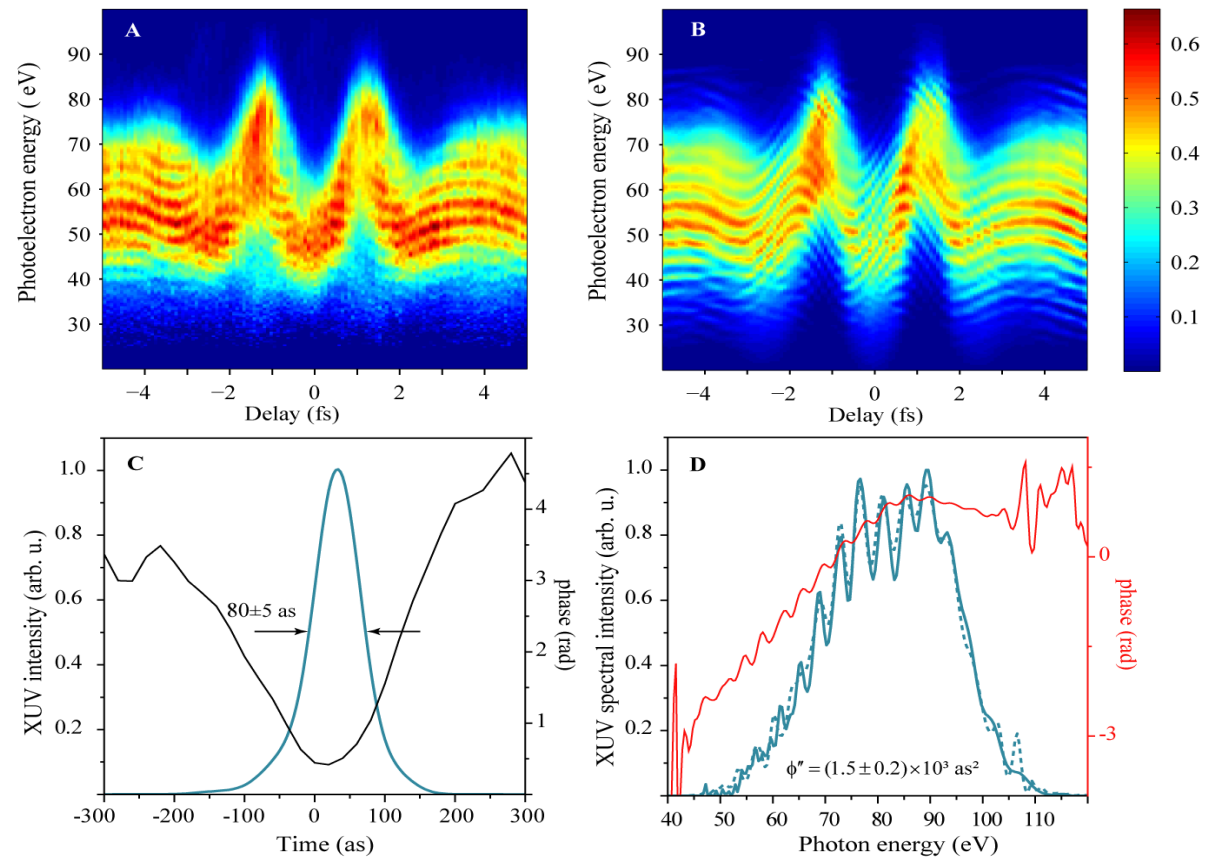


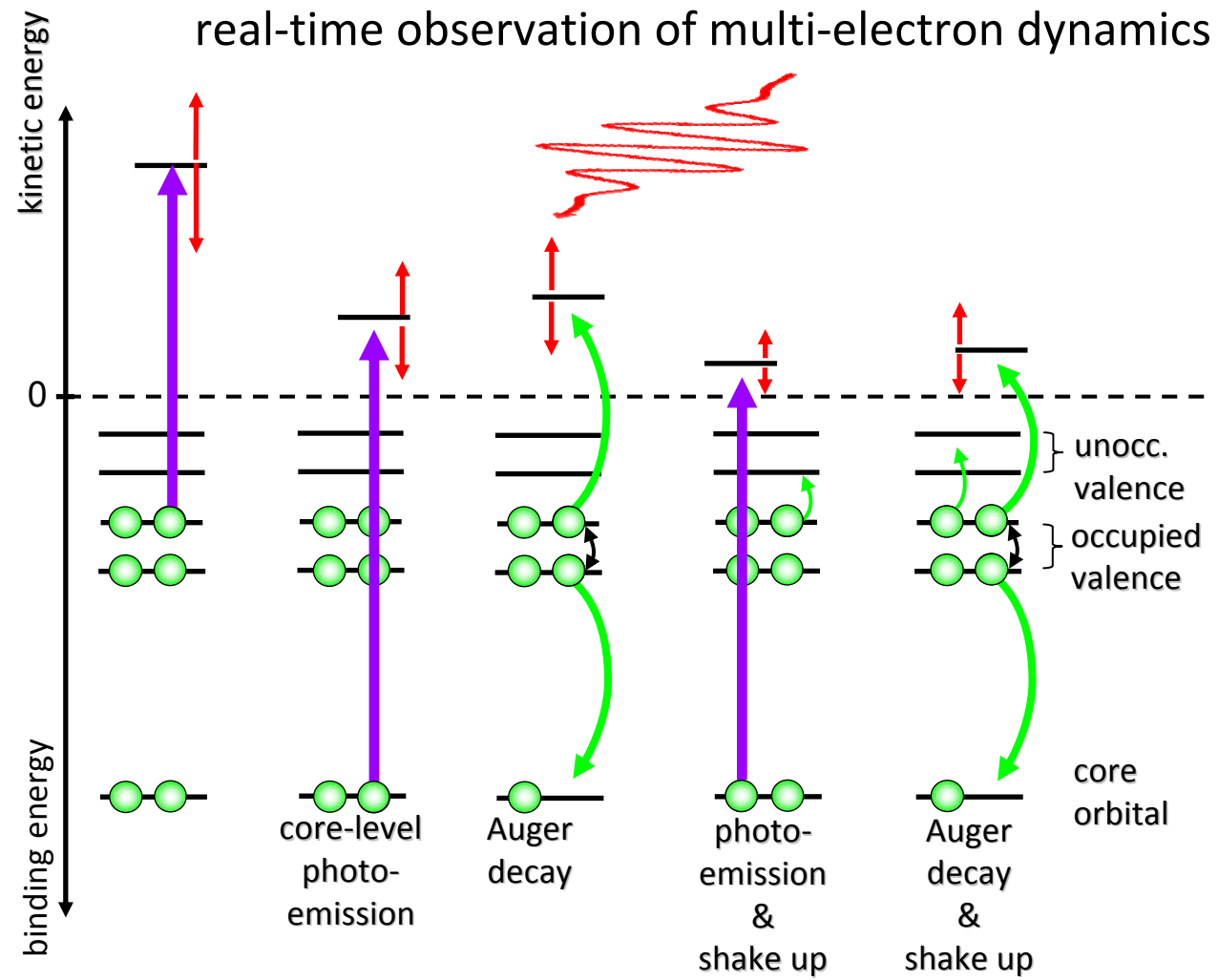
attosecond streak camera: complete measurement of a few-cycle light wave *and* a sub-fs xuv pulse





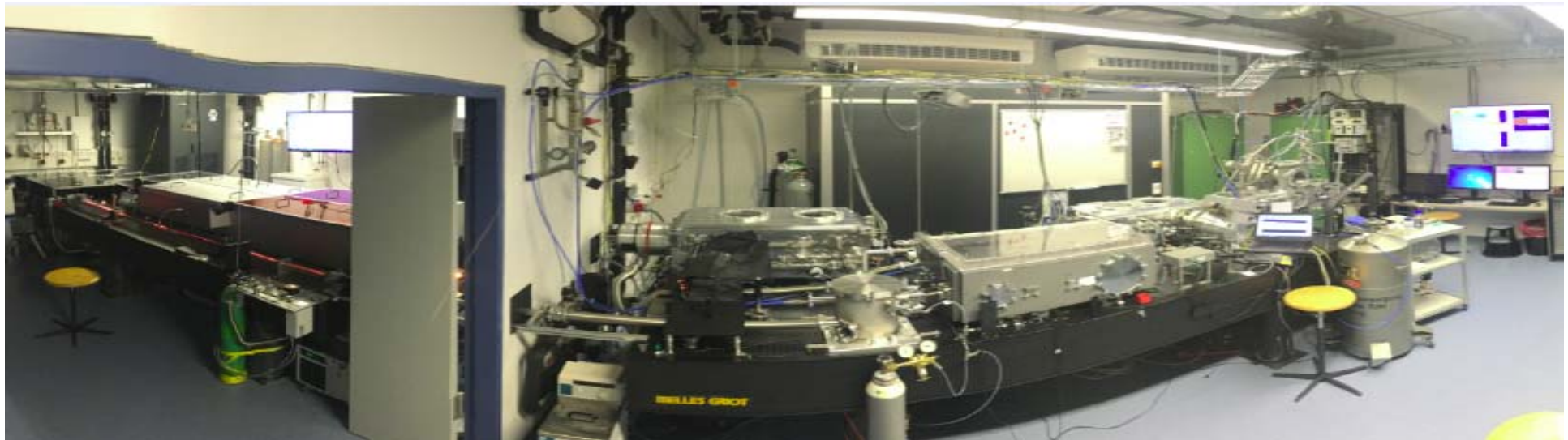
FROG-like spectrograms of sub-100-as pulses





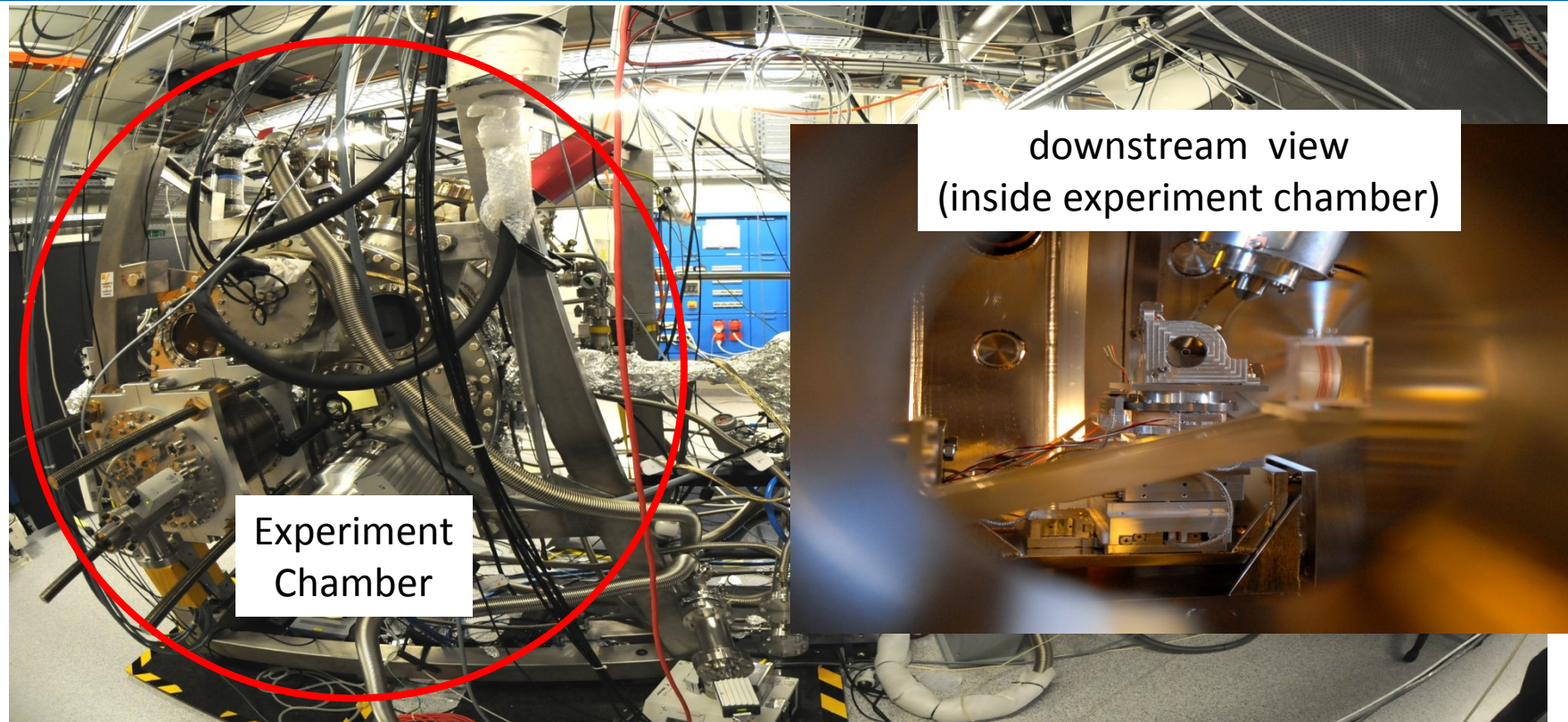


AS-beamline at TUM





AS3 beamline

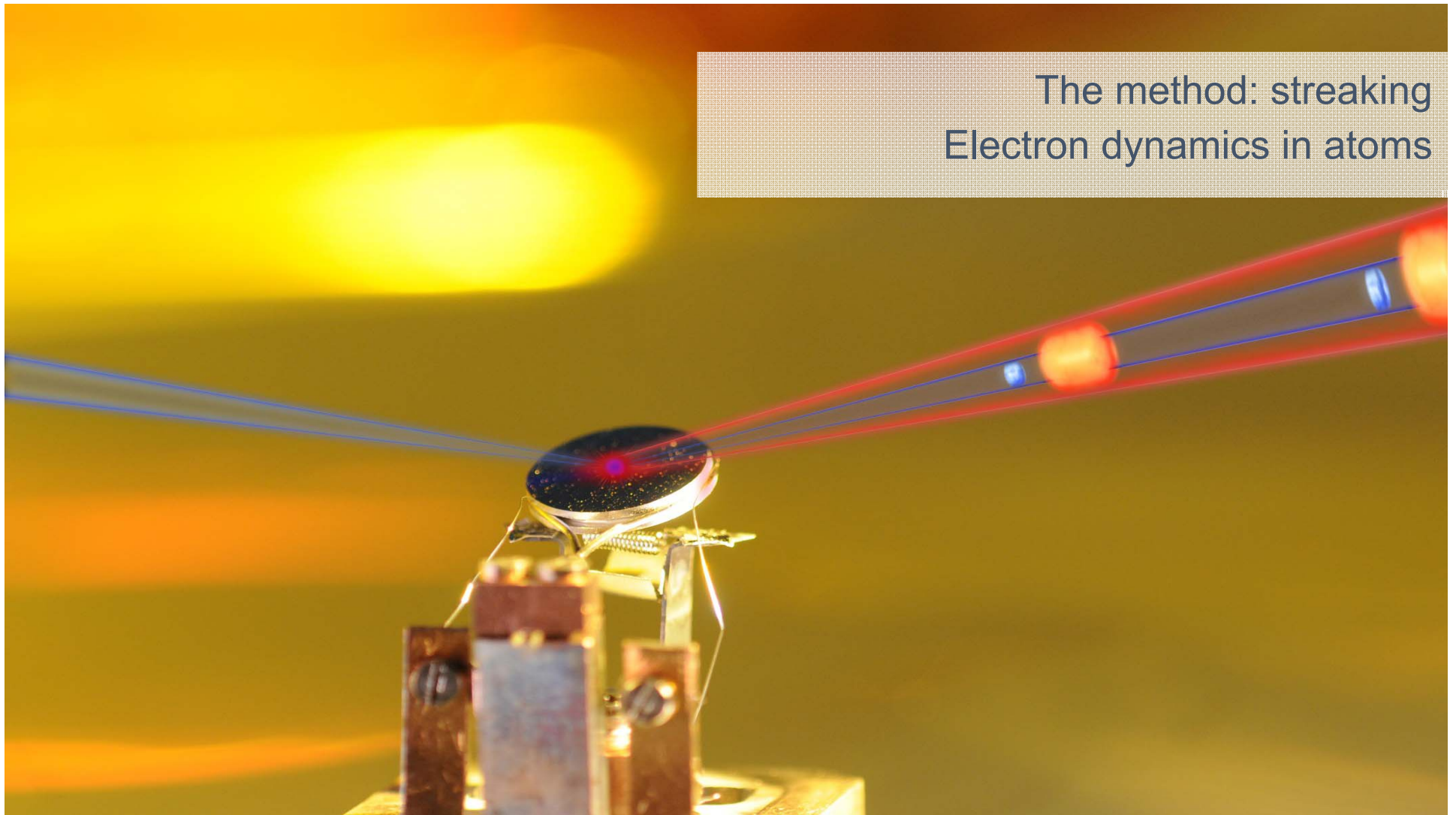


3×10^{-11} mbar

5×10^{-9} mbar

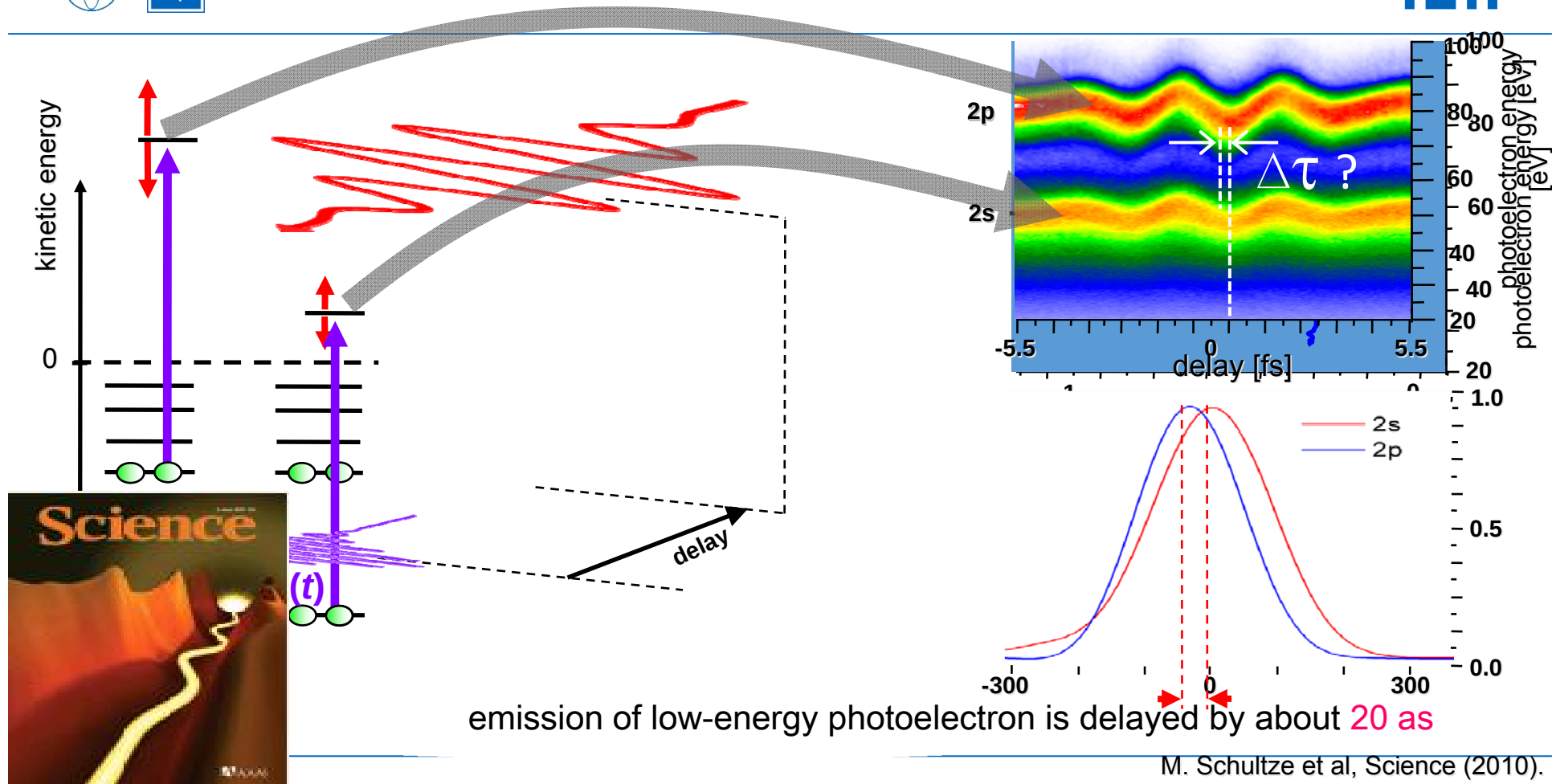
5×10^{-3} mbar

The method: streaking
Electron dynamics in atoms

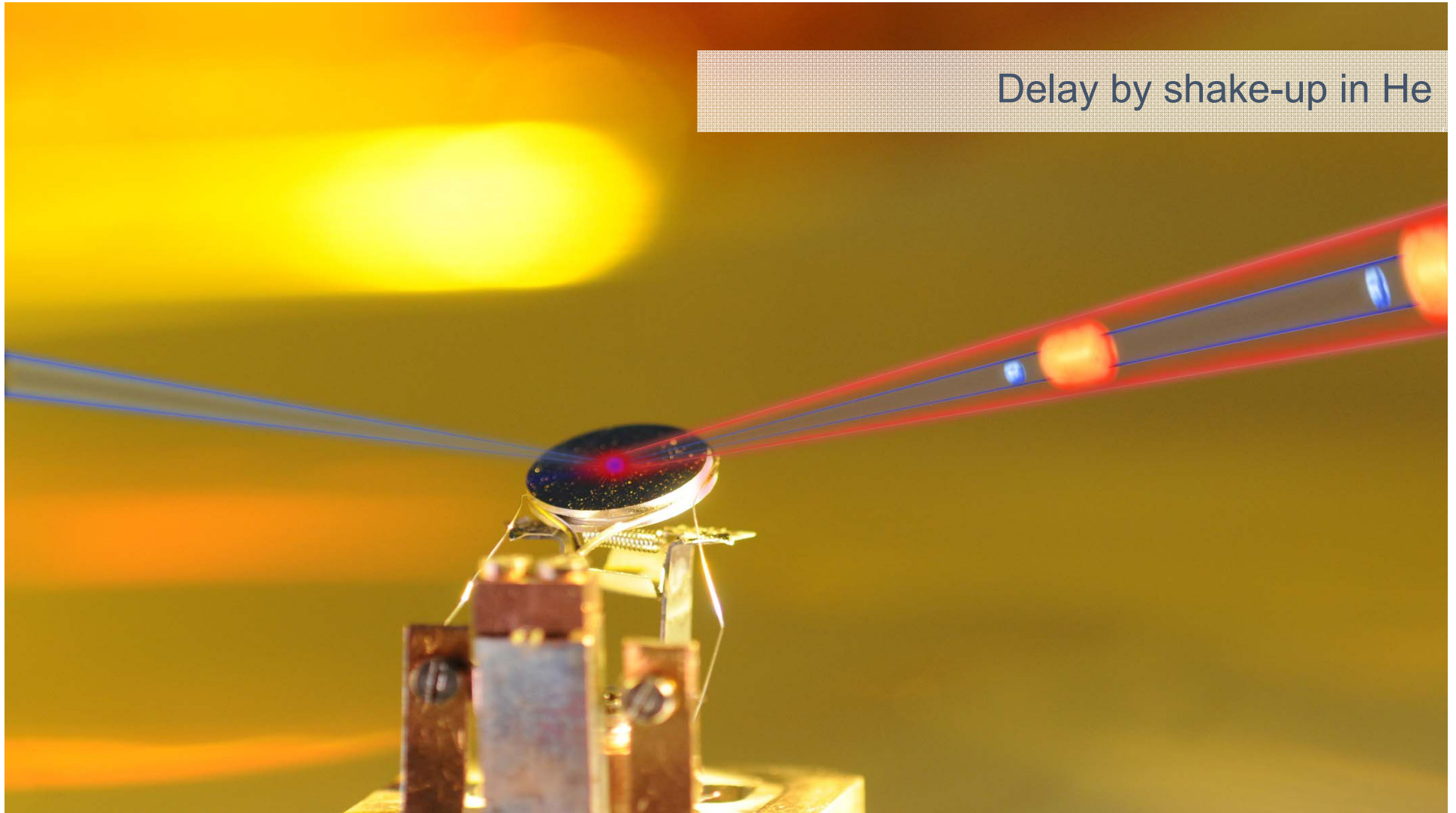




What about single photon ionization? Ne 2s / 2p

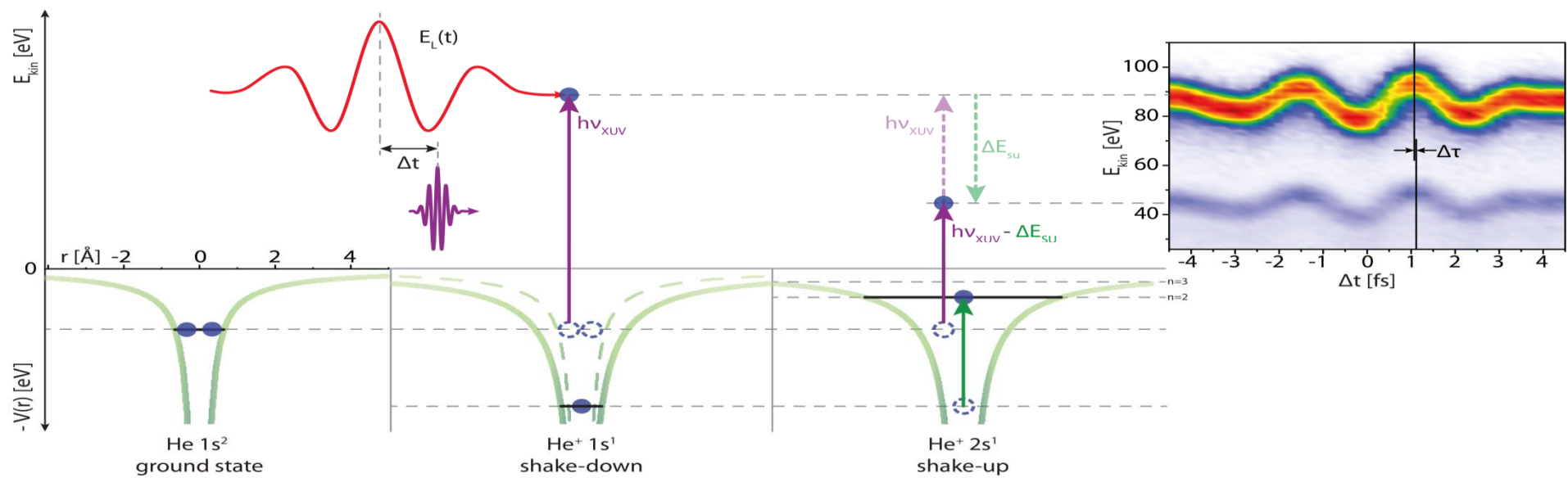


Delay by shake-up in He





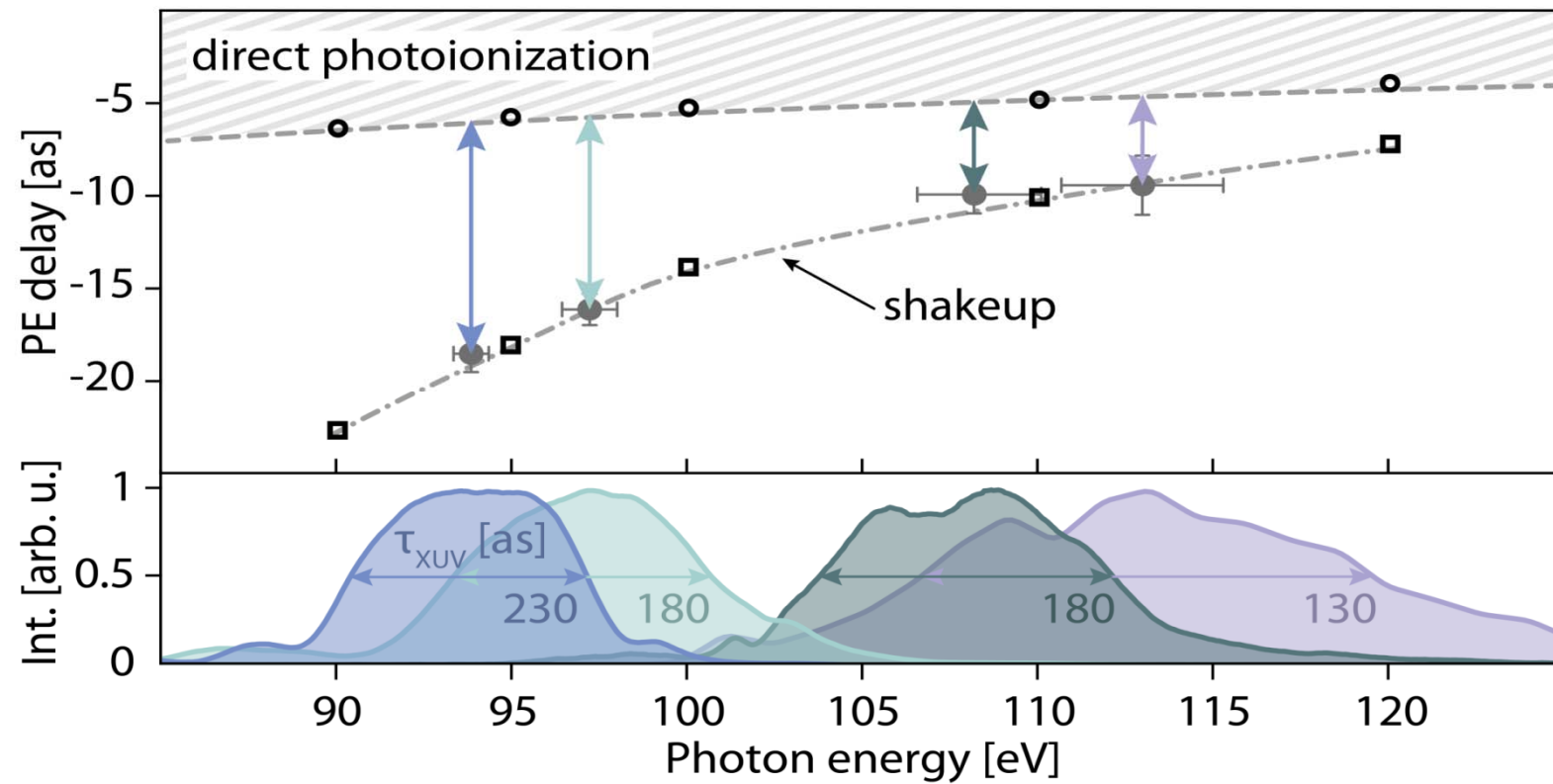
photoemission in Helium direct PE vs. shake-up



Ossiander et al., Science 2016

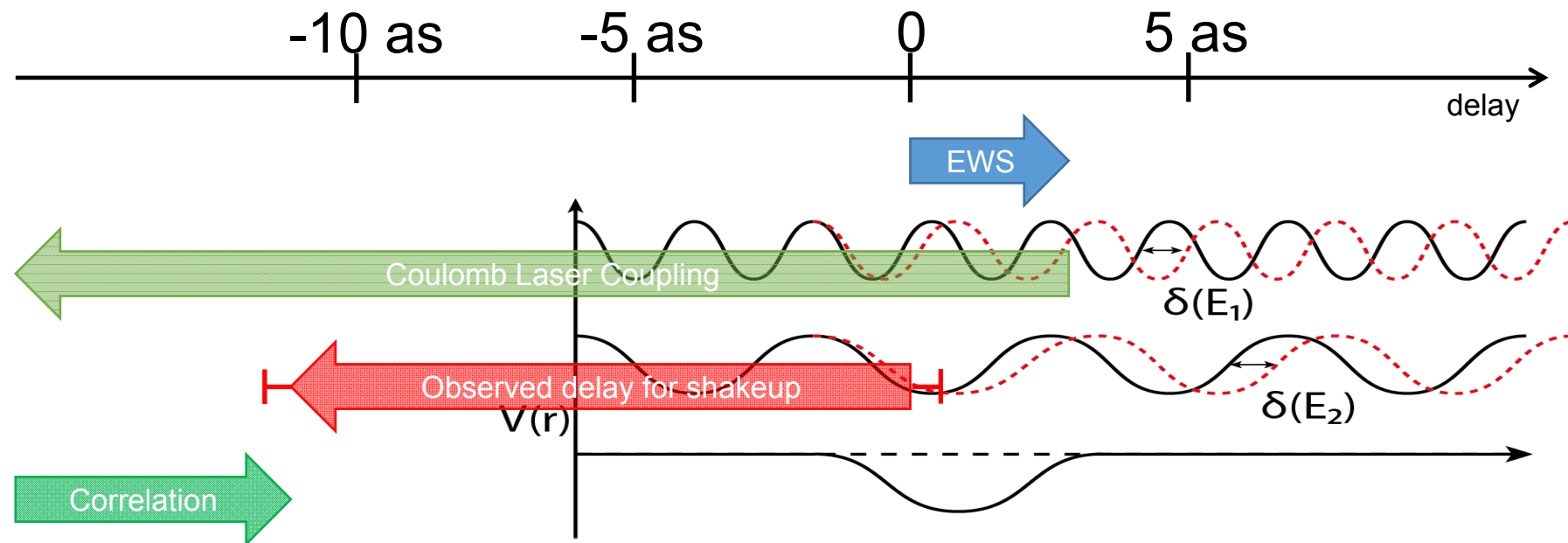


timing of photoemission in Helium





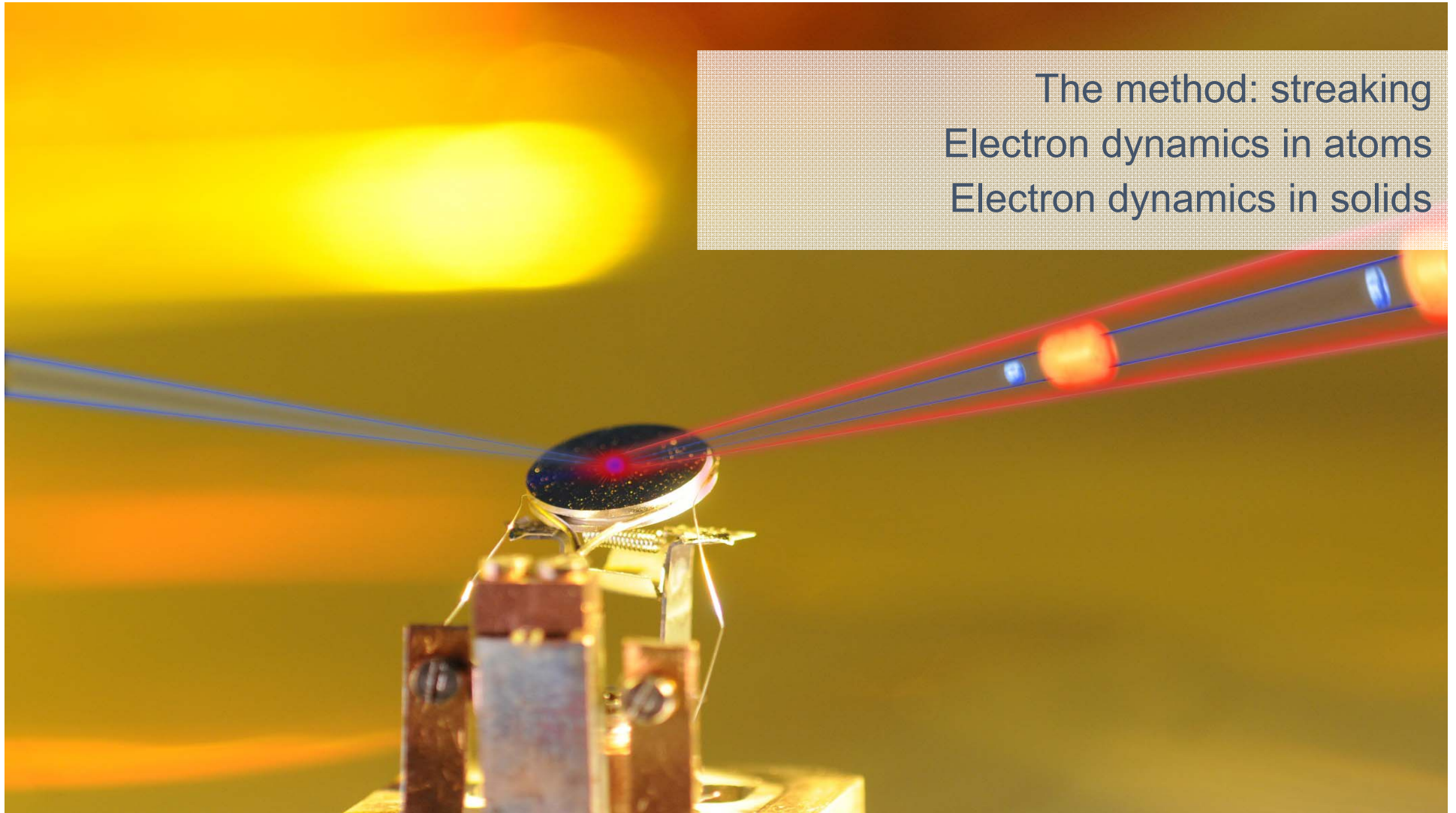
delay contributions



depending on energy: between 5as and 12 as

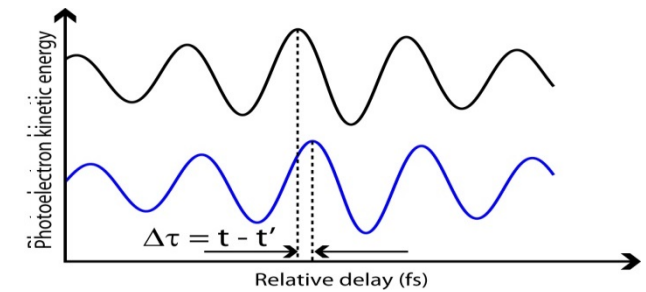
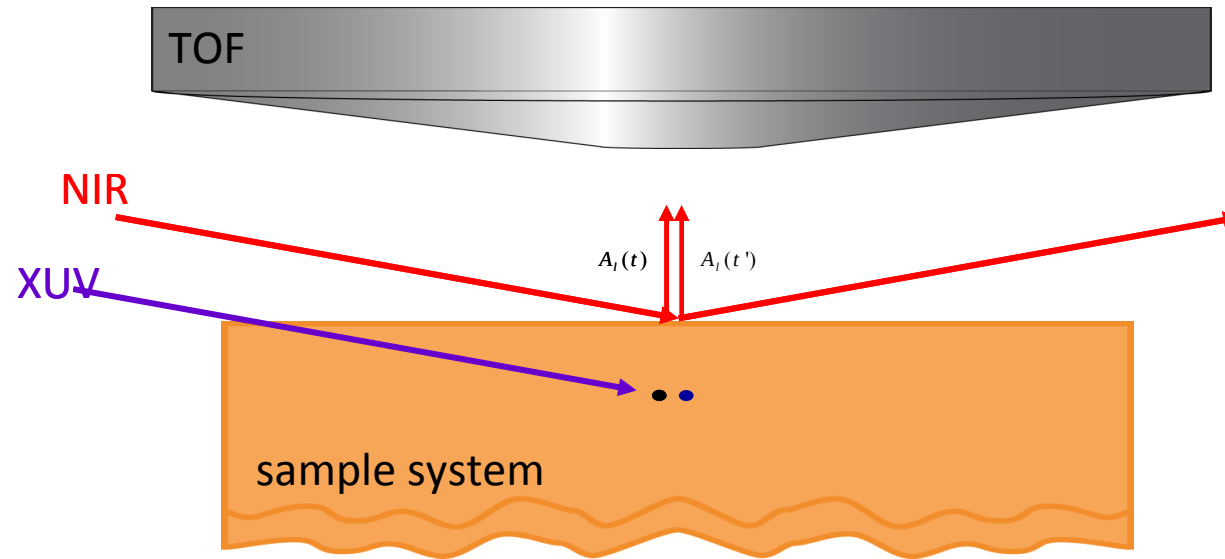
Pazourek, R., Time-resolved photoemission in one- and two-electron atoms. (TU Vienna, 2013). at <<http://katalog.ub.tuwien.ac.at/AC10774478>>

The method: streaking
Electron dynamics in atoms
Electron dynamics in solids



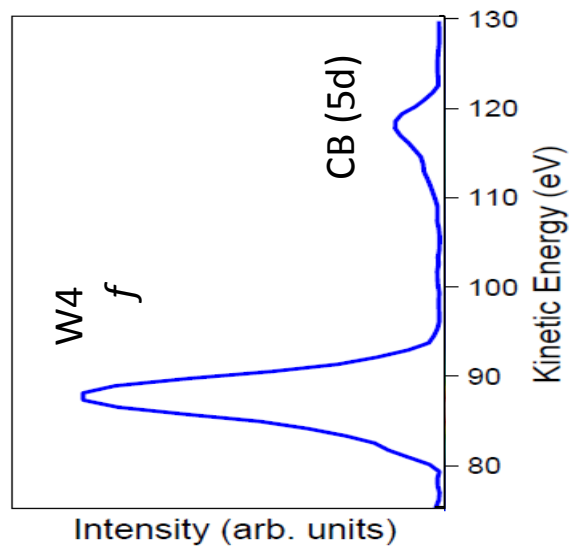
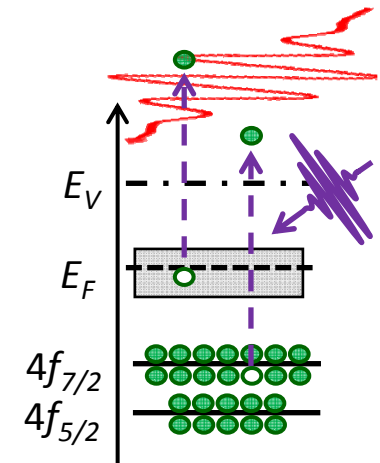
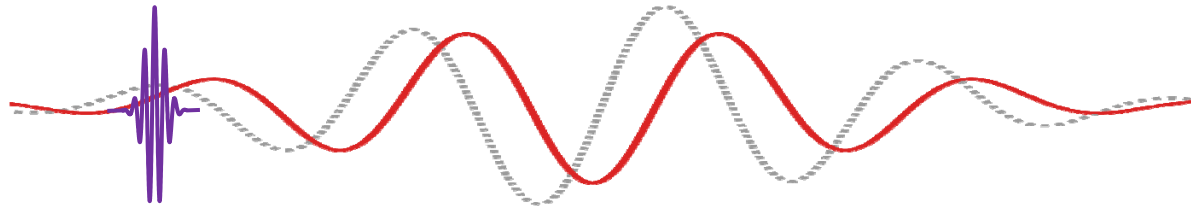


ionization of electrons in a solid



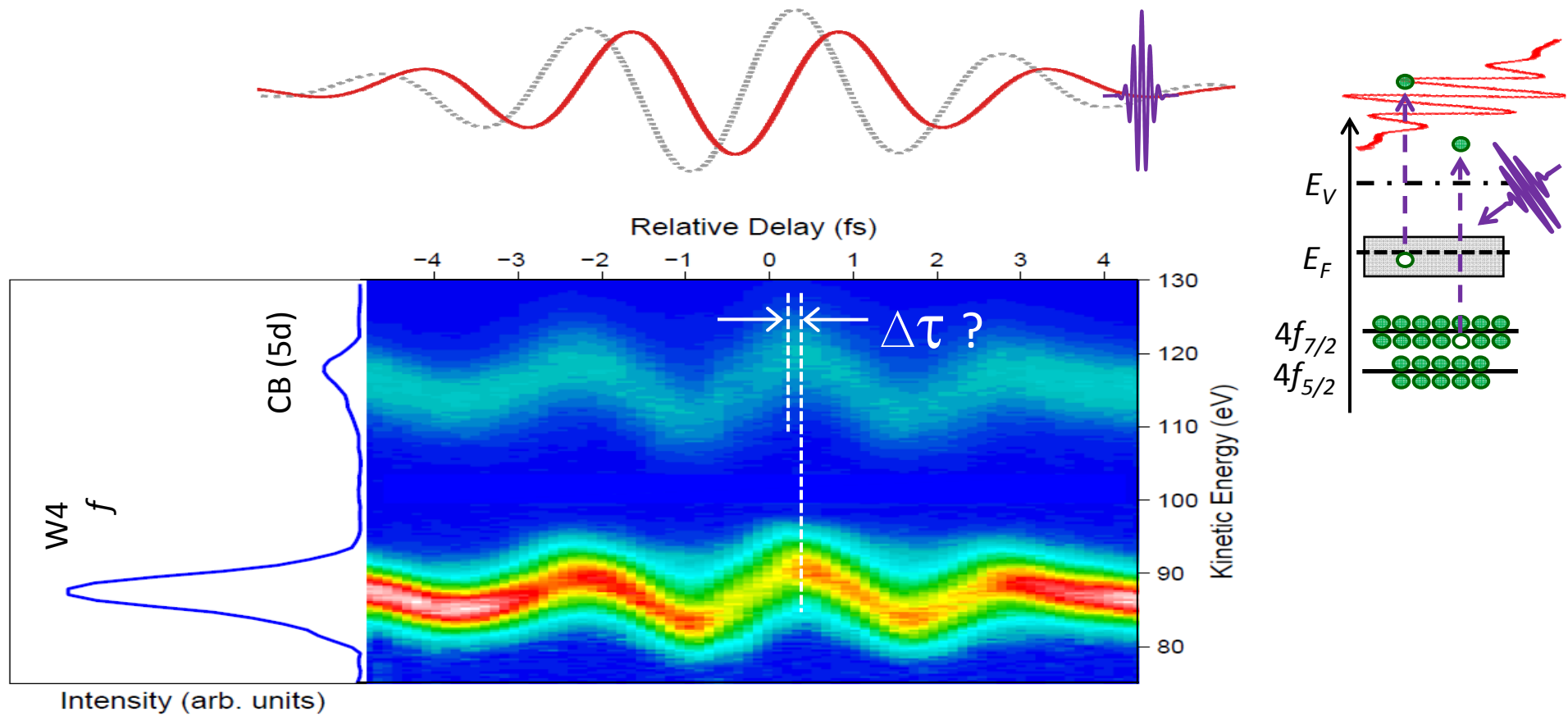


streaked photoemission from W(110)



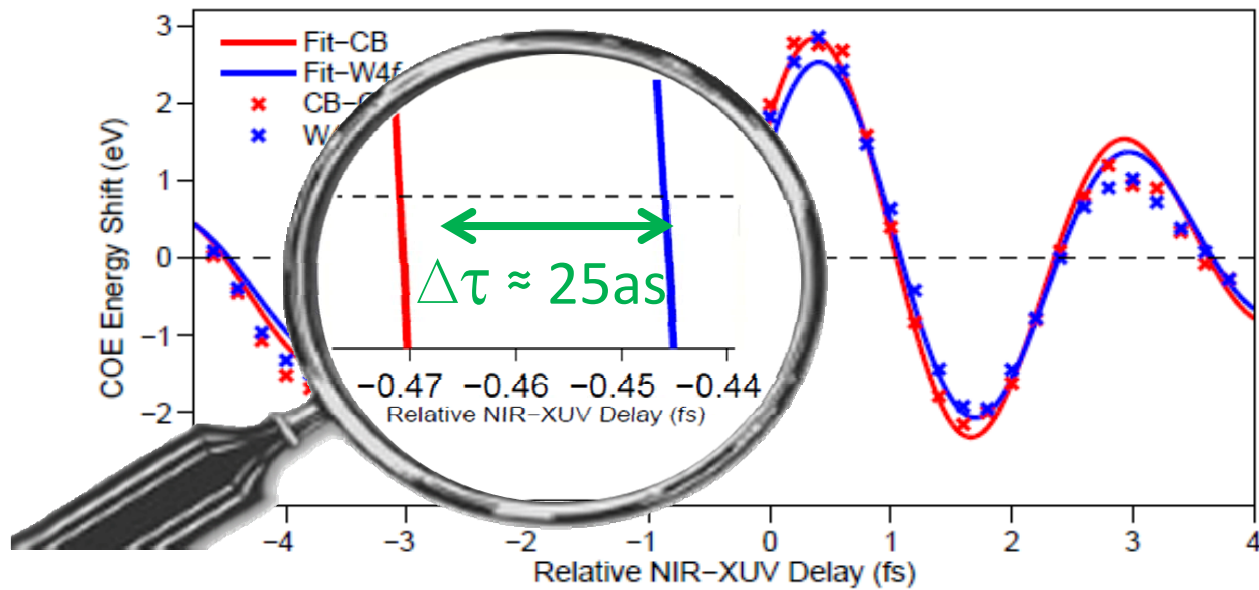


streaked photoemission from W(110)





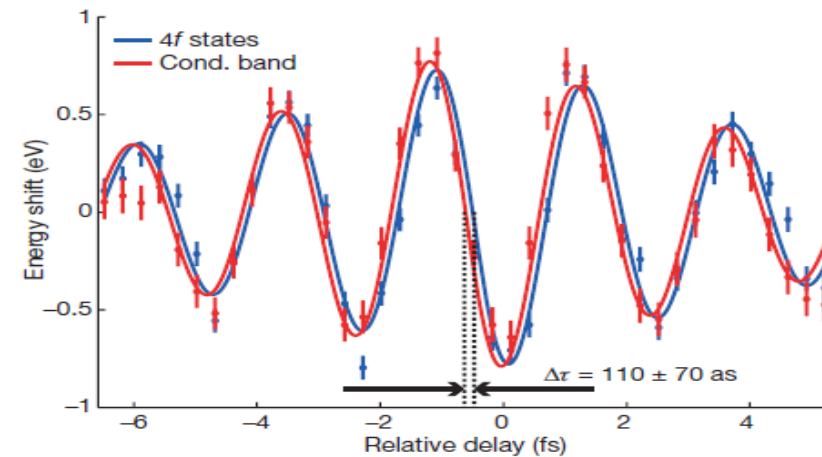
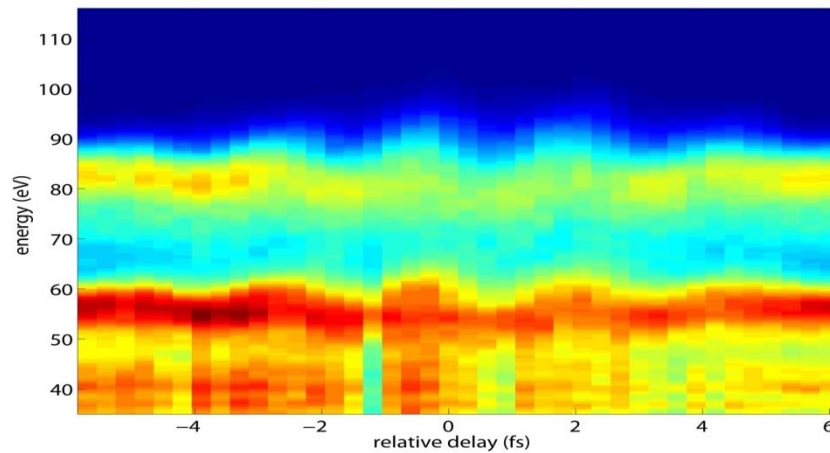
timing analysis



- Emission of the slower W4f core electrons is delayed



experiment from 2007



Sample: W(110) single crystal

Excitation: 300as XUV pulses centered at ~ 91 eV

- Emission of W4f electrons delayed by 110 ± 70 as with respect to emission of CB electrons (derived from a *single* measurement)





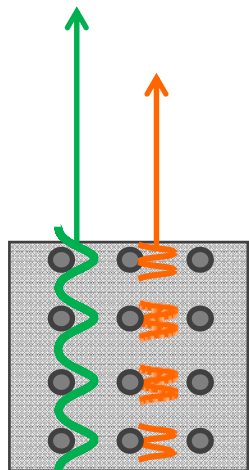
different ways of interpretation



Time delay in photoemission from solids is mainly due to:

...the difference in the spatial localization of the **initial-state** wave functions of CB and core-level electrons.

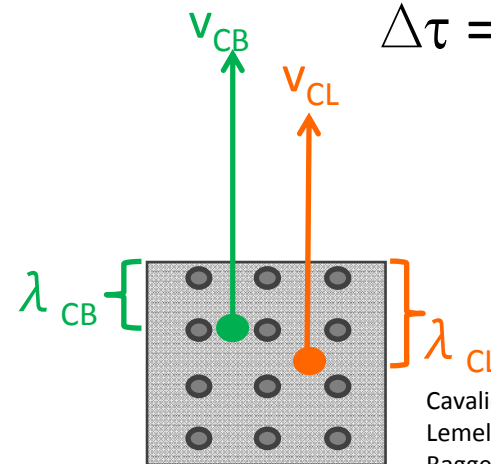
Surface states?



Zhang *et al*, PRA 84, 065403 (2011)
Zhang *et al*, PRL 102, 123601 (2009)
Kazansky *et al*, PRL 102, 177401 (2009)

...a propagation effect of the escaping electrons:

$$\Delta\tau = \lambda_{CL}/v_{CL} - \lambda_{CB}/v_{CB}$$



Impact of final-state band-structure?

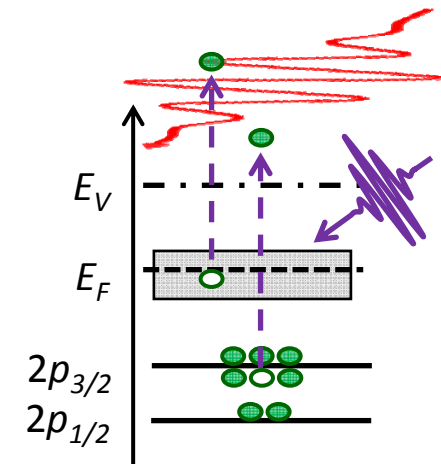
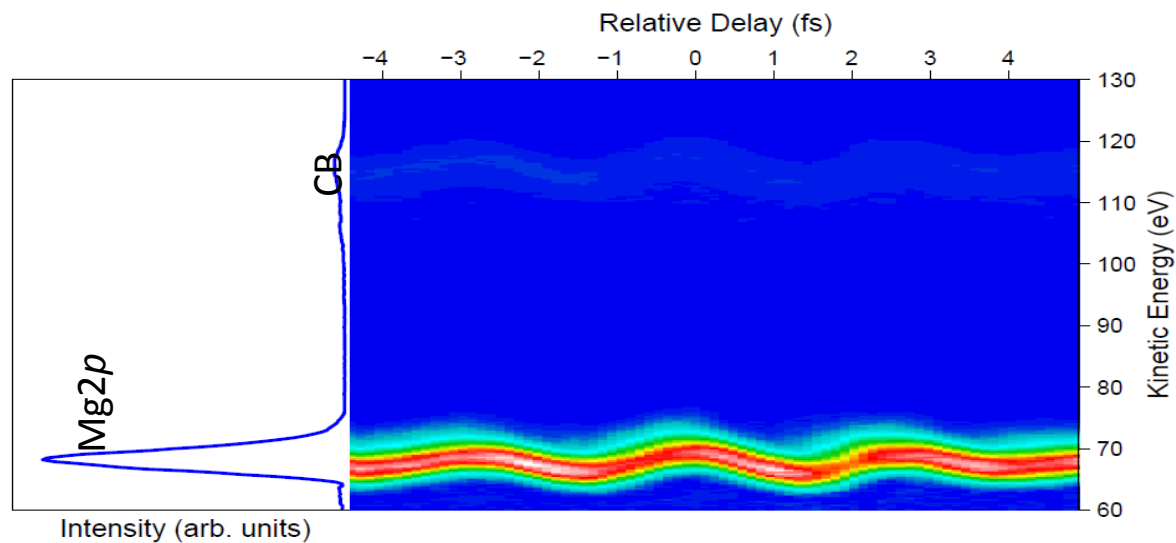
Cavalieri *et al*, Nature 449, 1029 (2007)
Lemell *et al*., PRA 79, 062901 (2009)
Baggesen *et al*., PRA 80, 030901 (2009)



attosecond streaking of Mg(0001)



- Ideal test case: truly delocalized CB electrons & localized $2p$ core-level electrons
- **Free-electron-like final-state band structure**
- Challenging: reactive material & low CB density of states

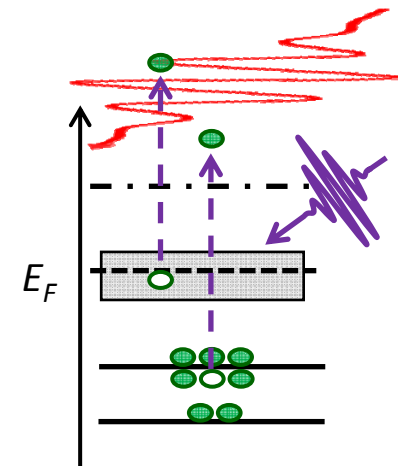
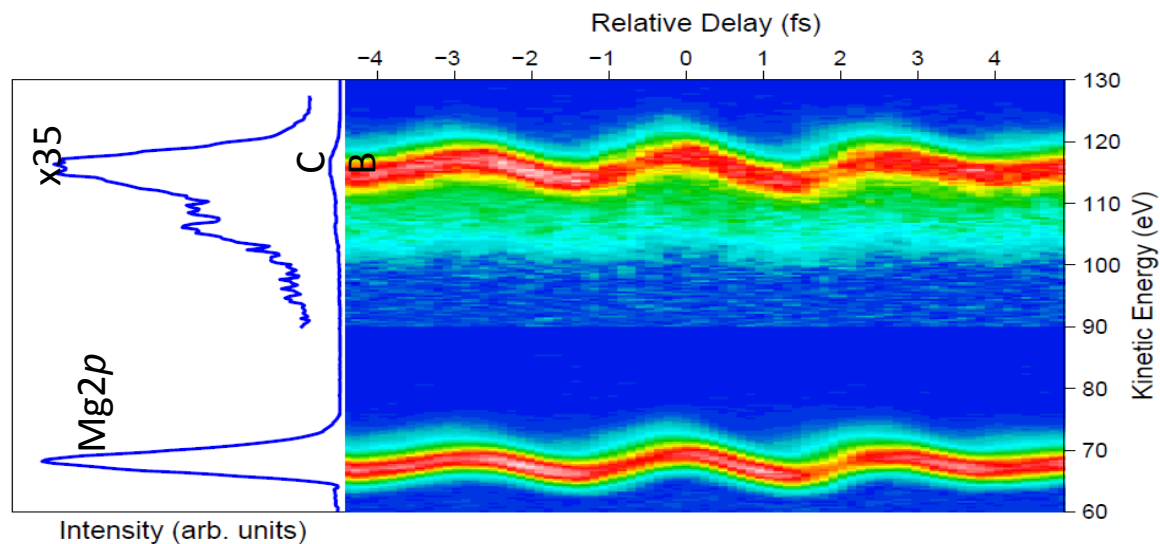




attosecond streaking of Mg(0001)

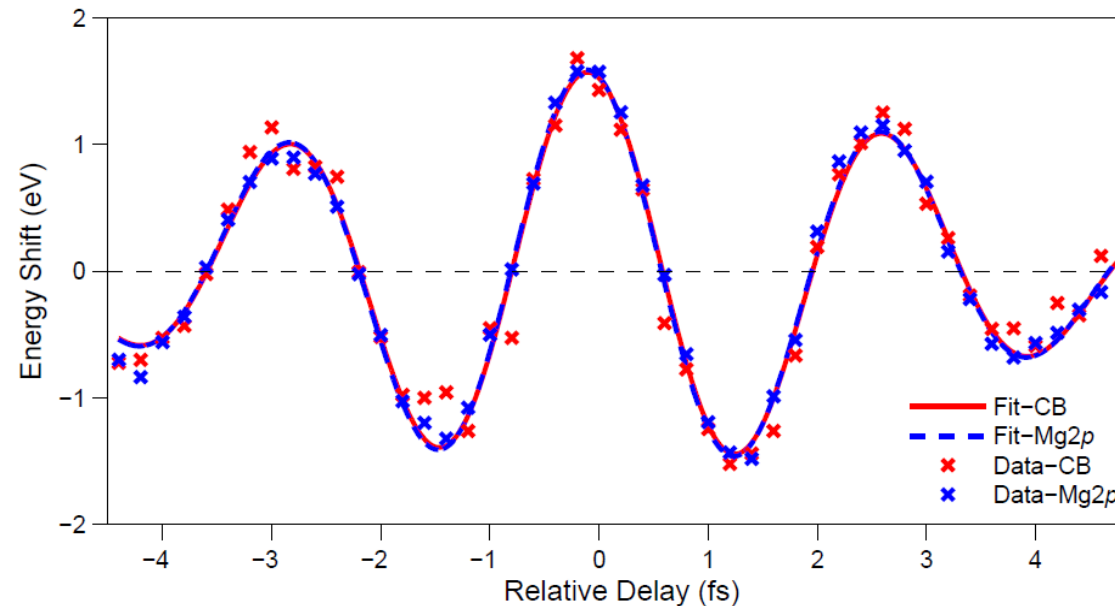


- Ideal test case: truly delocalized CB electrons & localized 2p core-level electrons
- **Free-electron-like final-state band structure**
- Challenging: reactive material & low CB density of states





attosecond streaking of Mg(0001)

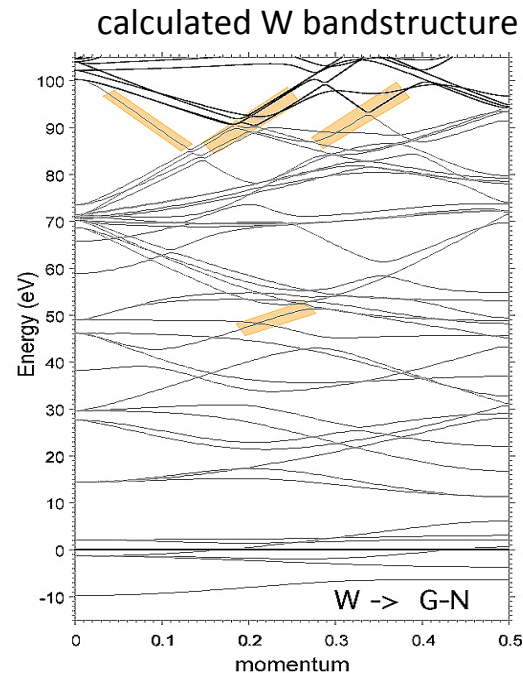


➤ Strong discrepancy with theories that emphasize the role of the initial-state localization ($\Delta\tau \rightarrow 100\text{as}$!)

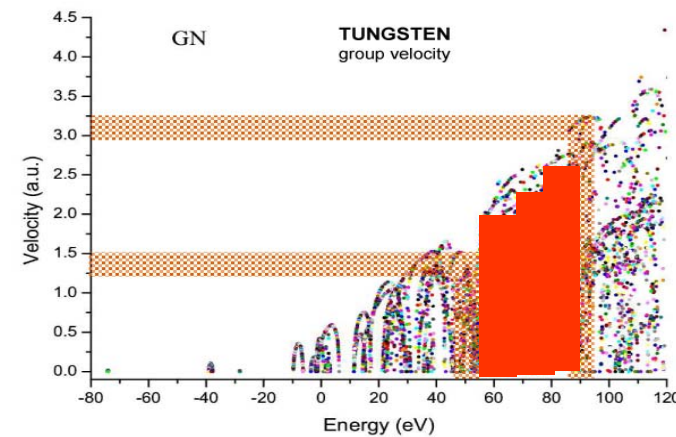
➤ Perfect agreement with the free-electron propagation limit: $\Delta\tau = \lambda_{\text{CB}}/v_{\text{CB}} - \lambda_{2p}/v_{2p} = 0\text{as}$



final-state bandstructure in the non-free-electron-like case



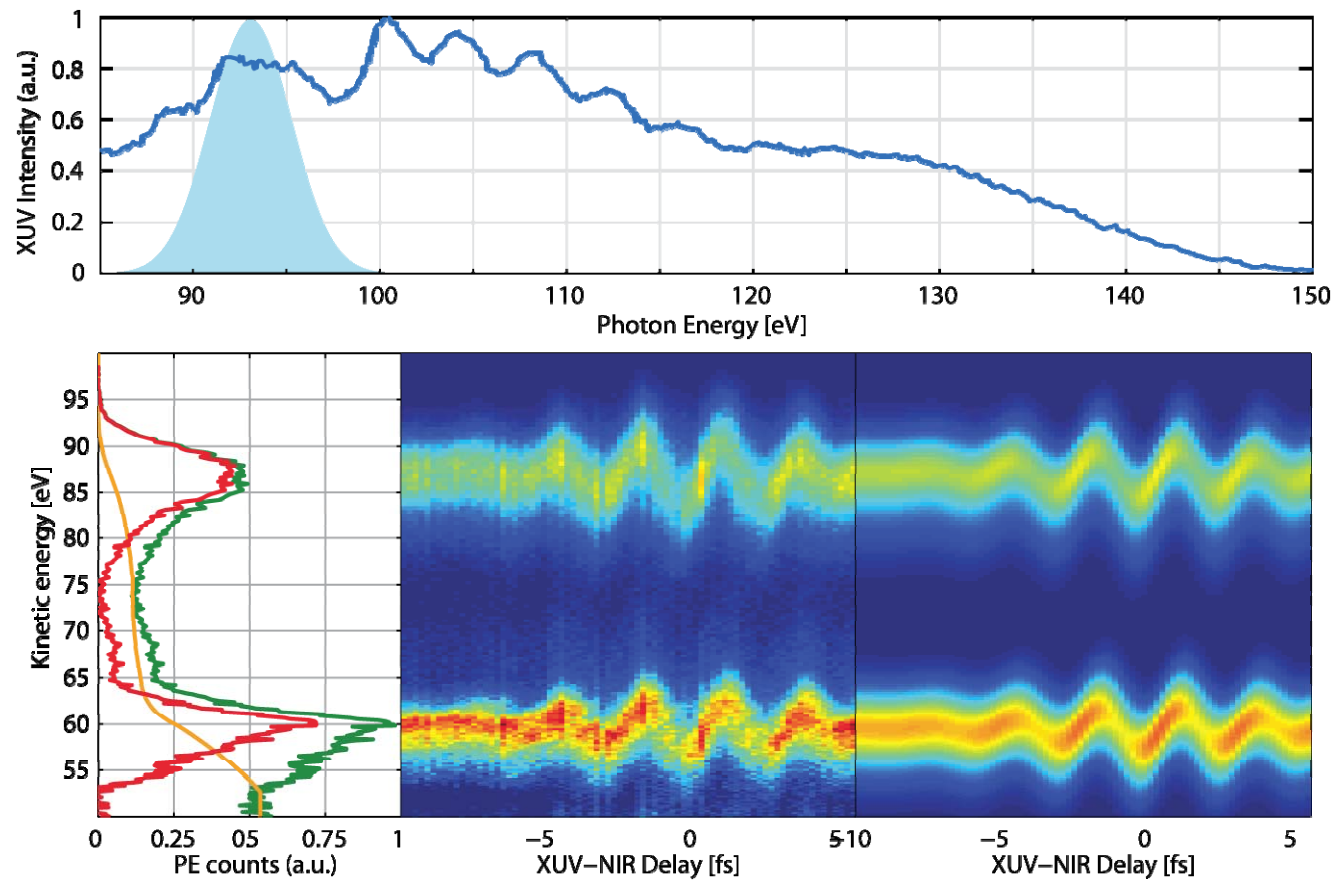
calculated group velocities for $h\nu = 95$ eV



- Conduction-band electrons (from near the Fermi energy) are excited into upper-conduction band states where their group velocity is $\sim$ twice that of the excited $4f$ -state electrons

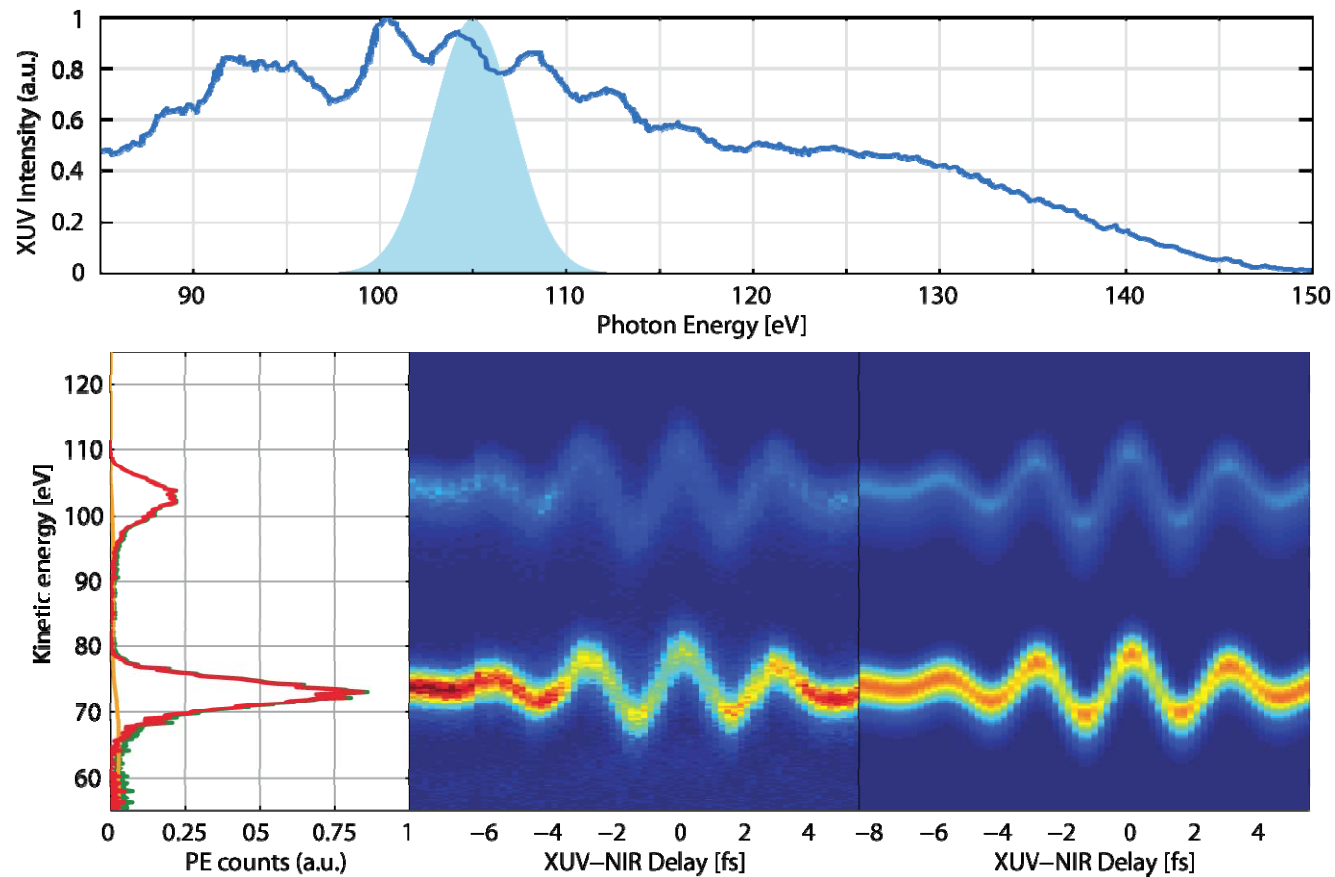


energy dependent *as* streaking spectroscopy



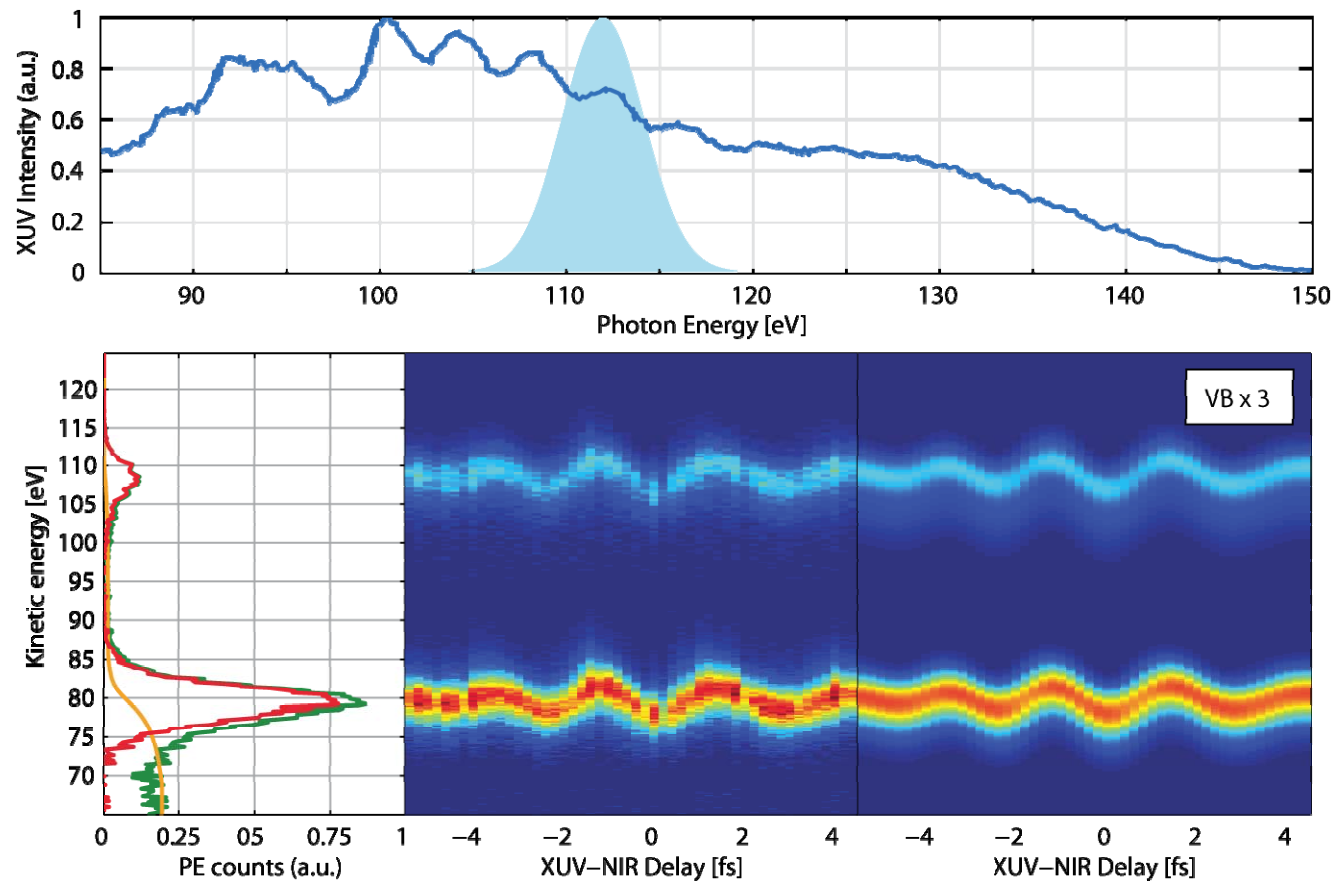


energy dependent *as* streaking spectroscopy



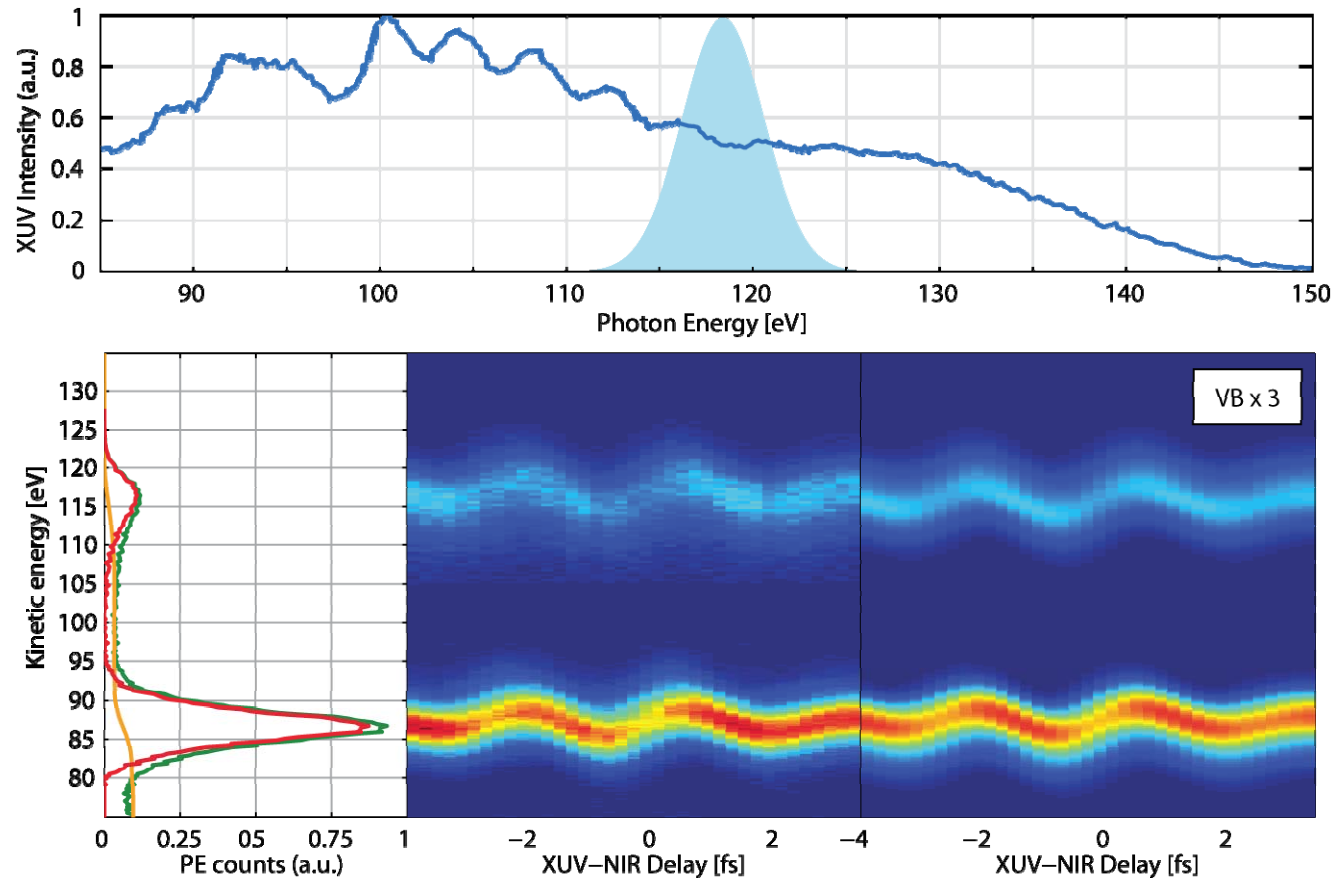


energy dependent *as* streaking spectroscopy



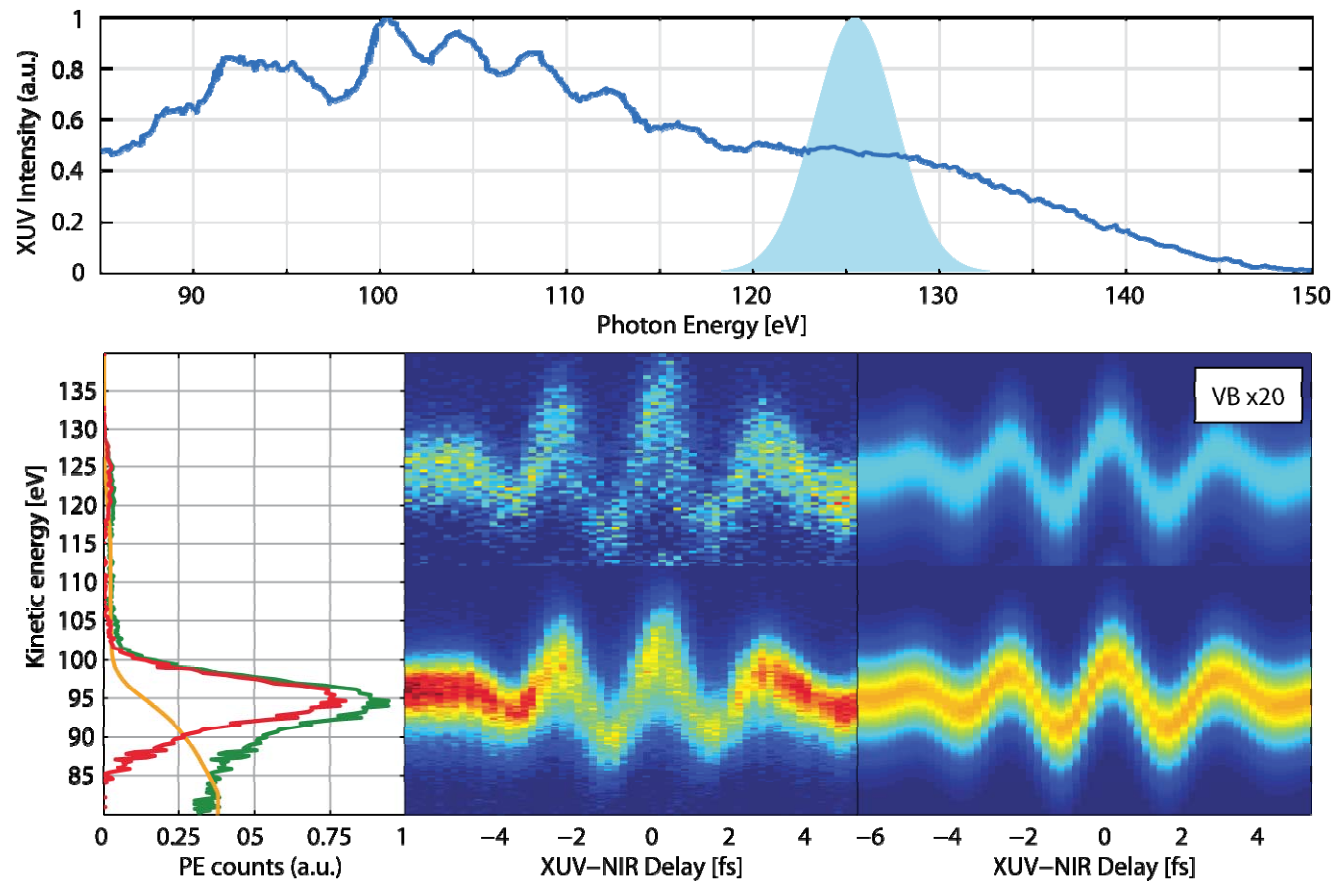


energy dependent *as* streaking spectroscopy



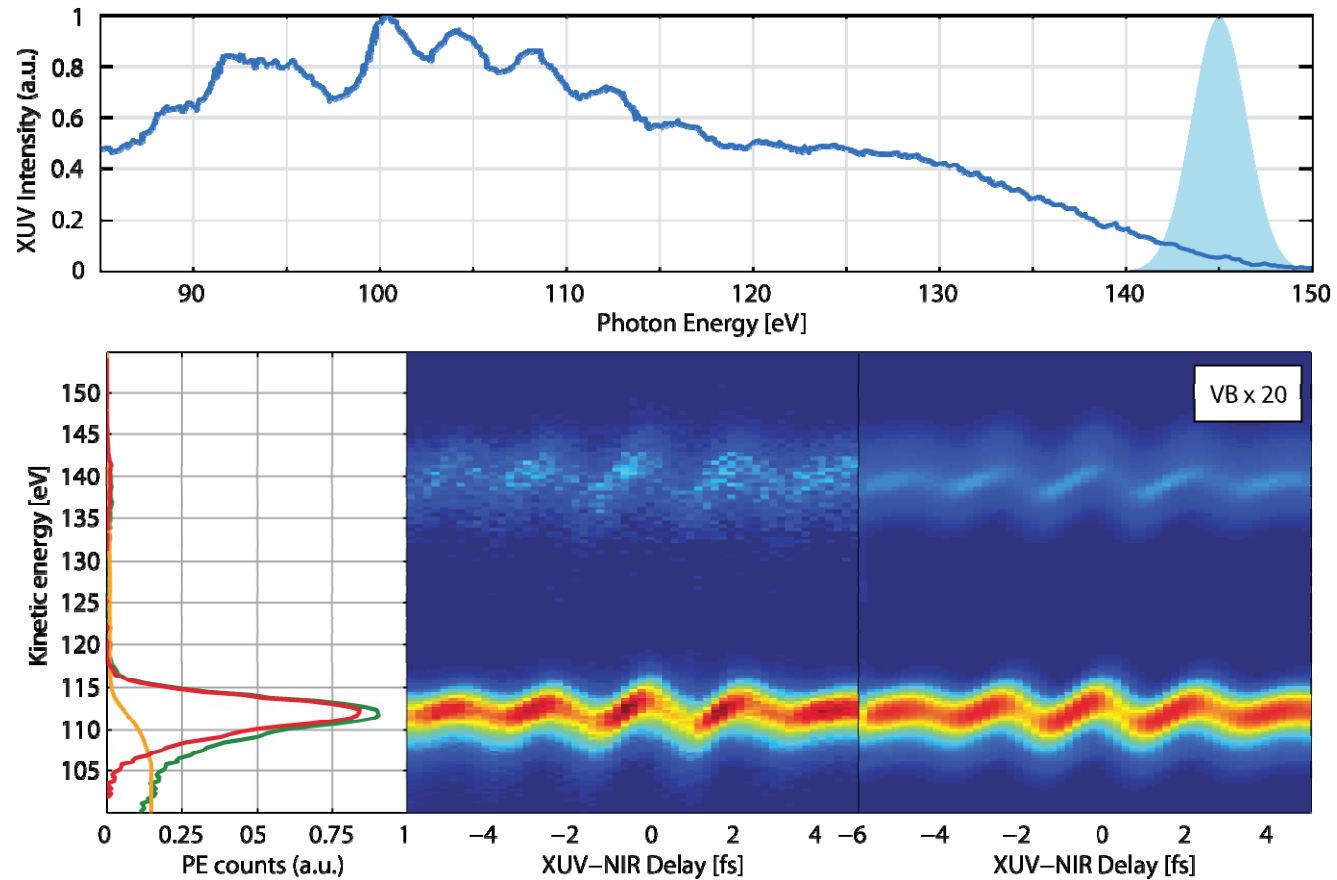


energy dependent *as* streaking spectroscopy



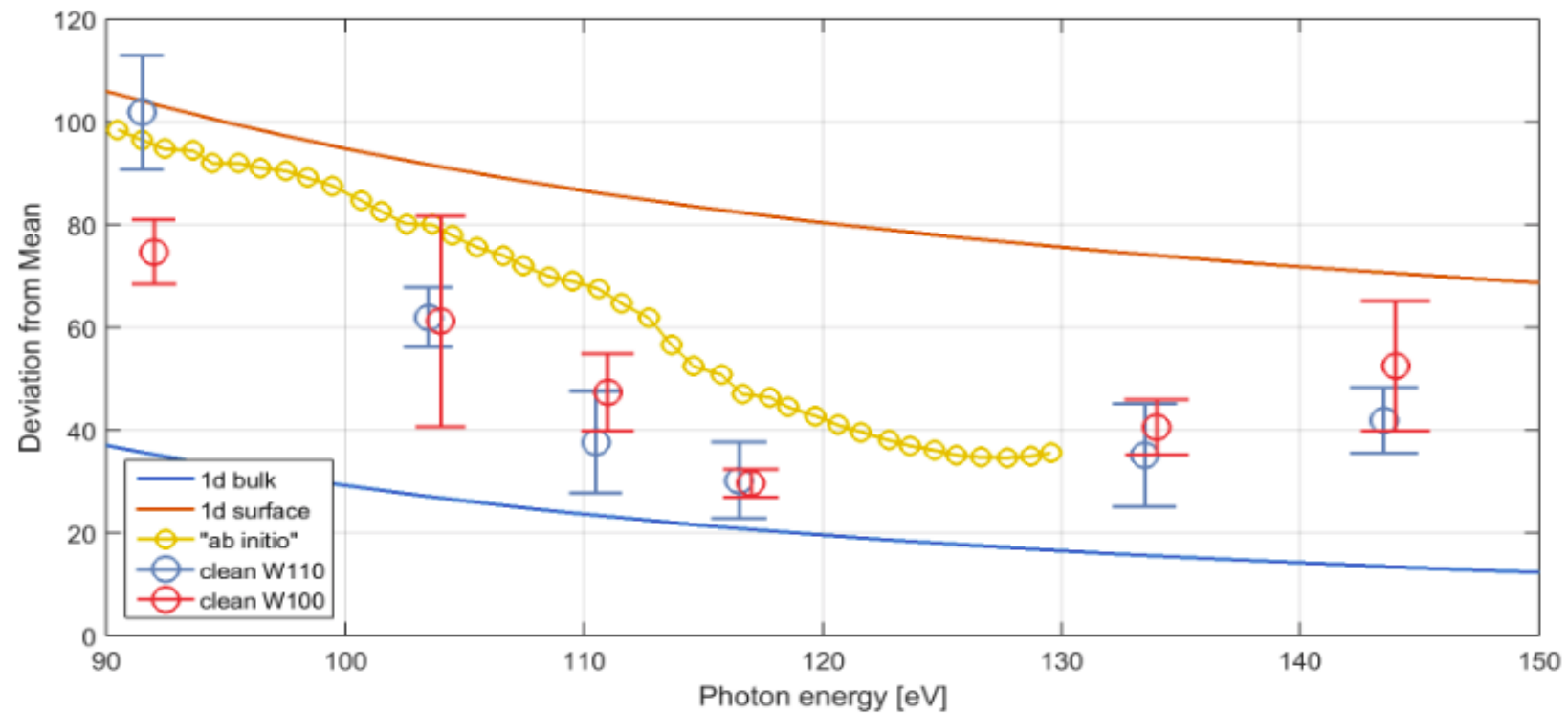


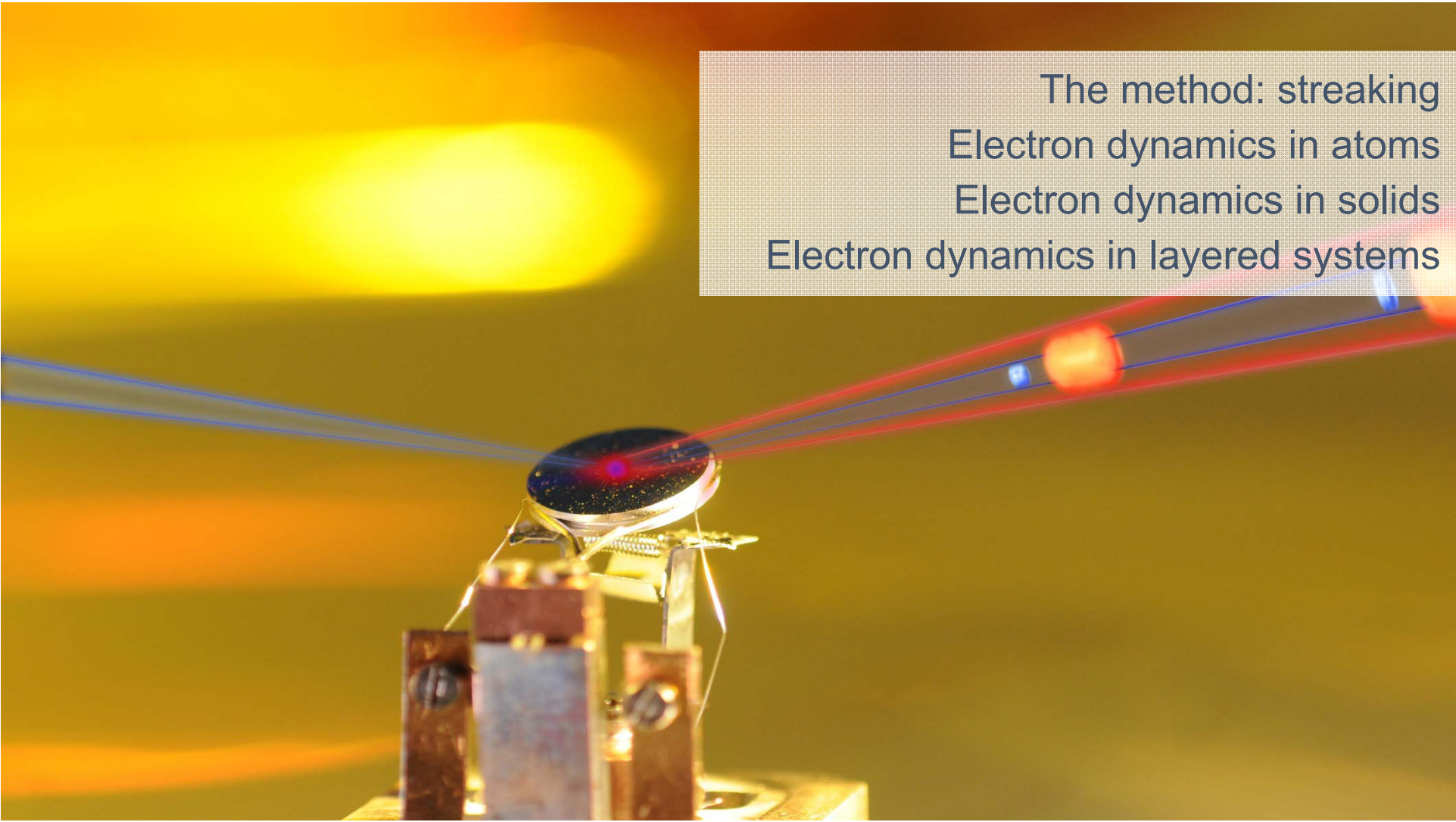
energy dependent *as* streaking spectroscopy





delay: simulation (San Sebastian) vs. experiment





The method: streaking
Electron dynamics in atoms
Electron dynamics in solids
Electron dynamics in layered systems

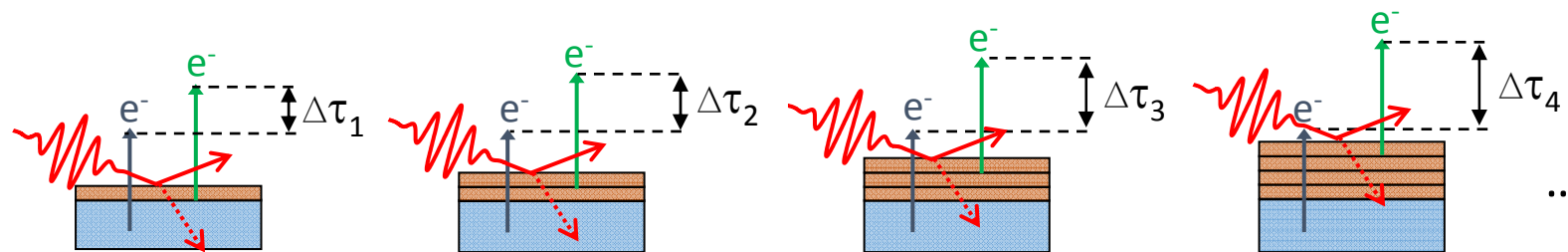
The background image shows a laboratory setup for streaking experiments. A bright yellow laser beam is directed at a small, dark, circular sample mounted on a metal stand. Two blue lines represent the streaking laser pulses, which are used to measure the time delay between the pump and probe pulses. The diagram overlaid on the image shows a red sphere (representing an electron) being accelerated by a blue line (representing the streaking laser pulse) to a final position (blue sphere) after a time delay.



real-time observation of electron propagation



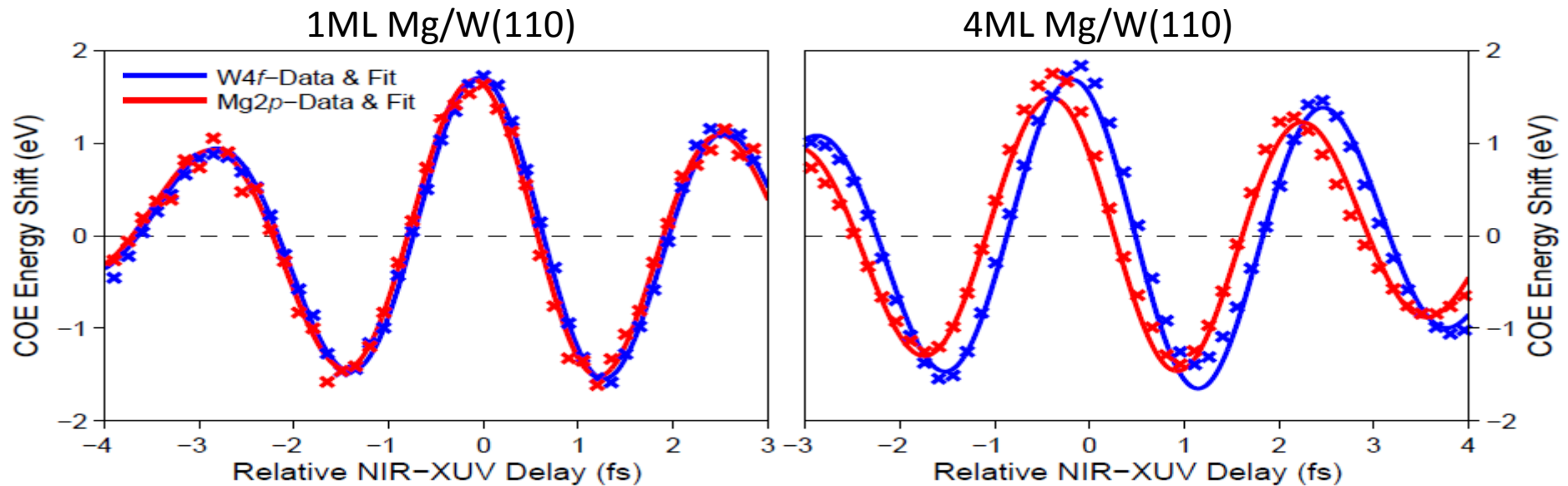
- Tuning the transport experienced by substrate electrons by introducing additional metal overlayers: $\Delta\tau_1 < \Delta\tau_2 < \Delta\tau_3 < \Delta\tau_4 \dots$?



- Investigated system: magnesium layers epitaxially grown on W(110)
- Simultaneous observation of W4f, Mg2p and CB photoelectrons possible



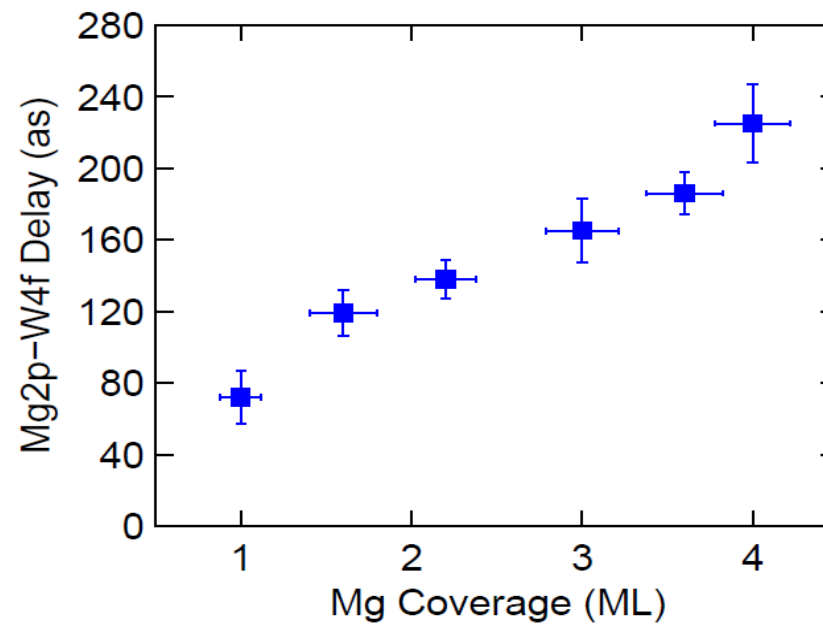
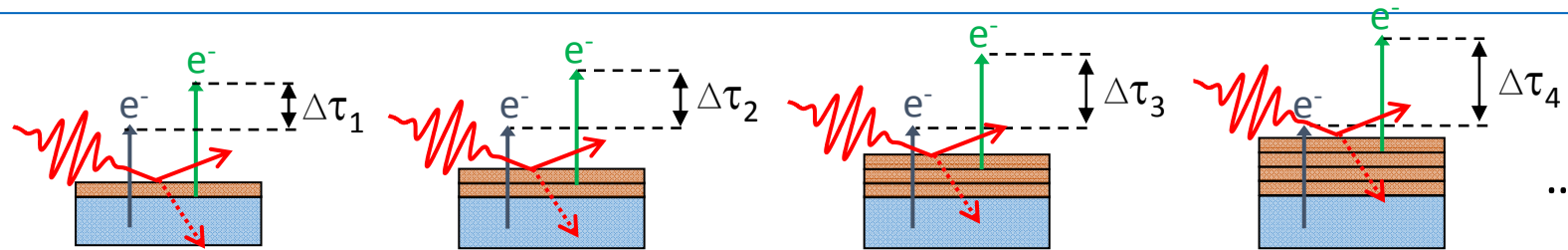
delay of W4f electrons in Mg/W(110)



- Mg2p-W4f time delay significantly increases as a function of the Mg overlayer thickness



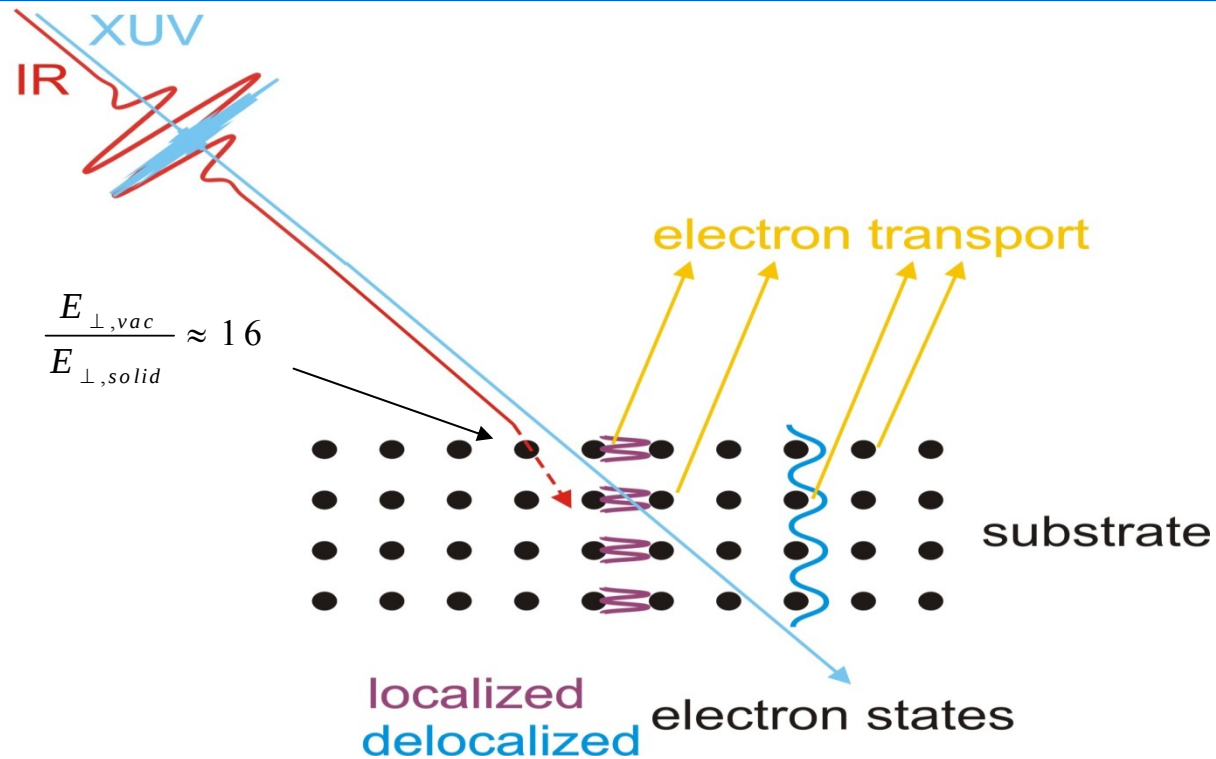
delay of W4f electrons in Mg/W(110)



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penetration of the normal field component



assumption: streaking acts outside the solid



propagation model for Mg/W(110)



$$\Delta\tau(d) = \tau_{2p} - \tau_{4f} = \frac{\langle z \rangle_{2p}}{v_{2p}} - \frac{\langle z \rangle_{4f} + d - \delta}{v_{4f}}$$

free electron dispersion

$$= \frac{1}{v_{2p}} \cdot \frac{\int_{\delta}^d z e^{-z/\lambda_1} dz}{\int_0^d e^{-z/\lambda_1} dz} - \frac{1}{v_{4f}} \cdot \left(\frac{\int_d^{\infty} (z-d) e^{-(z-d)/\lambda_2} dz}{\int_d^{\infty} e^{-(z-d)/\lambda_2} dz} + d - \delta \right)$$

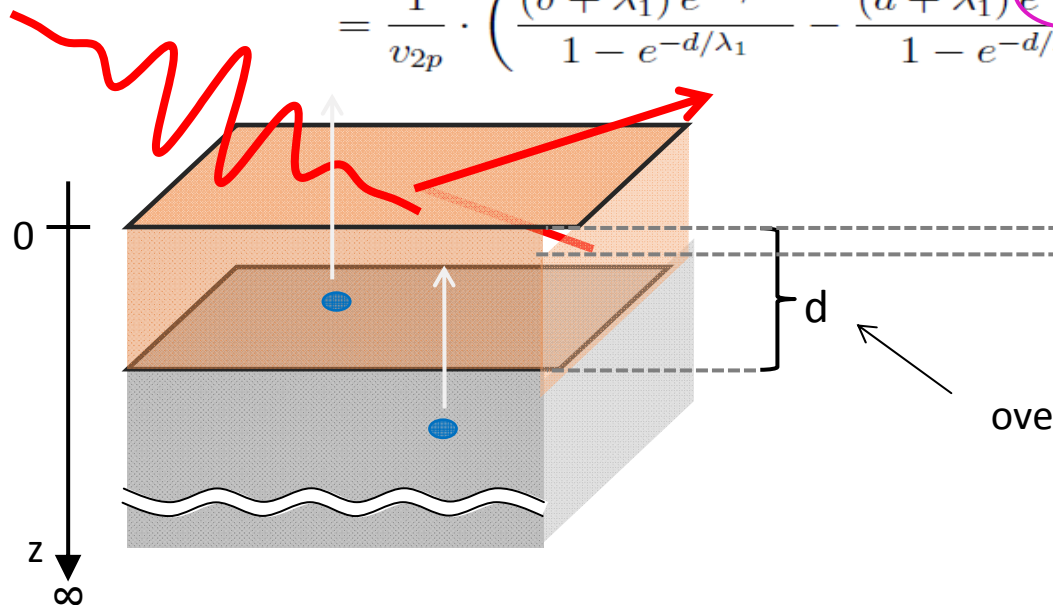
$$= \frac{1}{v_{2p}} \cdot \left(\frac{(\delta + \lambda_1) e^{-\delta/\lambda_1}}{1 - e^{-d/\lambda_1}} - \frac{(d + \lambda_1) e^{-d/\lambda_1}}{1 - e^{-d/\lambda_1}} \right) - \frac{\lambda_2 + d - \delta}{v_{4f}}$$

IMFP of W4f electrons in W (3.9Å)

IMFP of Mg2p electrons in Mg (3.7Å)

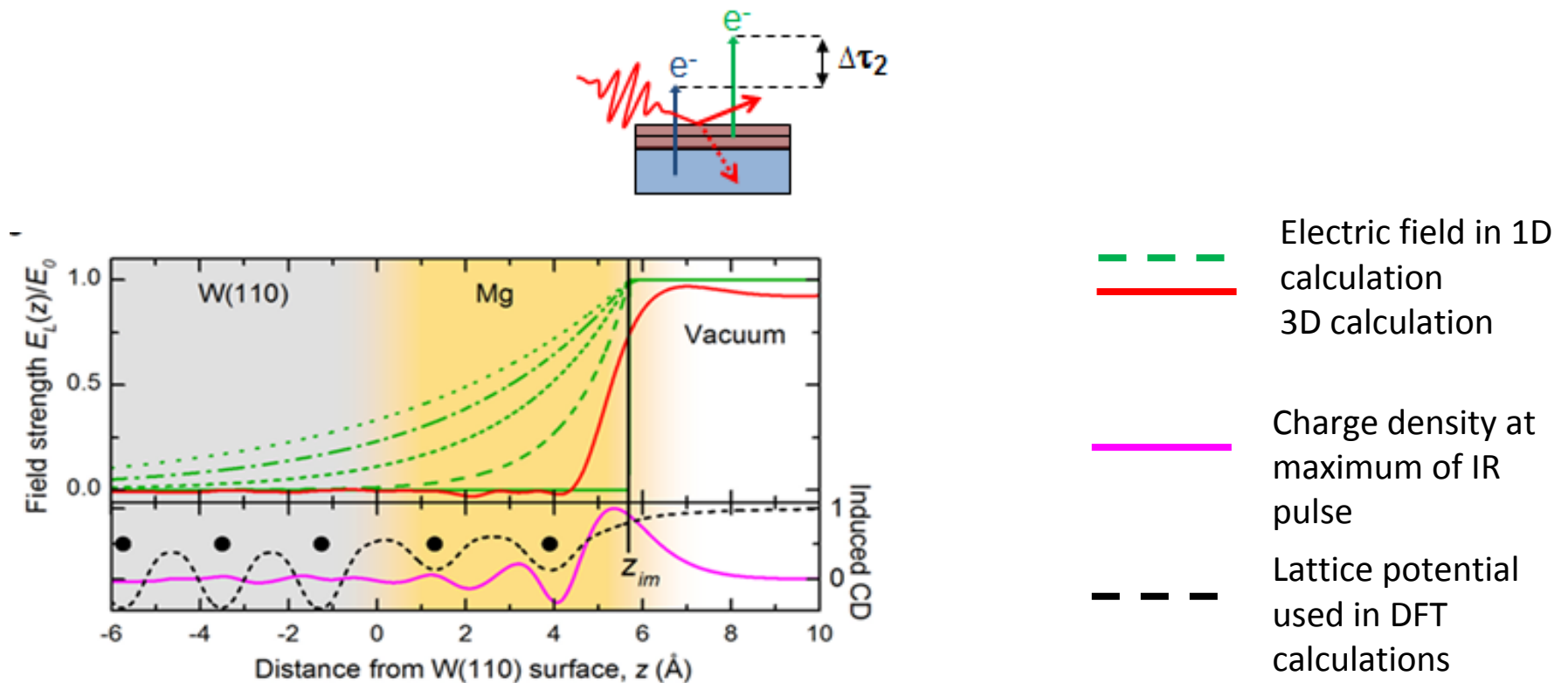
'effective' NIR screening length

overlayer thickness





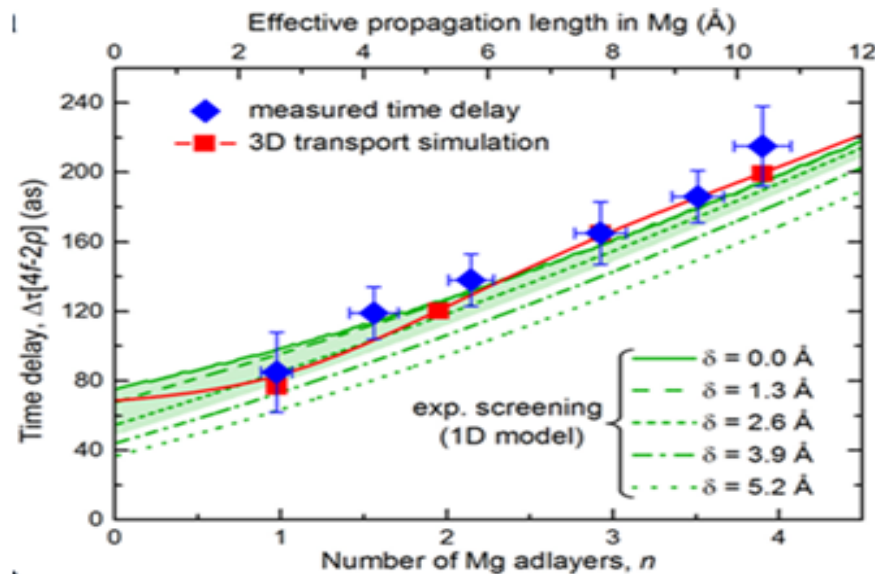
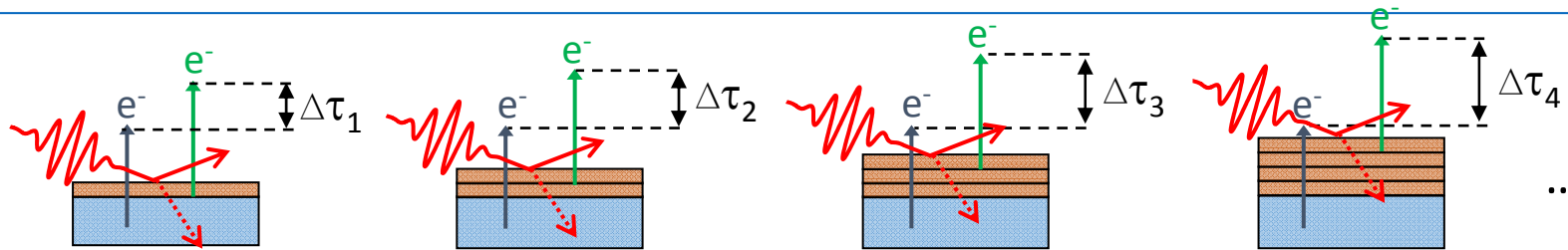
Delay of W4f electrons in Mg/W(110)



S. Neppl, al., *Nature* **517**, 342 (2015)



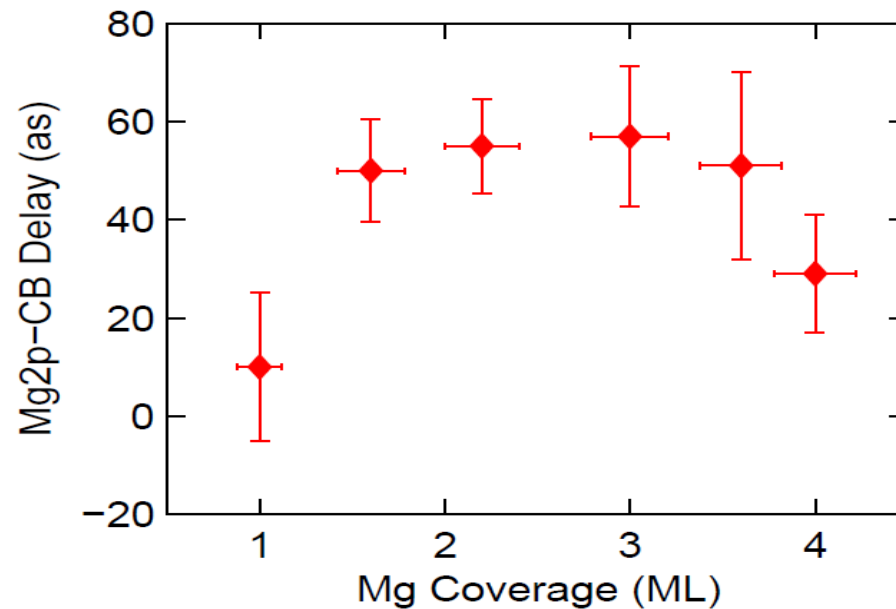
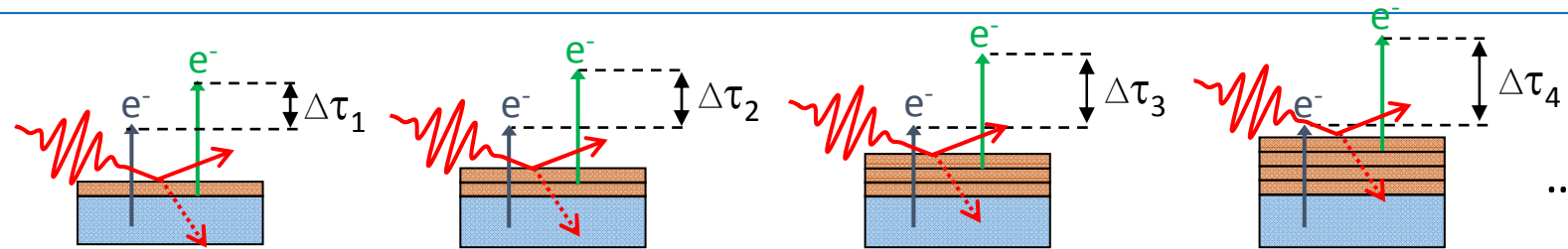
delay of W4f electrons in Mg/W(110)



exp. decaying field
& 3D calculations:
screening 1.3 - 2.6 Å



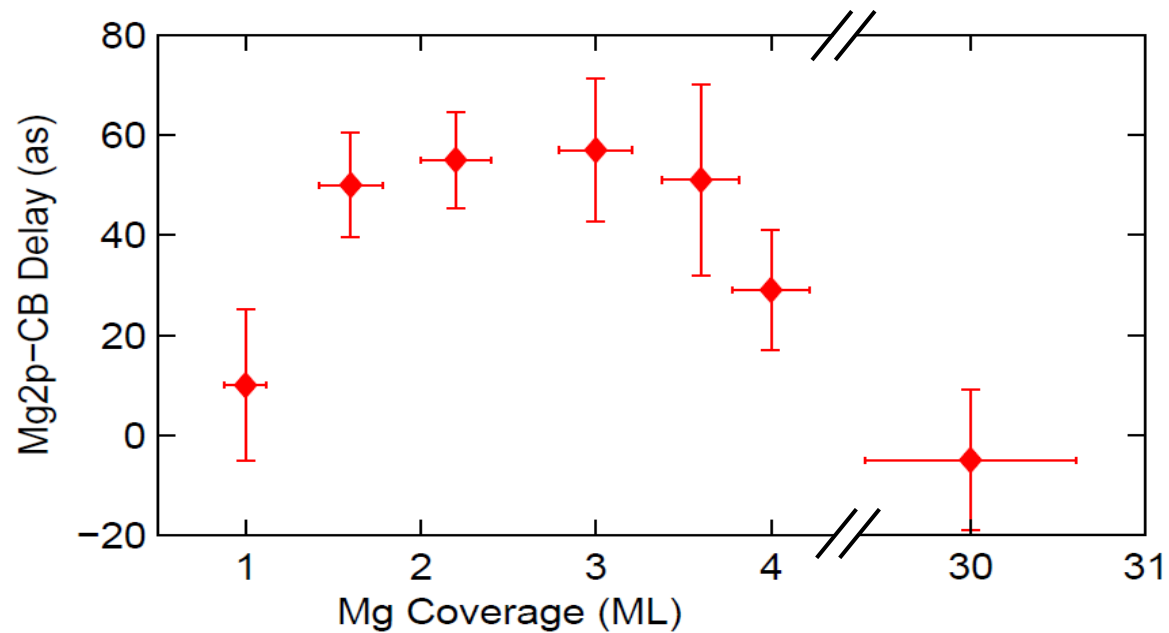
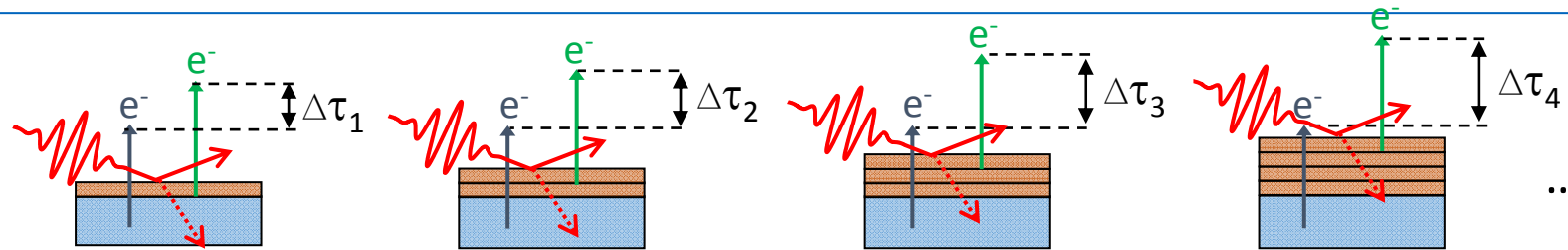
delay of **CB** electrons in Mg/W(110)



- More complex: superposition of W and Mg CB states
- Delay should converge to bulk Mg(0001) value for increasing coverage

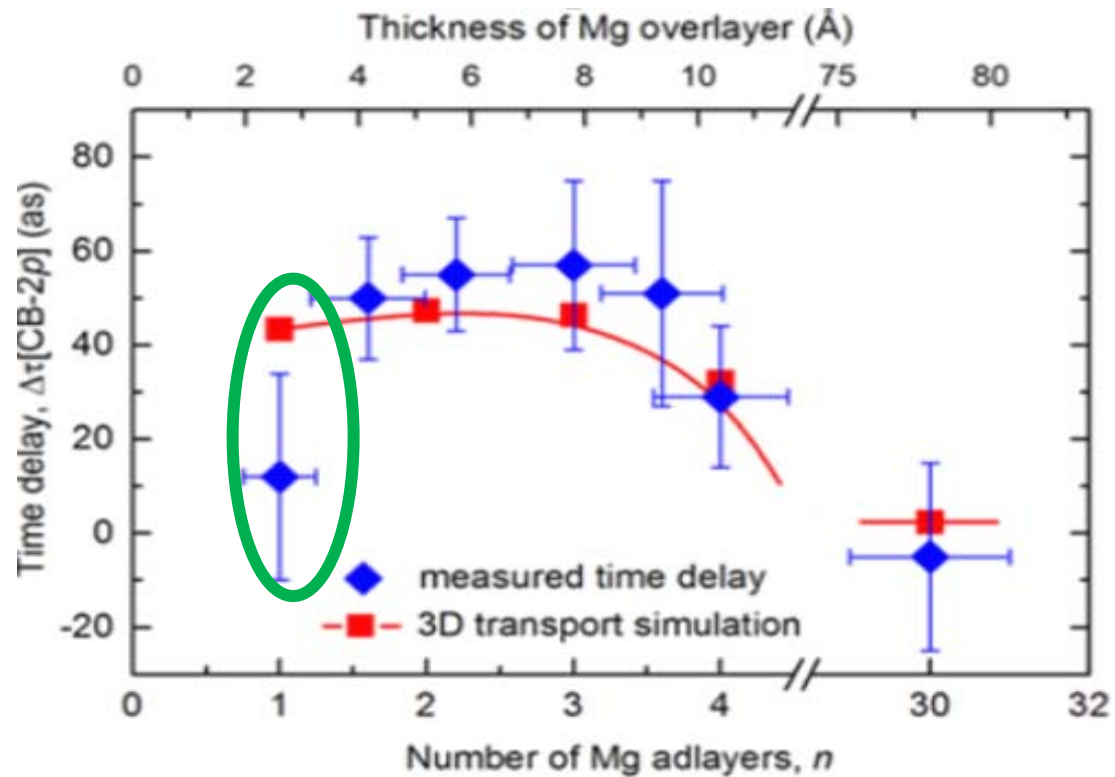


delay of **CB** electrons in Mg/W(110)





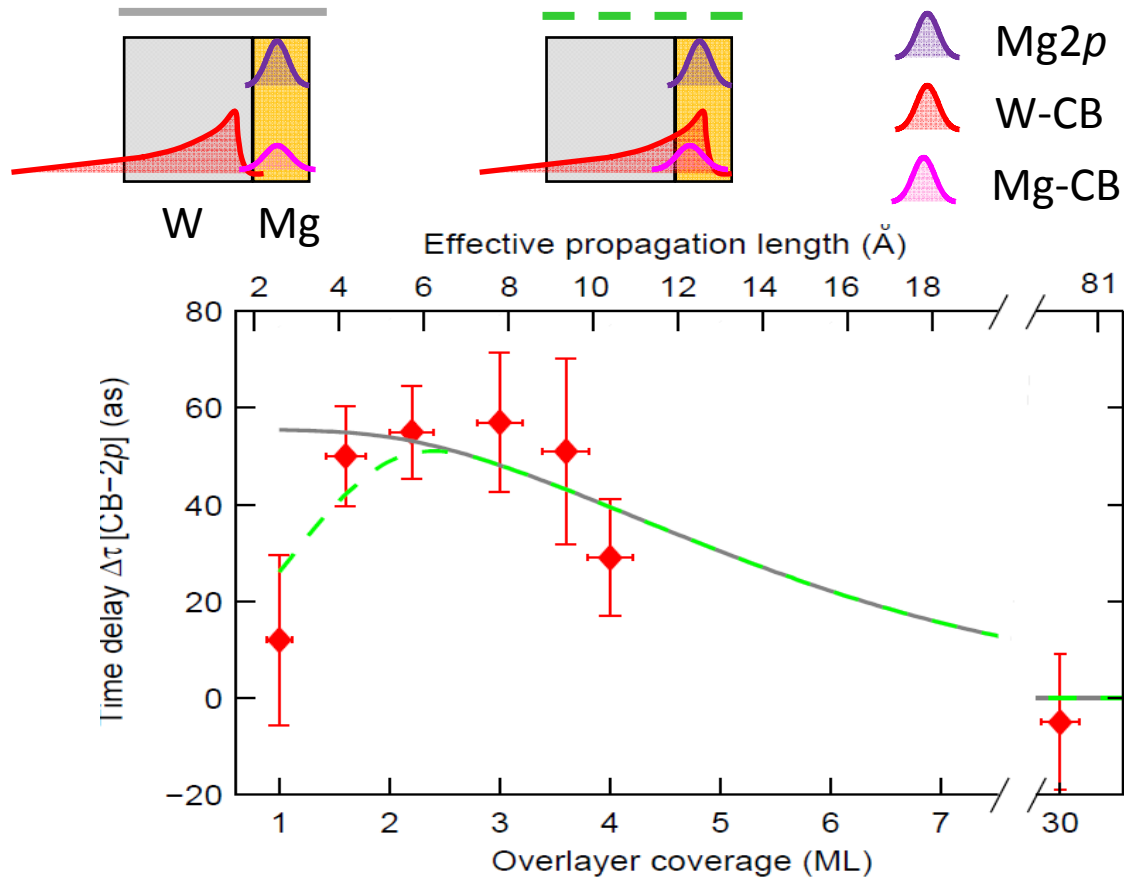
delay of **CB** electrons in Mg/W(110) 3D transport simulation



S. Neppl, al., *Nature* **517**, 342 (2015)



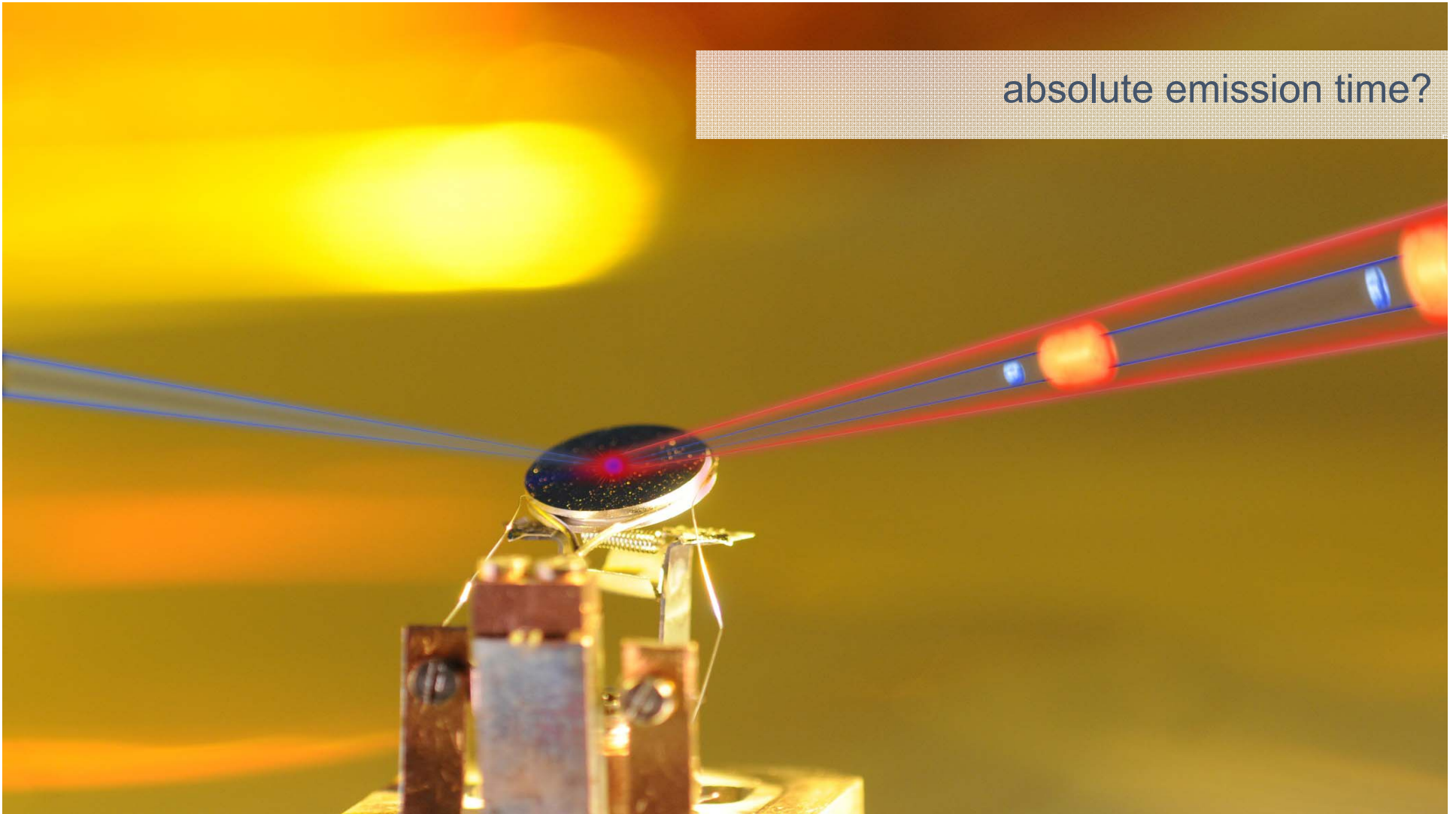
Delay of **CB** electrons in Mg/W(110)



hybridization?

penetration of the
W-CB electron
wavefunction
into the adlayer

absolute emission time?

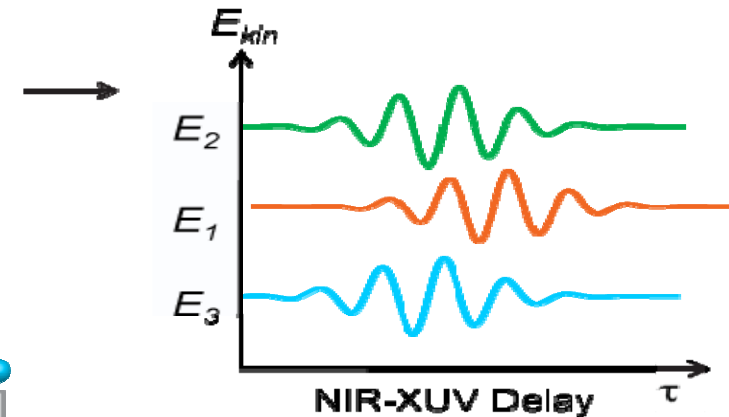
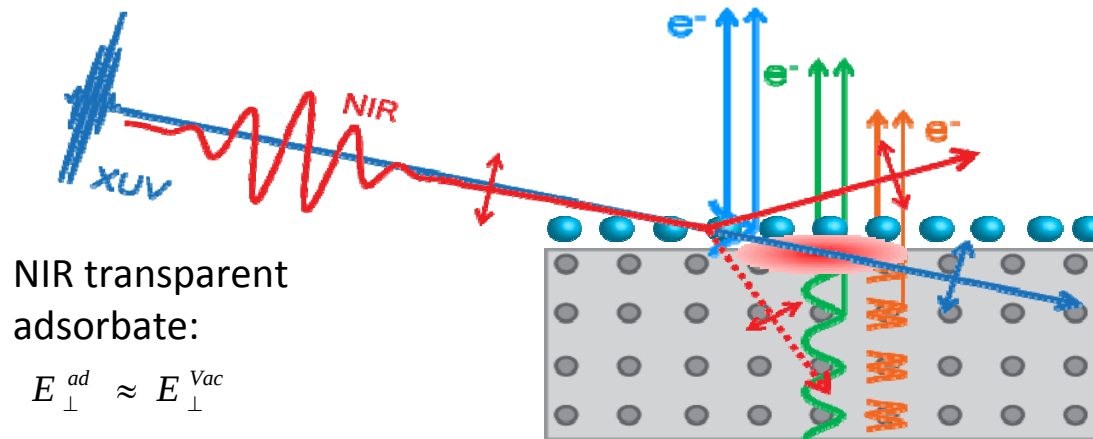




What about absolute emission times? – streaking of adsorbed iodine



The chronoscope / „time-zero-marker“



- Photoelectrons released from the adsorbate immediately respond to the streaking field at the surface
→ time “zero” for clocking substrate emission

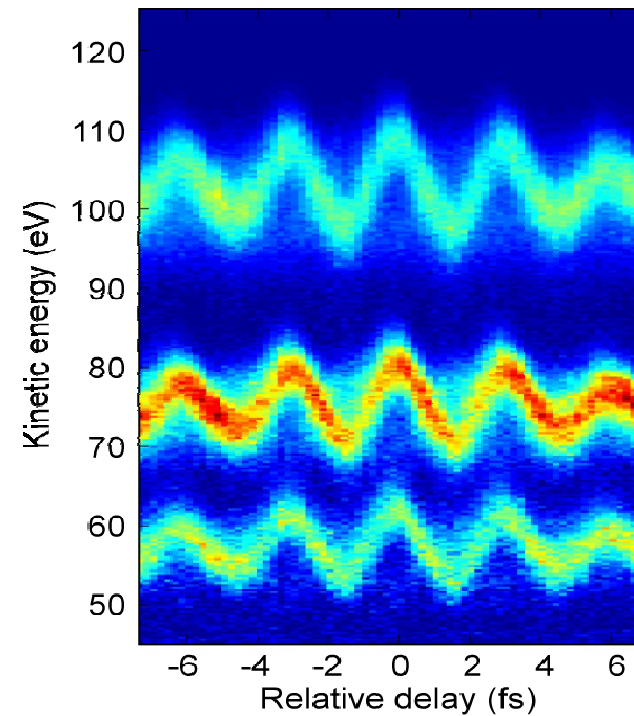


streaking experiments on I/W(110)



✓ Ideal system for streaking experiments:

- Sufficient energy separation
- Large cross sections
- Ease of sample preparation
- Long-term stable

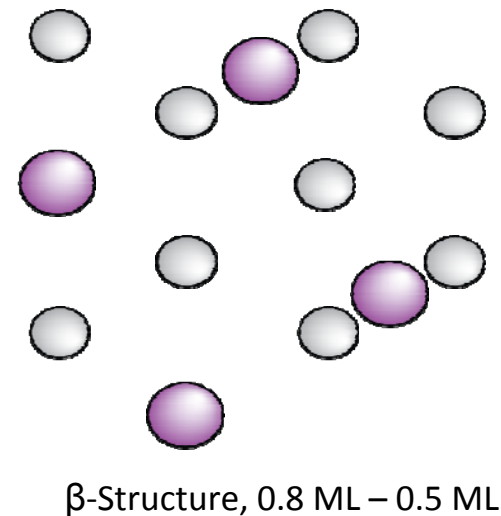
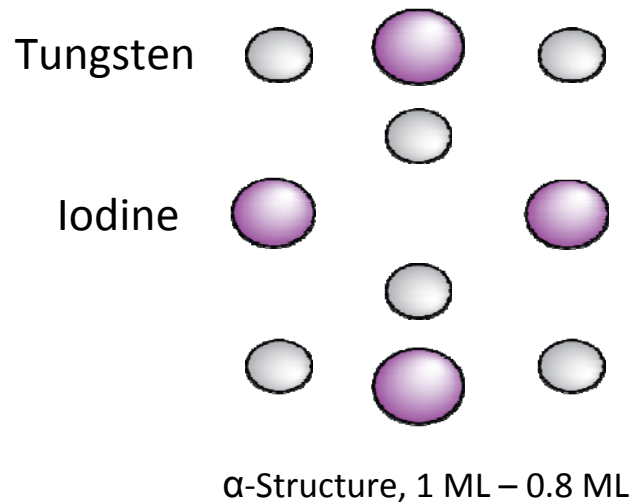




sub-monolayer structures of Iodine/Tungsten (110)



The surface coverage of the I/W(110) adlayer system is varied via thermal desorption after the growth of a saturated atomic monolayer at room temperature.



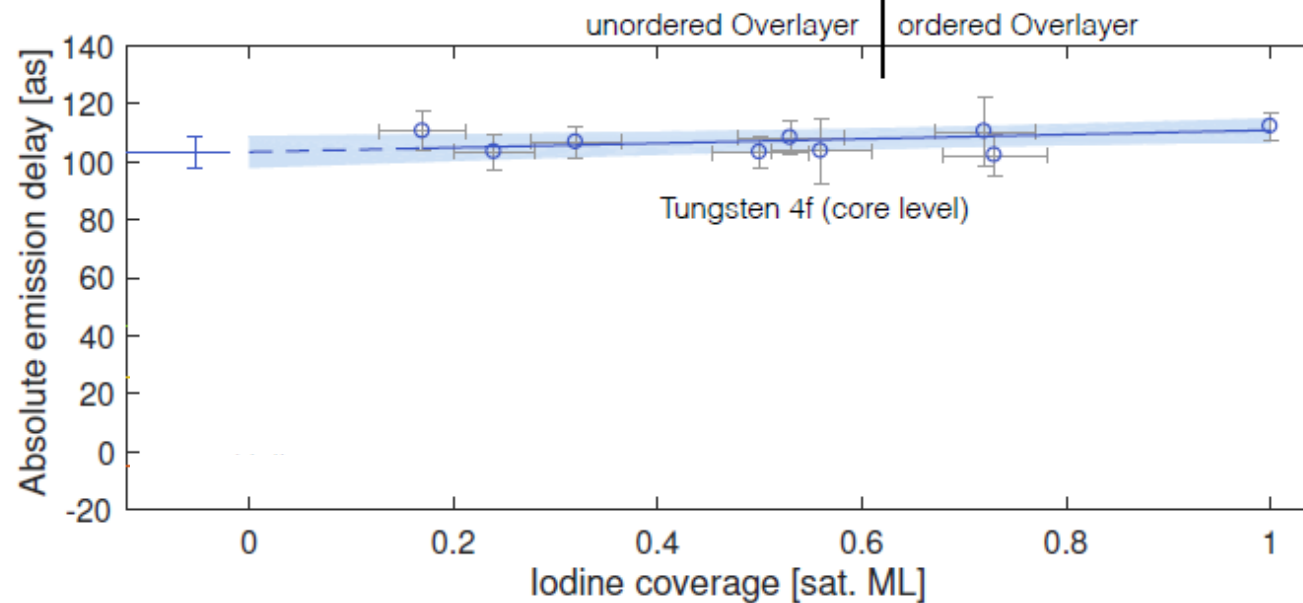
For a further increase in desorption temperature island formation is expected.



Extrapolation towards zero coverage – the undisturbing chronograph



Even for submonolayer adsorbate coverage residual effects of transport and electric field screening are expected to influence the observed delay → Extrapolation to 0 layers!



$$\tau_{W4f} = (103 \pm 6) \text{as} + \theta(8 \pm 7) \text{as}$$

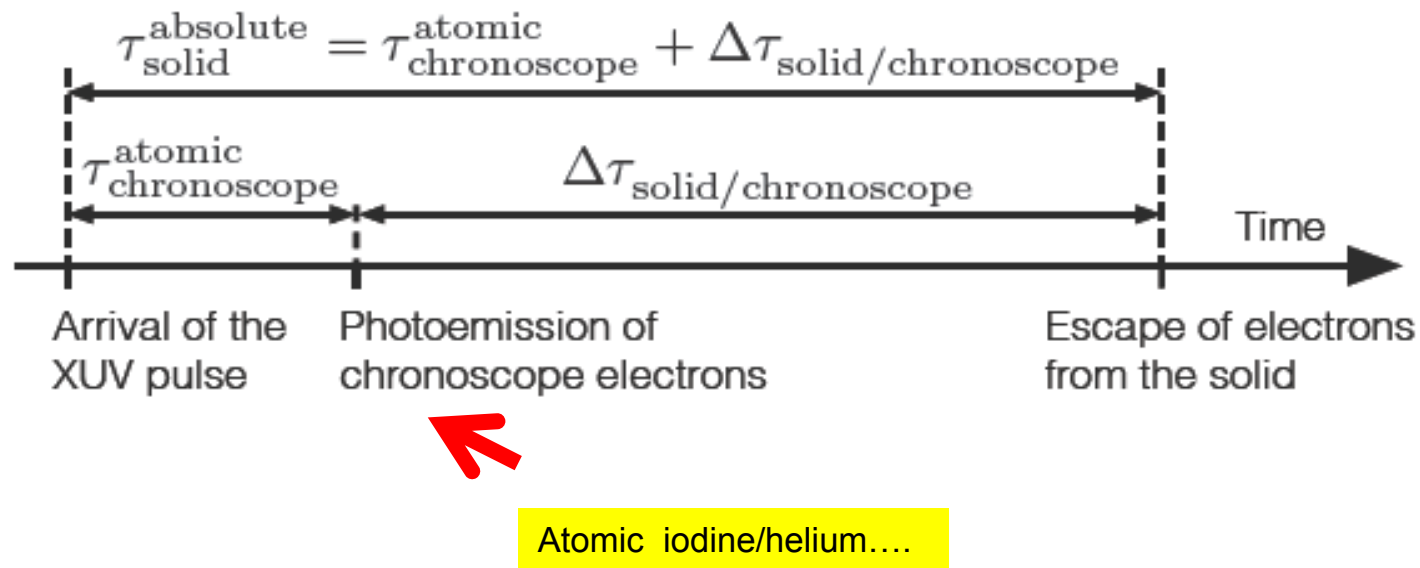
└──────────┘ Coverage in ML



Extrapolation towards zero coverage – the undisturbing chronograph

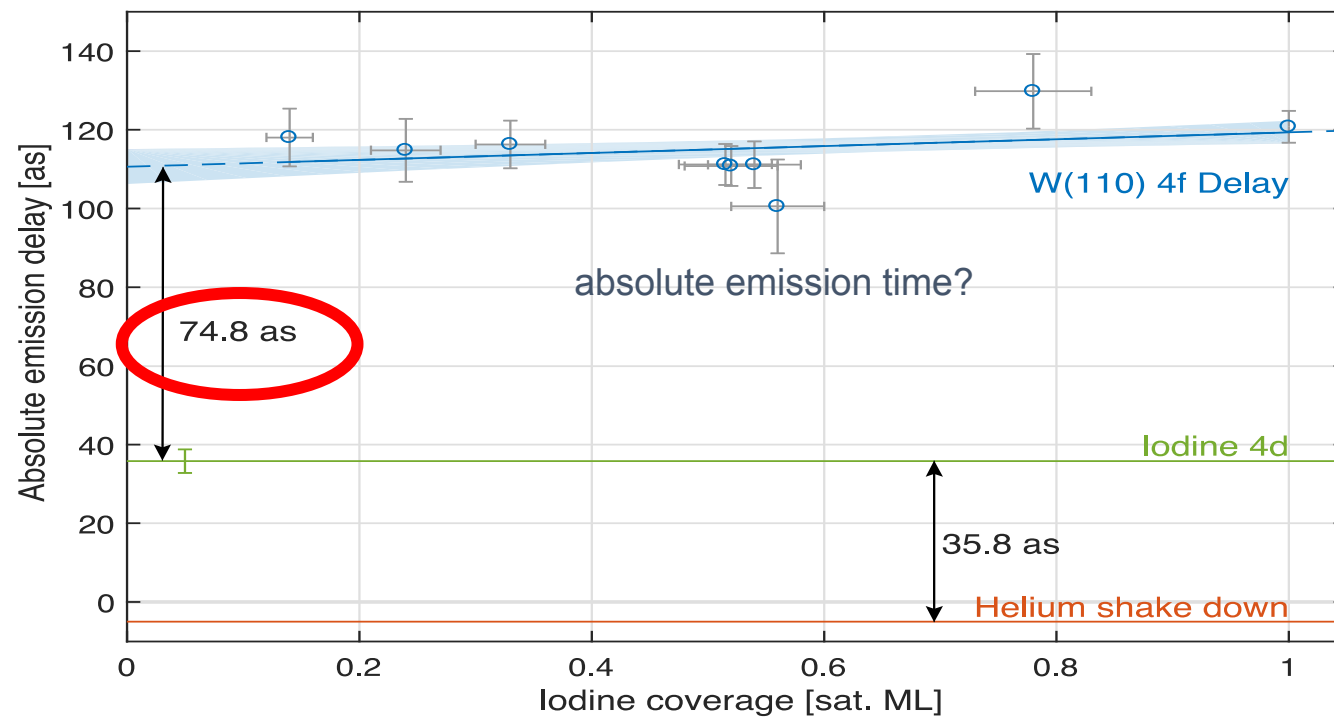


$$\tau_{W4f} = (103 \pm 6)\text{as} + \theta(8 \pm 7)\text{as}$$

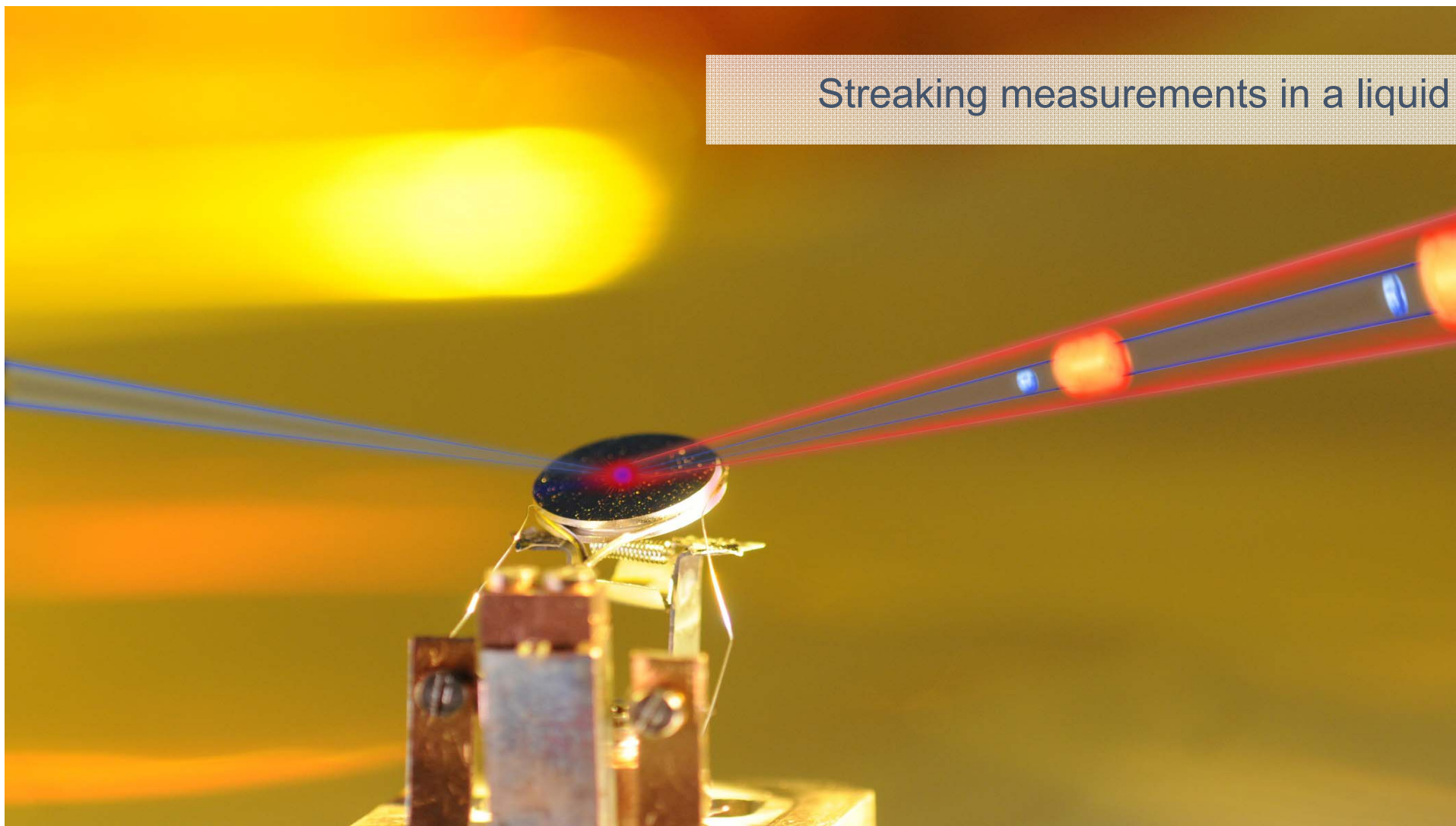




absolute timing of photoemission from surfaces

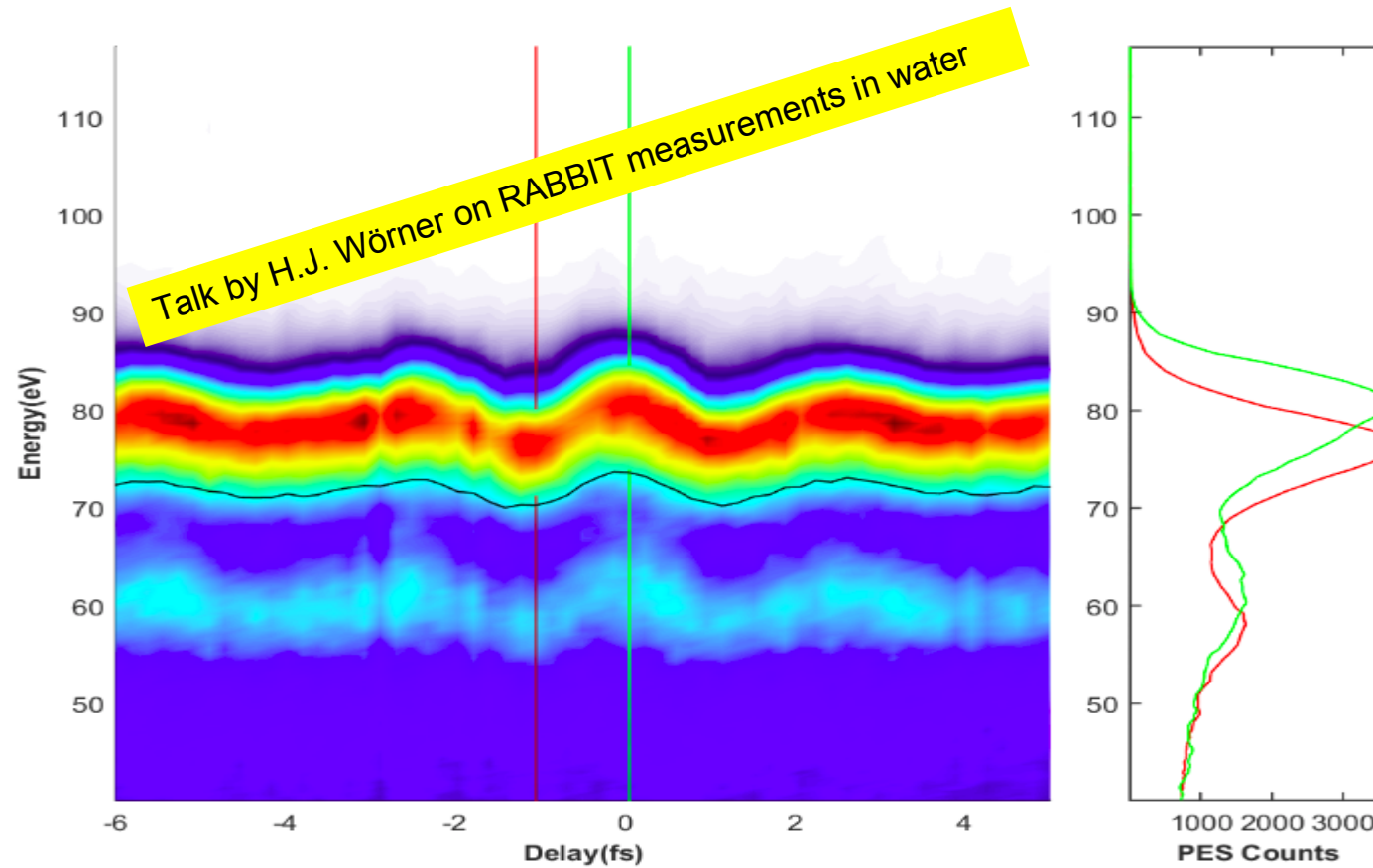


Streaking measurements in a liquid

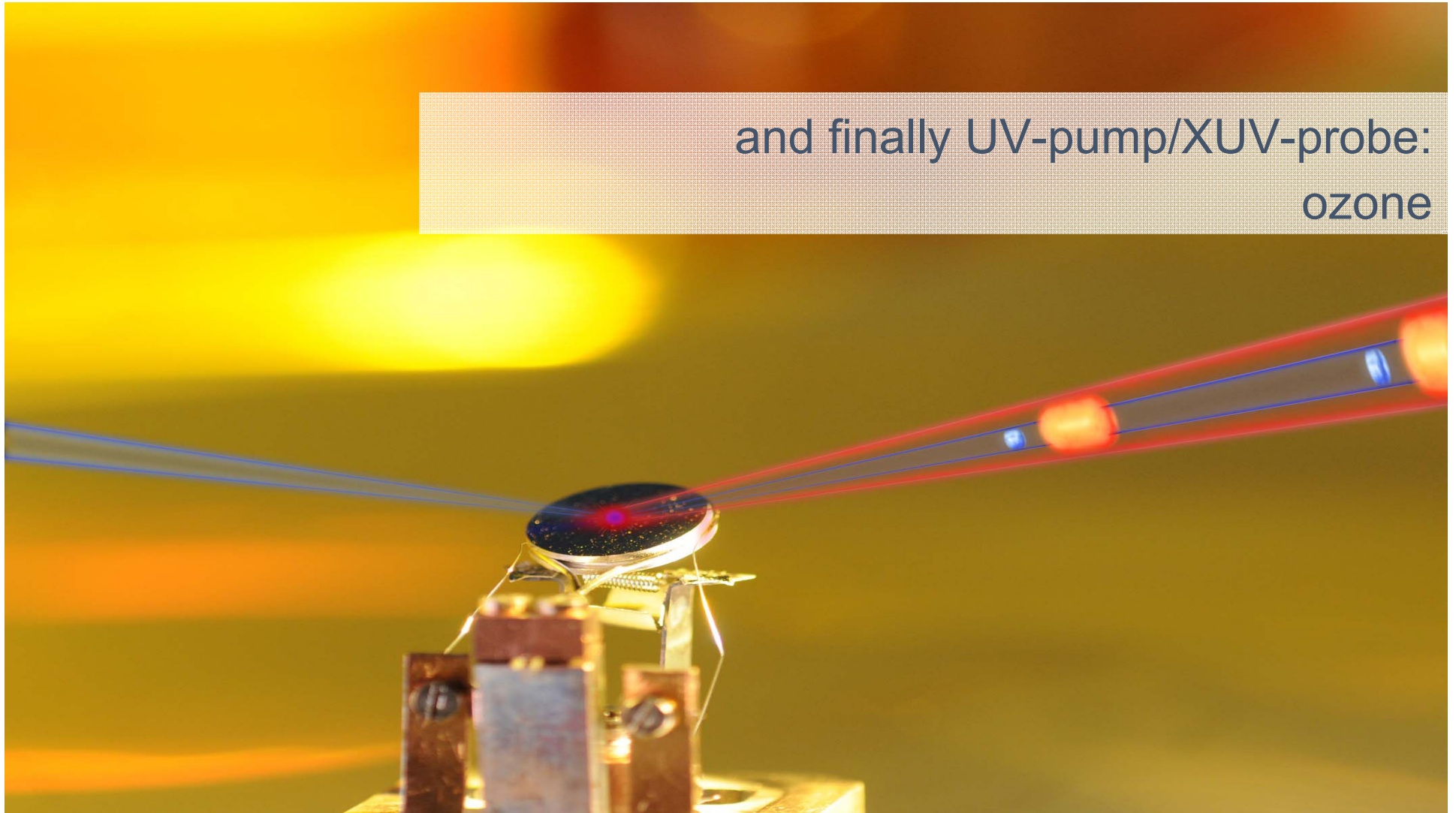




First streaking spectrogram of H₂O



and finally UV-pump/XUV-probe:
ozone

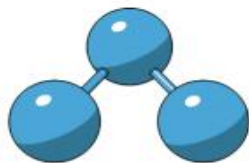




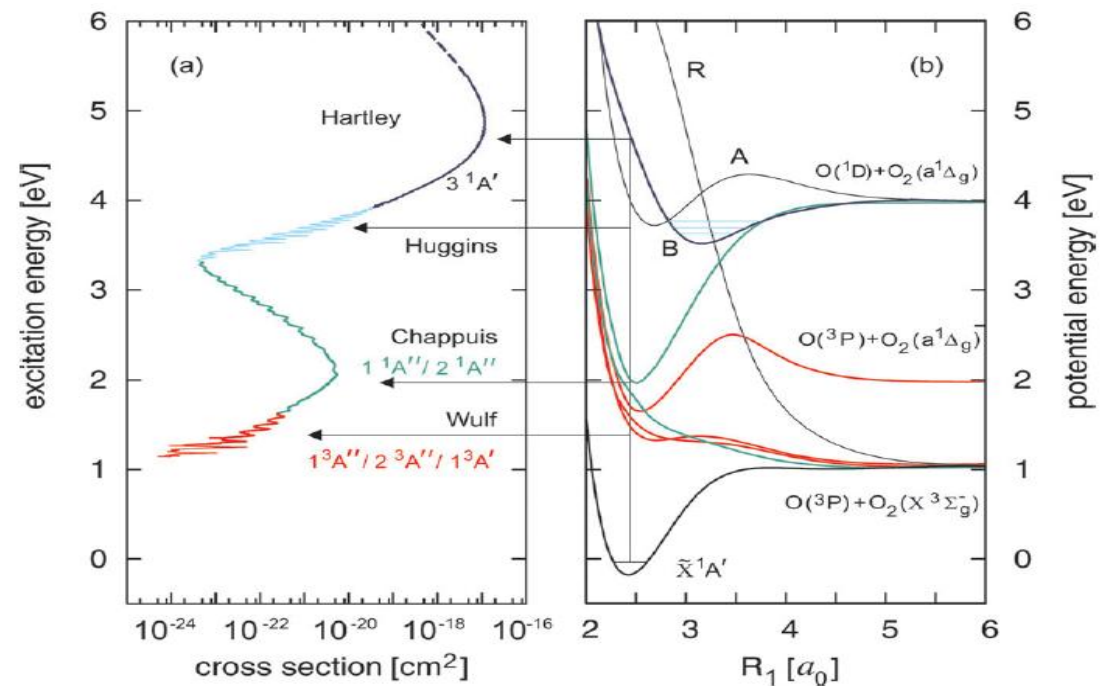
Exciting an electron wavepacket in a molecule UV+NIR pump / XUV probe



Sampling the evolution
with XUV as-pulses

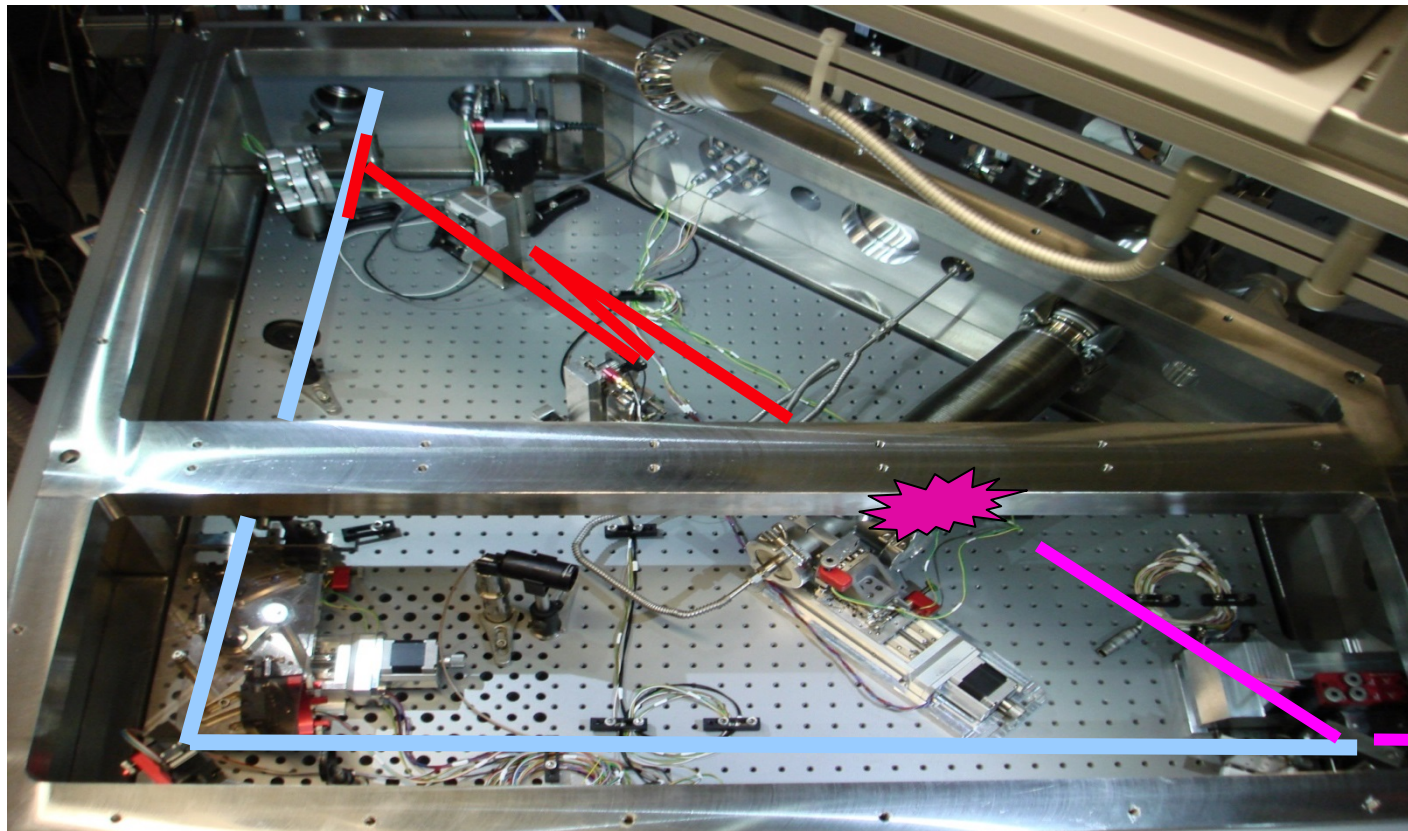


ozone





delay & UV-generation



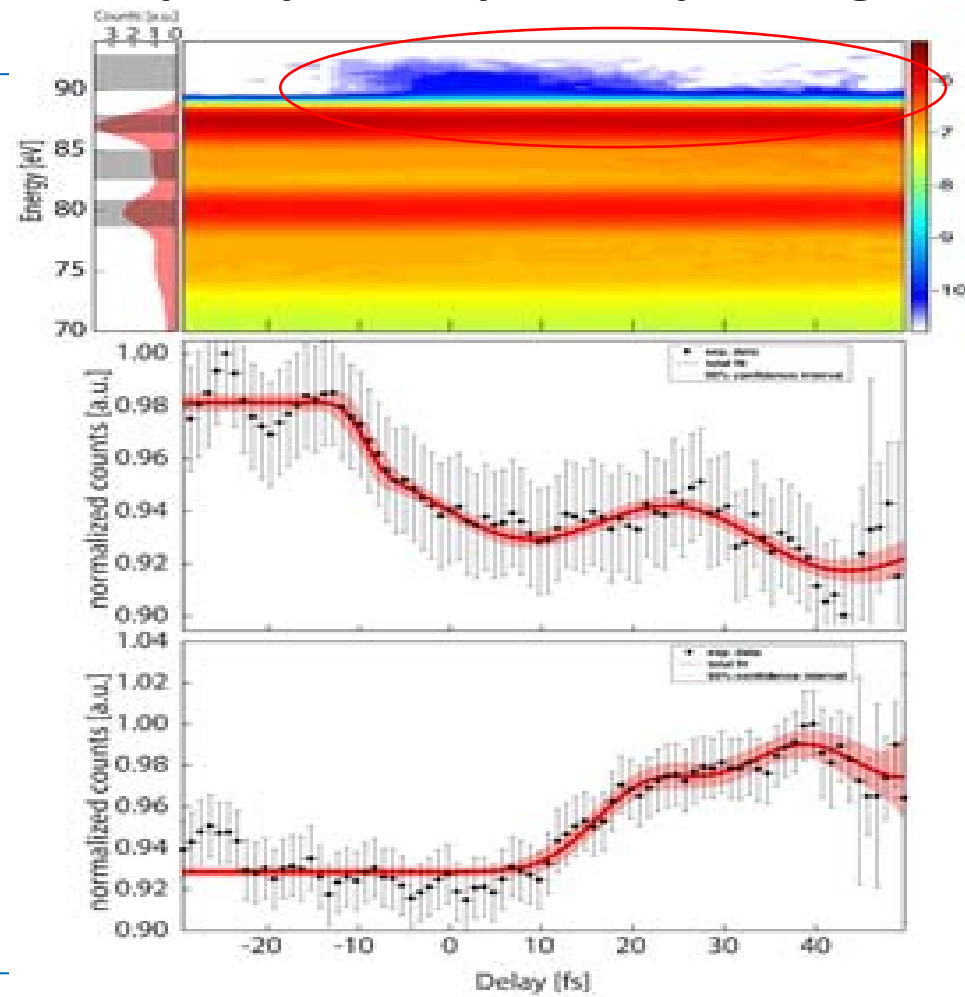


DUV-pump/XUV-probe spectrogramm



ground state

O₂...dissociation

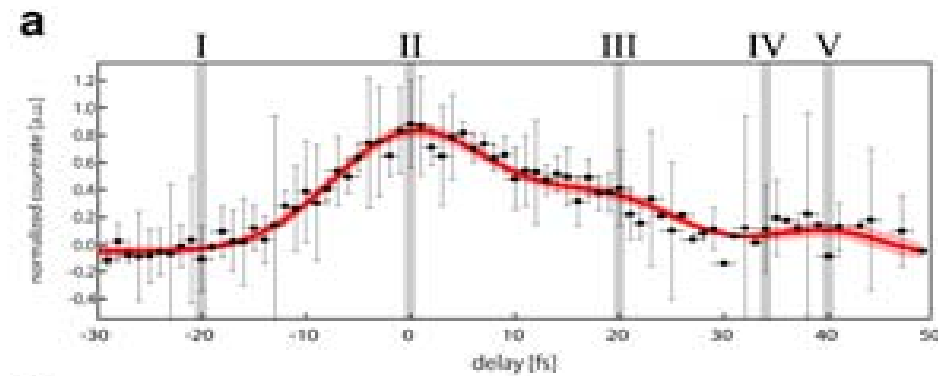




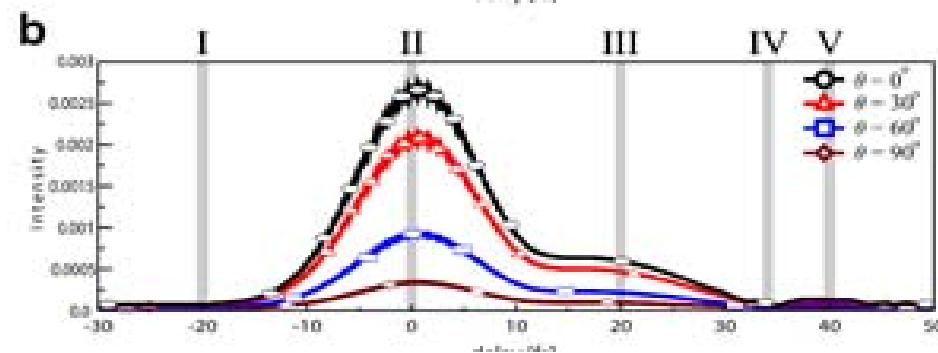
Revival of the excited state



measured



theory





Colleagues & Cooperations



H. Iglev , W. Helml, B. Bernhardt , J. Riemensberger , T. Latka, K. Hütten, R. Heider, D. Hutzler, V. Shirvanyan, K. Stallhofer, M. Wagner, M. Mittermair, K. Oberhofer, D. Potamianos, A. Dünsing

surface dynamics:

M. Schäffer, F. Allegretti, P. Feulner, D. Menzel, J. Barth

TU Munich, Germany

R. Ernstorfer, FHI Berlin, Germany

molecules:

A. Schiffrin, Monash University Australia

as- development:

M. Schultze, N. Karpowicz, F. Krausz, MPQ Garching, Germany

theory:

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Thanks for your attention



attosecond pulses from FELs

